

Operating on the single box this yields  $D \begin{matrix} W_\alpha & Y_\alpha \\ \diagdown & / \\ Z^\alpha & X^\alpha \end{matrix} =$

$$(\partial_X V)^{-1} V Y \left( (\partial_Z \partial_X)^{-1} Z \bar{X} \left( \begin{matrix} / & \diagdown \\ \diagdown & / \end{matrix} - \begin{matrix} / & / \\ \diagdown & / \end{matrix} \right) - 2(Y \partial_Z)^{-1} X Z \left( \begin{matrix} / & \diagdown \\ \diagdown & / \end{matrix} - (\partial_Z)^{-1} X Z \begin{matrix} / & \diagdown \\ \diagdown & / \end{matrix} \right) \right)$$

$$(\partial_X V)^{-1} V Y Z \bar{X} (W Y)^{-1} \left( \begin{matrix} / & \diagdown \\ \diagdown & / \end{matrix} - 2 \begin{matrix} / & / \\ \diagdown & / \end{matrix} + \begin{matrix} / & / \\ / & \diagdown \end{matrix} \right) = 0 \quad (9)$$

The amplitude corresponding to the a.s.d. part of  $F_{ab}$  leads to a twistor diagram independent of  $U^\alpha$  corresponding to the spinor expression

$$\Delta^{A'BB'C'} \phi(k_4) \chi_C(k_3) B_{BB'}(k_2) \psi_{A'}(k_1) = \frac{k_1 k_2 - k_1 k_3}{(k_1 k_2)(k_2 k_3)(k_3 k_1)} k_{1E'} k_3^{DE'} \chi^{A'} \psi_{A'} \Theta_{DB} \phi$$

Similarly in the non-abelian case a perturbative expansion in powers of coupling constants ( $g^n$ )

relates the exterior free fields  $\overset{\circ}{\Psi}_{(A'B')\Theta}^\Phi, \overset{\circ}{\Theta}_{(AB)\Theta}^\Phi$  of order  $g^0$  in

$$\overset{\circ}{F}_{ab\Theta}^\Phi + g \overset{\circ}{F} \dots = \epsilon_{AB} \overset{\circ}{\Psi}_{A'B'\Theta}^\Phi + \epsilon_{A'B'} \overset{\circ}{\Theta}_{AB\Theta}^\Phi + g \epsilon_{AB} \overset{\circ}{\Psi} \dots$$

$$= 2\partial_{[a} A_{b]\Theta}^\Phi - 2ig A_{X[a} A_{b]\Theta}^X \quad (10)$$

linearly to the gauge potentials (again taken to satisfy the Lorenz gauge condition  $\nabla_a A^a_\Theta = \partial_a A^a_\Theta = 0$ ). An analogous construction as in (6) can therefore be used.

One can apply this also to cases of higher helicities, such as for example in [2], as long as one is over flat space. If one looks at these potentials classically one has to ask, however, what their space-time version looks like (i.e. how they are to be 'contour-integrated') and they might turn out not to be very general.

[1] F. Müller (1991). Twistor Diagrams for some scattering amplitudes arising from the Standard Model. Qualifying dissertation.  
 [2] R. Penrose (1990). Twistor Theory for Vacuum Space-Times: a New Approach, T N 31, 6-8.

*Franz Müller*

To appear in: *P. Bergmann's Festschrift*, ed. N. Sanchez:  
**Twistors as Spin 3/2 Charges**

by Roger Penrose  
 Mathematical Institute  
 Oxford, UK

**Abstract** It is pointed out that twistors play a role as the charges for helicity 3/2 massless fields. Since such fields can be defined consistently in general Ricci-flat 4-manifolds, a possible new approach to defining twistors in vacuum