Some new scalar-flat Kähler and hyper-Kähler metrics

We know from the work of Claude Lebrun (J.Diff.Geom. 34 (1991) 223-253) that any scalar-flat Kähler metric with a Killing vector (or 'S'-action') arises from a solution u(x,y,z) of the equation

$$u_{xx} + u_{yy} + (e^{u})_{xx} = 0$$
 1

which is variously known as the 'SU(ω) Toda field equation' or 'Boyer-Finley equation'. What is more Claude tells us in detail how to go from the metric to the function u, though given a solution u there is a choice to be made on the way back to the metric. In particular, given u there is a choice which leads back to a hyper-Kähler metric.

Henrik Pedersen and Yat Sun Poon (Class.Quant.Grav.7 (1990) 1707) found a scalar-flat Kähler metric of Bianchi-type-IX, in the terminology of relativity. This means that the metric has a 3-parameter group of isometries, isomorphic to SU(2) and transitive on 3-surfaces. In particular then the Pedersen-Poon (PP) metric comes from a solution of (1) and it is possible follow Lebrun's direction to find out which. It turns out that the PP-metric arises from an ansatz for (1) which has a simple generalisation to a wider class of solutions of (1). Thus the PP-metric can be generalised to a wider class of scalar-flat Kähler or hyper-Kähler metrics, and this is what I want to describe here.

The idea, with the benefit of hindsight, is to seek a solution of (1) for which u is constant on central ellipsoids. In other words, define u implicitly by the equation

$$Q(u,x,y,z) \equiv X^{t}M(u)X = 1$$

where $X^* = (x,y,z)$ and M(u) is a (symmetric, positive-definite) matrix function of u, to be found. Differentiate (2) implicitly and substitute into (1), then what remains is a second-order matrix ODE in u which can be integrated once. To write the resulting first-order equation out, first define the 3x3 matrix

$$g = diag(1,1,e^{u})$$
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then the equation we want turns out to be

$$VM_{\omega} = MgM$$

using a subscript u to denote $\frac{d}{du}$

where V is an integral of

$$V_{\perp} = \text{Mitrace(gM)}$$
 5

Because of (5), equation (4) is still quite complicated, but there is a dramatic simplification if we work with minus the inverse of M:

set $N = -M^{-1}$ then (4) and (5) reduce to

$$VN_{\perp} = g$$
 6

and

$$V_{ij} = -Metrace(gN^{-1}).$$
 7

From (3) and (6), the off-diagonal terms in N are constant so set

$$N = \begin{bmatrix} a & v & \mu \\ v & b & \lambda \\ \mu & \lambda & c \end{bmatrix}$$

where λ, μ, ν are constants, then (6) in components reduces to the 3 equations

$$a_{u} = b_{u} = 1/V$$
; $c_{u} = e^{u}/V$ 9

At once, from (9)

$$a-b = 2\zeta$$
 constant,

where we retain the convention of using Greek letters for constants. Define a new constant η by

$$\eta^2 = \zeta^2 + v^2$$

then a convenient parametrisation, in terms of a new variable w, turns out

$$a = \eta \left(\frac{\omega+1}{\omega-1}\right) + \zeta$$
 $b = \eta \left(\frac{\omega+1}{\omega-1}\right) - \zeta$ 10

Next we note from (7) that

$$V^2$$
detN = ξ constant 11

Our aim is to obtain an equation for w. We solve (10) and (11) for c in terms of w, V and constants, then eliminate V in favour of w_ using (9) and (10). This gives c in terms of w, w_ and constants and then the last part of (9) gives a second-order ODE for w. A change of independent variable

$$z = e^{u}$$
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enables this final ODE to be written as

$$\frac{d^2w}{dz^2} = \left(\frac{1}{2w} + \frac{1}{w-1}\right) \left(\frac{dw}{dz}\right)^2 - \frac{1}{z} \frac{dw}{dz} + \frac{(w-1)^2(\alpha w + \beta)}{z^2} + \frac{\gamma w}{w} + \frac{\delta w(w+1)}{z}$$
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where $\alpha,\beta,\gamma,\delta$ are constants which can be expressed in terms of the constants which we already have ie $\lambda,\mu,\nu,\xi,\eta,\zeta$. Painlevé-buffs will recognise the ODE in (13) as Painlevé-V (P-V for short), with the usual conventions. In fact, the constant δ turns out to be zero in this case, and Fokas and Ablowitz (J.Math.Phys. 23 2033 (1982)) explain how to transform this special case of P-V to P-III.

Given a solution w(z) of (13), we can work our way back to the matrix N and so to M and to u(x,y,z). To be able to write down a scalar-flat Kähler metric, we have still to solve a linear PDE, but to write down a hyper-Kähler metric, there are no more choices to be made: given u(x,y,z) the metric is immediate. In particular, given the function u(x,y,z) appropriate to the PP-metric, there is a hyper-Kähler metric with the same u (apparently a new one, but it won't be of Bianchi-type-IX unless its already known!).

The special case of this procedure when the matrix M in (2) is diagonal, or equivalently when the constants λ,μ,ν in (8) are all zero, is the one which leads to the PP-metric, which was already known to involve P-III. What is obscure in this approach is why the PP solution of (1), which must lead to a scalar-flat Kähler metric with an S'-action, actually leads to a far more symmetric metric with an SU(2)-action. This raises the possibility that these more general solutions of (1) also lead to scalar-flat Kähler metrics with SU(2)-action. If these exist, and I am grateful to Andrew Dancer for the suggestion that they do, then their metrics will look like

$$ds^2 = Fdt^2 + A_{i,j}\sigma^i\sigma^j$$

where $A_{i,j}$ is a matrix function of t, and the σ^i are a basis of left-invariant one-forms for S^3 . The question is whether one can relate the matrix A to the matrix M of (2). Certainly for the PP-metric, when M and A are both diagonal and t in (14) is proportional to u, this is possible.

One could also try different generalisations of the ansatz (2): for example with hyperboloids in place of ellipsoids (ie with an indefinite M(u) rather than a positive-definite one) which presumably leads to a Bianchi-type-VIII metric, or with paraboloids or even with planes (This case, which is fairly easy to do, leads to a Riccati equation in place of (13)). And, of course, free with every solution of (1) one obtains a hyper-Kähler metric too.

Acknowledgement

I gratefully acknowledge that this calculation was precipitated by Andrew Dancer's suggestion to me that the PP-metrics have non-diagonal generalisations.

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