

Apolarity between  $P_1, \dots, P_n$  and  $Q_1, \dots, Q_n$  is now simply the condition that  $P_1, \dots, P_{n-1}$  and  $R_1, \dots, R_{n-1}$  be apolar.

Proof Consider  $\overset{n}{P} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ C \\ \vdots \\ X^N \end{pmatrix} \overset{1}{Q} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ C \\ \vdots \\ Q_N \end{pmatrix} = 0$ . suffices have now become numbers above

Solutions for  $X^A$  give  $\overset{1}{R} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ C \\ \vdots \\ R_N \end{pmatrix}$ . Hence

$$\overset{n}{P} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ C \\ \vdots \\ Q_N \end{pmatrix} = \overset{1}{R} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ C \\ \vdots \\ R_N \end{pmatrix} \leftarrow \text{taking multiplying factor} = 1$$

Apolarity between  $P_1, \dots, P_{n-1}$  and  $R_1, \dots, R_{n-1}$  is

$$\overset{1}{P} \begin{pmatrix} B \\ \vdots \\ P^N \end{pmatrix} \overset{1}{R} \begin{pmatrix} B \\ \vdots \\ R^N \end{pmatrix} = 0$$

which is the same as the required condition

$$\overset{n}{P} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ P^N \end{pmatrix} \overset{1}{Q} \begin{pmatrix} A \\ \vdots \\ B \\ \vdots \\ C \\ \vdots \\ Q_N \end{pmatrix} = 0.$$

Q.E.D.

### A Simple Observation Concerning {22} Vacuums

It is well known that every {22} vacuum ("type D"), with Weyl spinor

$$\Psi_{ABCD} = M r^{-3} \alpha_A \alpha_B \beta_C \beta_D$$

(where  $M$  is a constant and  $\alpha_A \beta^A = 1$ ) possesses a Killing spinor - a valence-2 symmetric twistor! -

$$\chi_{AB} = r \alpha_{(A} \beta_{B)} \quad \text{satisfying} \quad \nabla_{A'} \chi_{BC} = 0.$$

(See Walker & Penrose, *Comm. Math. Phys.* **18** (1970) 265-74; also *Spinors & Space-Time* Vol. 2, p. 107.) In the Schwarzschild solution,  $r$  is the standard "radial coordinate", but in general  $r$  is a complex "radial" quantity.

I am not aware that anyone has pointed out the following simple but striking consequence:

Proposition Along every null geodesic, with parallelly propagated tangent spinor  $O^A$ , the null datum

$$\Psi_0 = \Psi_{ABCD} O^A O^B O^C O^D = C_{abcd} l^a m^b \bar{l}^c m^d$$

has the precise form

$$\Psi_0 = \frac{K}{r^5}$$

where  $K$  is a complex constant depending on the choice of null geodesic (with its choice of scale for  $O^A$ ).

Proof This is an immediate consequence of the constancy of  $\chi_{AB} O^A O^B$  along the null geodesic ( $O^A \bar{O}^{A'} \nabla_{AA'} (\chi_{BC} O^B O^C) = 0$ ) and the fact that  $\Psi_0 = M r^{-3} (\alpha_A O^A)^2 (\beta_B O^B)^2$  and  $\chi_{AB} O^A O^B = r (\alpha_A O^A) (\beta_B O^B)$ , whence  $\Psi_0 = M / (\chi_{AB} O^A O^B)^2 r^5$ . Q.E.D.

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