

Outline:

I Localization in Twistor Space

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hep-th/0404095

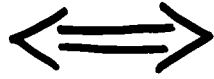
II Topological M-theory

R. Dijkgraaf, S.G., A. Neitzke, ...

hep-th/0411072

Witten:

perturbative $\mathcal{N}=4$
super Yang-Mills



Topological String
(B-model)
in twistor space $\mathbb{C}\mathbb{P}^{3|4}$

$$\bullet \mathbf{Z}^A := (\lambda^1, \lambda^2, \mu^1, \mu^2 \mid \psi^1, \psi^2, \psi^3, \psi^4) \in \mathbb{C}^{4|4}$$

holomorphic 4/4 - form :

$$\Omega = d\lambda^1 d\lambda^2 d\mu^1 d\mu^2 d\psi^1 d\psi^2 d\psi^3 d\psi^4$$

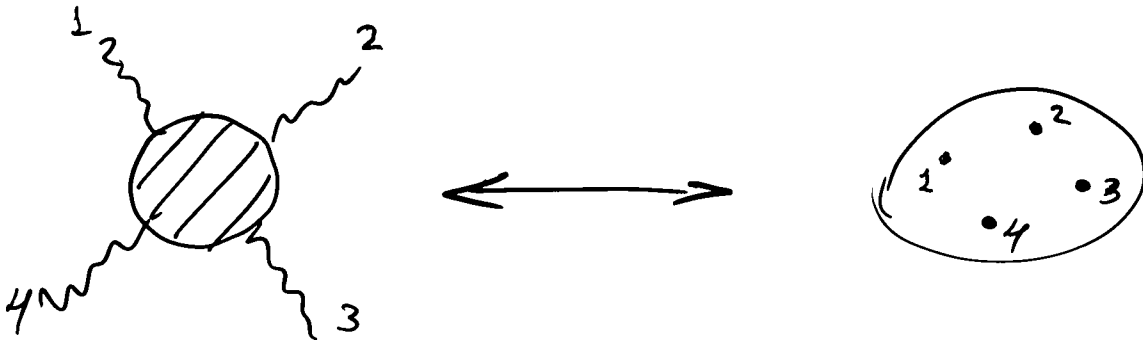
identification :

$$(\lambda, \mu \mid \psi) \sim (t\lambda, t\mu \mid t\psi) \quad t \in \mathbb{C}^*$$

\Rightarrow Calabi-Yau supermanifold $\mathbb{C}\mathbb{P}^{3|4}$

$$\mathcal{A} \left(\underbrace{+++}_{n} \overbrace{---}^q \right) = \int \omega$$

moduli space of hol. curves
in $\mathbb{C}P^{3|4}$ of degree $d=q-1$
with n marked points

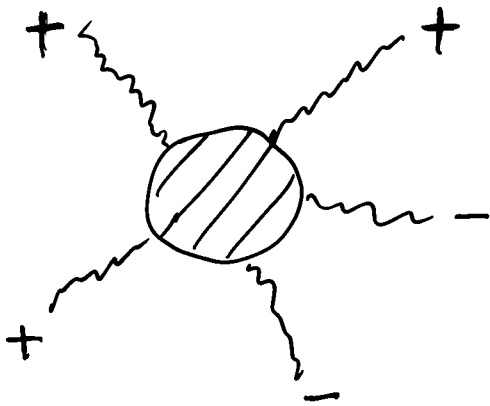


Remark: in the B-model, hol. curves
are interpreted as D1-brane instantons.

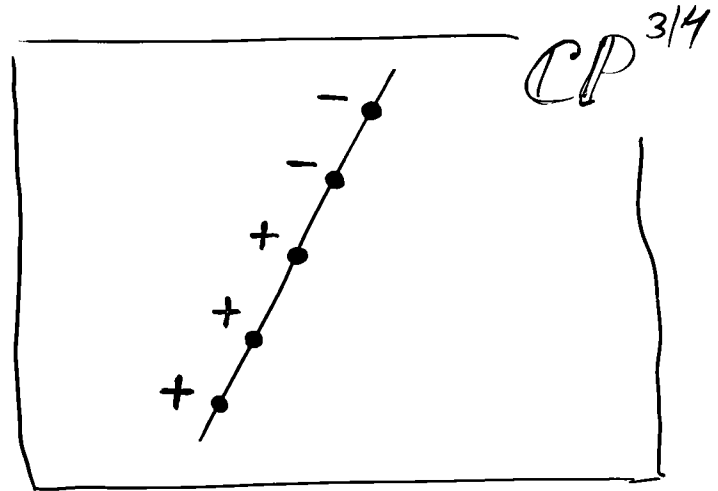
Example 1:

$$A(+++--)$$

$$d = q - 1 = 1$$



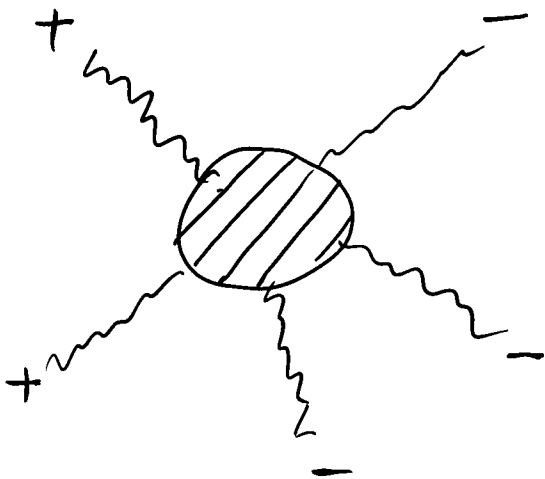
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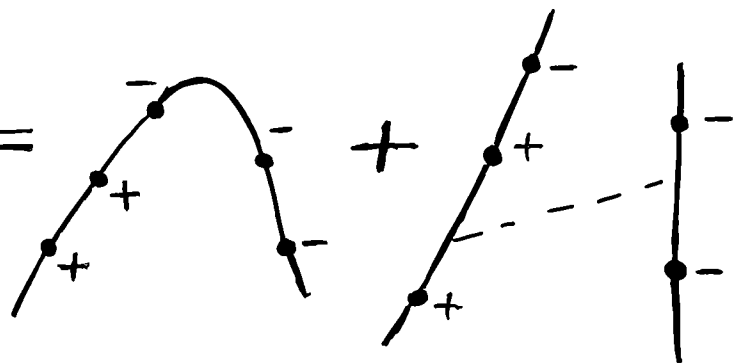
Example 2:

$$A(++----)$$

$$d = 2$$



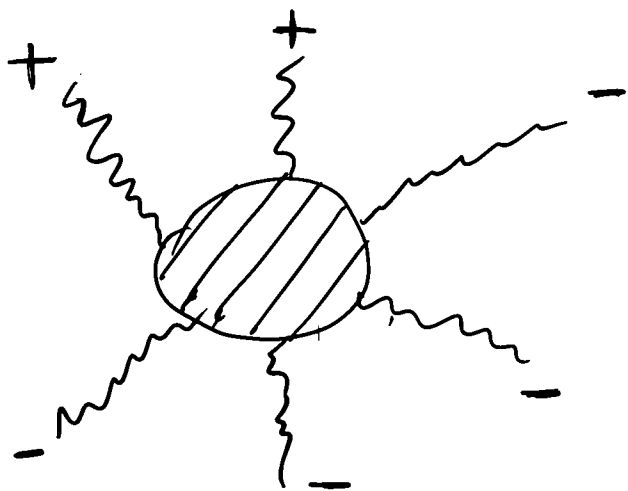
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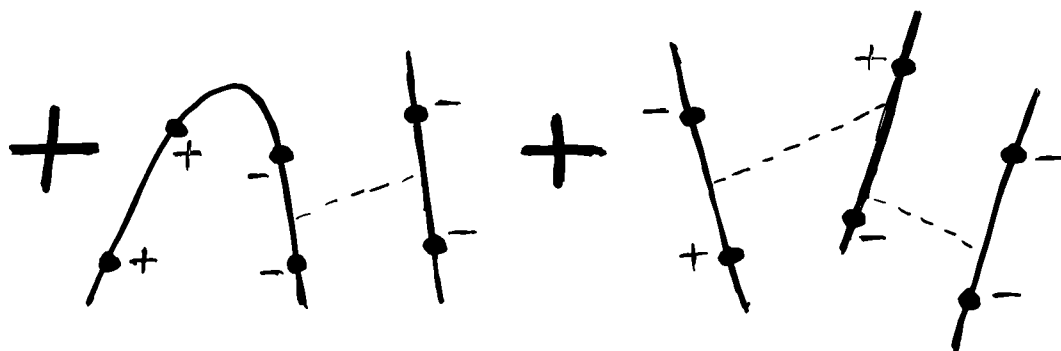
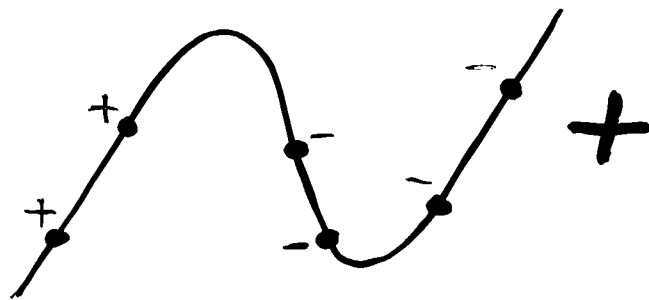
Example 3:

$A(++----)$

$d=3$



$=$



Connected Instantons:

- $\mathcal{M}_{0,n,d}(\mathbb{C}P^{3/4})$ moduli space of genus 0, n -pointed curves of degree d in $\mathbb{C}P^{3/4}$

- $\dim \mathcal{M}_{0,n,d}(\mathbb{C}P^{3/4}) = (4d+n) - (4d+4)$

- Following [Witten], we realize it as the space of automorphism classes of maps

$$p: \mathbb{C}P^1 \rightarrow \mathbb{C}P^{3/4}$$

- $\mathcal{M}_{0,n,d}(\mathbb{C}P^{3/4})$ is non-compact due to deg. curves, e.g.



\Rightarrow moduli space of stable maps

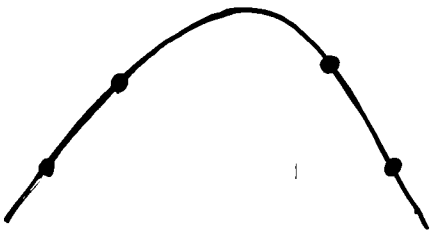
$$\mathcal{M}_{0,n,d}(\mathbb{C}P^{3/4})$$

$$\begin{cases} P: \mathbb{C}P^1 \rightarrow \mathbb{C}P^{3/4} \\ P^A(\sigma) = \sum^A = \sum_{k=0}^d \beta_k^A \sigma^k \end{cases}$$

moduli $\beta_k^A \in \mathbb{C}^{4d+4/4d+4}$

natural measure

$$\mu_d = \prod_{k,A} d\beta_k^A$$



$$\mathcal{A}_{YM}^{\text{conn}} = \int_{\mathcal{M}_{0,n,d}} \prod_{k,A} d\beta_k^A \cdot \underbrace{\prod_{i=1}^n \frac{d\sigma_i}{\sigma_i - \sigma_{i+1}}}_{\text{hol. } n\text{-form on } (\mathbb{C}P^1)^n} \prod_{i=1}^n ev_i^*(\phi_i)$$

$\omega(\sigma_i)$

- ϕ_i $\bar{\partial}$ -closed $(0,1)$ -form on $\mathbb{C}P^{3/4}$

• $GL(2, \mathbb{C}) : \sigma \mapsto \sigma' = \frac{a\sigma + b}{c\sigma + d}$

$$ad - bc \neq 0$$

$$\Rightarrow \mathcal{A}_{YM}^{\text{conn}} = \int_{\mathcal{M}_{0,n,d}} \frac{\prod_{K,A} d\beta_K^A \prod_{i=1}^n \frac{d\sigma_i}{\sigma_i - \sigma_{i+1}}}{GL(2, \mathbb{C})} \wedge \prod_{i=1}^n e_{V_i}^*(\phi_i)$$

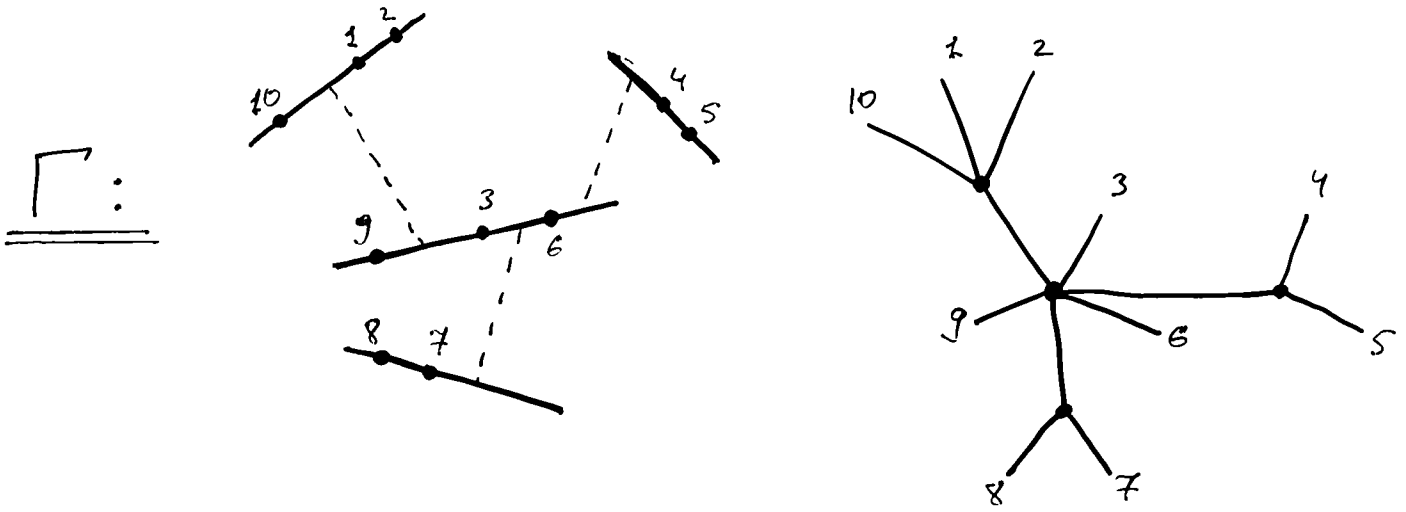
$$= \int_{\mathcal{M}_{0,n,d}} \frac{\mu_d \wedge \omega(\sigma_1, \dots, \sigma_n)}{GL(2, \mathbb{C})} \wedge \Phi$$

* Surprise:

$$\mathcal{A}_{YM}^{\text{conn}} = \mathcal{A}_{YM}$$

R. Roiban,
M. Spradlin,
A. Volovich

Completely Disconnected Instantons



- $Q_i: \mathbb{C}P^1 \rightarrow \mathbb{C}P^{3/4}$

$$Q_i^A(\sigma) = \sum_{k=0}^1 \beta_{k,i}^A \sigma^k$$

- moduli $\beta_{k,i}^A \in \mathbb{C}^{8d|8d}$

- $GL(2, \mathbb{C})^d \Rightarrow 4d | 8d$

$$\mathcal{M}_{\text{lines}}^\Gamma = \prod_{i=1}^d \mathcal{M}_{0,n,2}(\mathbb{C}P^{3/4})$$

(Holomorphic) Measure:

$$i) \mu_{\text{lines}} = \prod_{K, A, i} d\beta_{K, i}^A$$

$$ii) \omega_i := \omega(\sigma_1, \dots, \sigma_{n_i})$$

$$iii) \Phi = \prod_i \text{ev}_i^*(\phi_i)$$

$$iv) D(x, y) \text{ (0,2)-form on } \mathbb{C}P^{3/4} \times \mathbb{C}P^{3/4}$$

“holomorphic CS propagator”

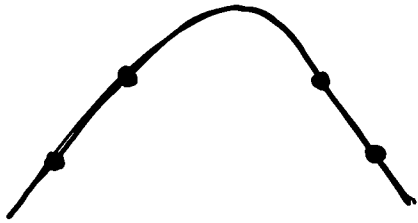
$$\mathcal{A}_{\text{YM}}^{\text{disc}} = \sum_{\Gamma} \int_{\mathcal{M}_{\text{lines}}^{\Gamma}} \frac{\mu_{\text{lines}} \wedge \left(\prod_{i=1}^d \omega_i \right) \wedge \Phi \wedge D}{(\text{GL}(2, \mathbb{C}))^d}$$

* Yet another surprise:

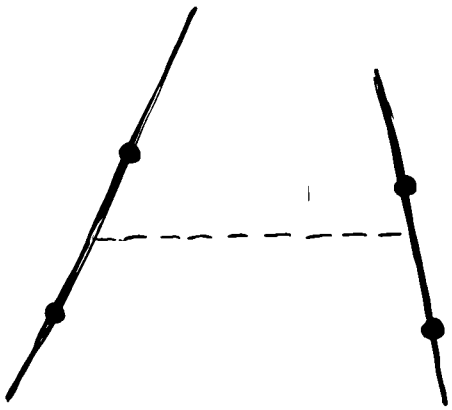
$$\mathcal{A}_{\text{YM}}^{\text{disc}} = \mathcal{A}_{\text{YM}}$$

F. Cachazo,
P. Svrcek,
E. Witten.

Puzzle:

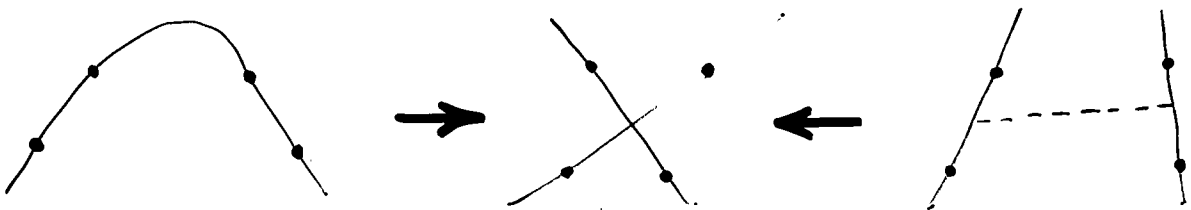


$$A_{YM} = \frac{\int \omega_{\text{conn}}}{\mathcal{M}_{0,n,d}}$$



$$A_{YM} = \frac{\int \omega_{\text{disc}}}{\mathcal{M}_{\text{lines}}}$$

- Both integrals have a simple pole along \mathcal{M}_{int} :

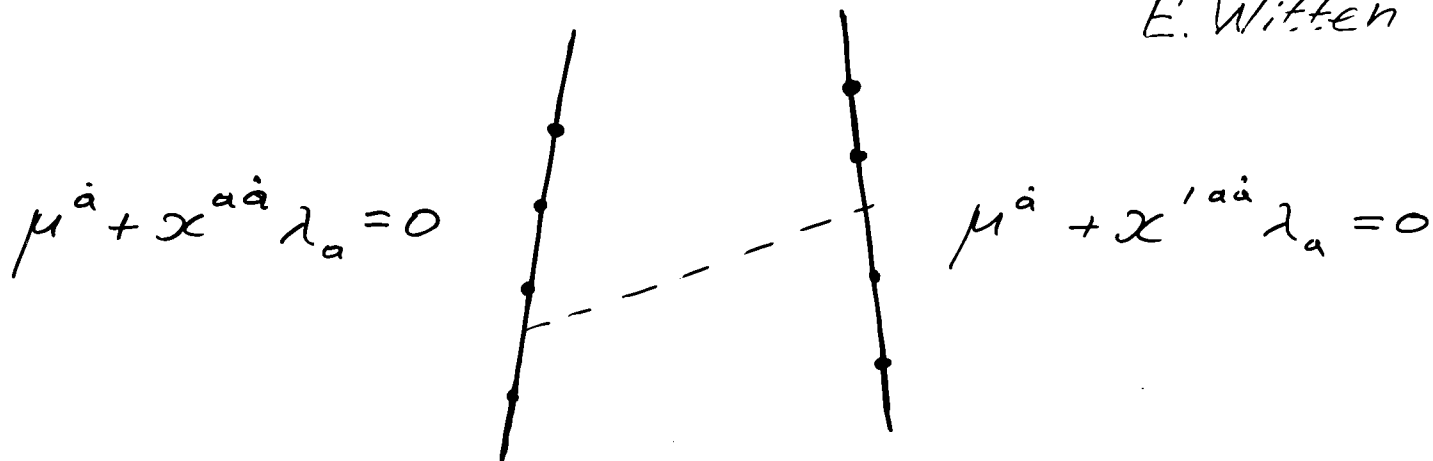


- same residue \Rightarrow

$$\frac{\int \omega_{\text{conn}}}{\mathcal{M}_{0,n,d}} = \frac{\int \omega_{\text{disc}}}{\mathcal{M}_{\text{lines}}}$$

• Residue from ω_{disc}

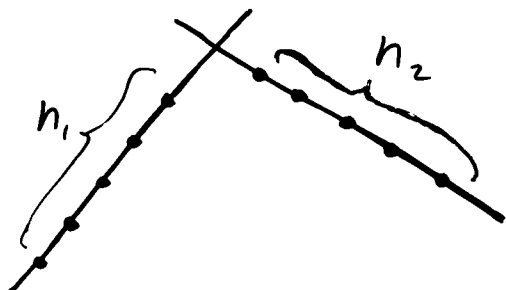
F. Cachazo, P. Srceek,
E. Witten



$$y^{aa} := x'^{aa} - x^{aa} \Rightarrow$$

$$\omega_{disc} \propto \frac{d^4 y}{y^{aa} y'^{aa}} = \frac{dy^{11} dy^{22} dy^{12} dy^{21}}{2(y^{11} y^{22} - y^{12} y^{21})}$$

$$\Rightarrow \sum_{\Gamma} \int_{\mathcal{M}_{int}^{\Gamma}} \frac{\mu_{int} \wedge \prod_{i=1}^{n_1} \frac{d\sigma_i}{(\sigma_i - \sigma_{i+1})} \wedge \prod_{i=1}^{n_2} \frac{d\sigma_i}{(\sigma_i + \sigma_{i+1})}}{\text{vol}(GL(2, \mathbb{C}))^2} \wedge \Phi$$

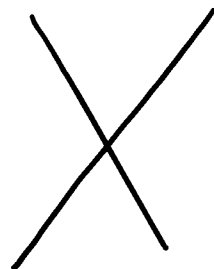


- Residue from ω_{conn}

ε deformation parameter

$\varepsilon=0$:

$$\boxed{xy = \varepsilon}$$



- $\omega_{\text{conn}} = \frac{\mu_d \wedge \omega(\sigma_1, \dots, \sigma_n)}{\text{vol}(GL(2, \mathbb{C}))} \wedge \Phi$

- $\mu_d \sim \frac{d\varepsilon}{\varepsilon^3}$

- $\omega(\sigma_1, \dots, \sigma_n) \sim \varepsilon^2$

\Rightarrow $\boxed{\omega_{\text{conn}} \sim \frac{d\varepsilon}{\varepsilon}}$

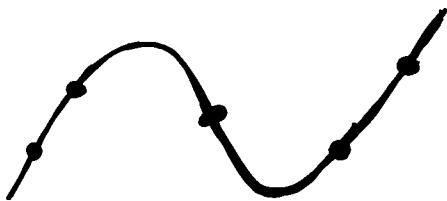
- simple pole if the distribution of marked points into two groups preserves the cyclic order

- same residue as in the disconnected case

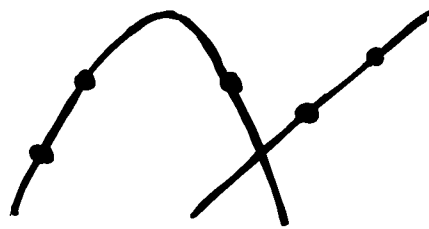
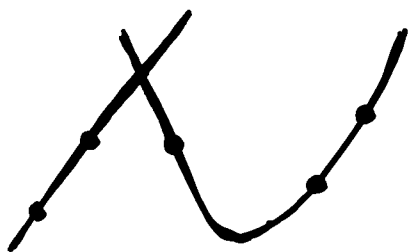
Higher Degree Curves and

Intermediate Prescriptions

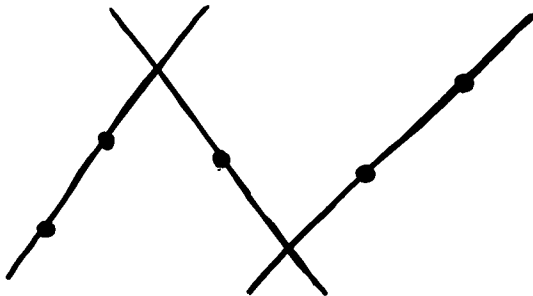
$\overline{\mathcal{M}}_{0,n,d}$



$\mathcal{M}_{int}^{\Lambda_1} \cup \mathcal{M}_{int}^{\Lambda_2}$



$\mathcal{M}_{int}^{\Lambda_1} \cap \mathcal{M}_{int}^{\Lambda_2}$



Open Questions

- Choice of contour / integral
- External wavefunctions
- Loops and higher genus
- Choice of prescriptions
- Recursion relations
- \vdots
- Physical Interpretation?
- Dual String Theory?

Many Faces of Topological

M - theory

- Geometry of $\mathcal{N} = 1$ String Vacua
- Form Theories of Gravity in
Diverse Dimensions
- Black Hole Entropy & Attractors
- Topological String Theories
and Dualities Between Them

Special Holonomy Manifolds

X n -dim'n'l Riemannian manifold

$$\text{Hol}(g_X) \subseteq \text{SO}(n)$$

Metric	Holonomy	n	\mathcal{NUSY} $\nabla\zeta = 0$	Invariant p -forms
Calabi-Yau	$SU(3)$	6	$\frac{1}{4}$	$p=2: K$ $p=3: \Omega$
	G_2	7	$\frac{1}{8}$	$p=3: \Phi$ $p=4: G = * \Phi$
Exceptional	$Spin(7)$	8	$\frac{1}{16}$	$p=4: \Psi$

Invariant Forms

\exists covariantly constant spinor ($\nabla \zeta = 0$)

$$\omega^{(p)} = \zeta^\dagger \Gamma_{i_1 \dots i_p} \zeta$$

$\omega^{(p)}$ covariantly constant
invariant under $\text{Hol}(g_X)$

- \mathcal{S} minimal (supersymmetric) cycle
 $\dim(\mathcal{S}) = p$

$$\Rightarrow \text{Vol}(\mathcal{S}) = \int_{\mathcal{S}} \omega^{(p)}$$

Example:

$$\text{Hol}(g_X) = G_2 \Leftrightarrow \begin{cases} d\Phi = 0 \\ d*\Phi = 0 \end{cases}$$

$$\Rightarrow ds^2 = \sum_{i=1}^7 e^i \otimes e^i$$

$$\Rightarrow \Phi = \frac{1}{3!} \psi_{ijk} e^i \wedge e^j \wedge e^k$$

$$\sigma_i \sigma_j = -\delta_{ij} + \psi_{ijk} \sigma_k, \quad \sigma_i \in \text{Im}(\mathcal{D})$$

$$\Leftarrow B_{jk} = -\frac{1}{144} \Phi_{j_1 i_1 i_2} \Phi_{k_1 i_2 i_3} \Phi_{i_3 i_4 i_5} \Sigma^{i_1 \dots i_7}$$

$$g_{ij} = \det(B)^{-1/9} B_{ij}$$

$$\rightsquigarrow *\Phi, \text{ etc.}$$

Form Theories of Gravity

in D dimensions

$D=2$: 2D Topological Gravity

$D=3$: Chern-Simons Gravity

$D=4$: 2-form Gravity

(self-dual sector of loop quantum gravity)

$D=6$:
{ Kähler Gravity (A-model) $dK=0$
Kodaira-Spencer Gravity (B-model) $d\Omega=0$

!Mitchin:
{ 4-form theory $dK=0$
3-form theory $d\Omega=0$

4D 2-form Gravity

M 4-manifold

A $SU(2)$ gauge connection

$$\Sigma \in \Omega^2(M) \otimes \text{adj}(SU(2)) \quad k=1, 2, 3$$

$$\boxed{S_{4D} = \int_M \Sigma^k \wedge F_k - \frac{\Lambda}{24} \Sigma^k \wedge \Sigma_k}$$

$$\begin{cases} D_A \Sigma = 0 \\ F_k = \frac{\Lambda}{12} \Sigma_k \end{cases}$$

$$\Rightarrow \sqrt{g} g_{ab} = -\frac{1}{12} \Sigma_{aa_1}^i \Sigma_{ba_2}^j \Sigma_{a_3 a_4}^k \epsilon^{ijk} \epsilon^{a_1 \dots a_4}$$

self-dual Einstein:

$$R_{ab} = \Lambda g_{ab}, \quad W_{abcd}^{(+)} = 0$$

3-form theory in $D=7$

X 7-dimn'l manifold

N. Hitchin

$$\Phi \in \Omega^3(X)$$

• action:

$$S(\Phi) = \frac{1}{2} \int_X \Phi \wedge *_{\Phi} \Phi = \int_X \sqrt{g_{\Phi}} =$$

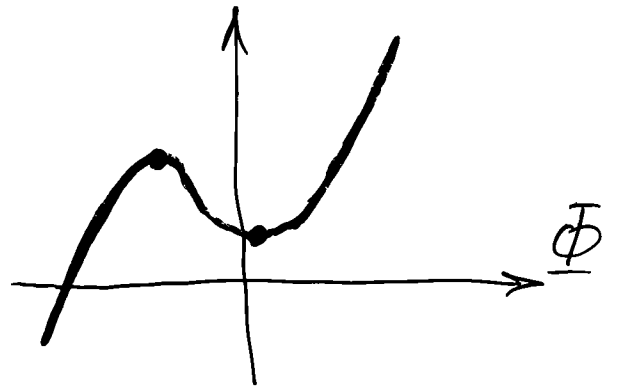
$$= \int_X (\det B)^{1/9}$$

homog. of degree $7/3$ in Φ

• assume $d\Phi = 0$, fix $[\Phi] \in H^3(X)$

vary $\Phi = \Phi_0 + dB$

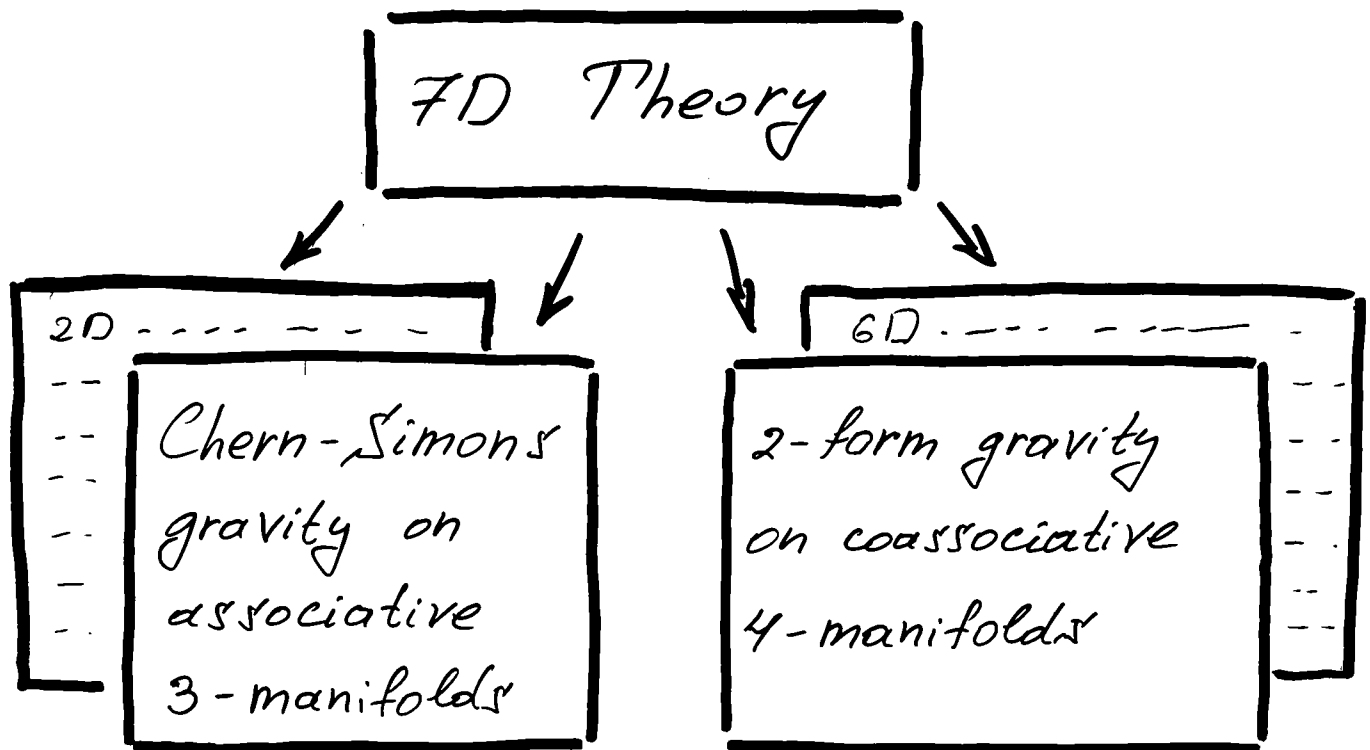
• crit. pts:
$$\begin{cases} d\Phi = 0 \\ d*\Phi = 0 \end{cases}$$



$$\Rightarrow \sqrt{g} g_{jk} = -\frac{1}{144} \Phi_{j_1 i_1 i_2} \Phi_{k_1 i_3 i_4} \Phi_{i_5 i_6 i_7} \varepsilon^{i_1 \dots i_7}$$

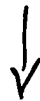
metric with G_2 holonomy!

- FD theory on X reduces to n -dimensional form theory of gravity on a calibrated cycle $M \subset X$, $\dim(M) = n$



"local model":

$$\mathbb{R}^m \rightarrow \bar{X}$$



$$M^n$$

$$m+n=7$$

- $SO(m)$ - invariant ansatz:

y^i local coordinate on \mathbb{R}^m

$$\alpha^i := D_A y^i = dy^i + (A y)^i$$

Example ($m=3, n=4$):

$$\Phi = \alpha^{123} + \alpha^1 \wedge \Sigma^1 + \alpha^2 \wedge \Sigma^2 + \alpha^3 \wedge \Sigma^3$$

$$\begin{cases} d\Phi = 0 \\ d*\Phi = 0 \end{cases} \Rightarrow \begin{cases} D_A \Sigma = 0 \\ F = \frac{\Lambda}{12} \Sigma \end{cases}$$

$\sqrt{g} g_{ij}$ along M reduces to

$$\sqrt{g} g_{ab} = -\frac{1}{12} \sum_{aa_1}^i \sum_{ba_2}^j \sum_{a_3 a_4}^k \varepsilon^{ijkl} \varepsilon^{a_1 \dots a_4}$$

- $\mathbb{R}^3 \rightarrow X$
 \downarrow
 M^4
- bundle of self-dual 2-forms on M^4 , $X = \Lambda_+^2(M^4)$
 $\hookrightarrow X =$ bundle of self-dual 2-forms of unit norm
 $= T(M^4)$

- In fact, g_X is asymptotic to a cone on $T(M)$

$T(M^4)$ has integrable eplx. structure $\iff M^4$ self-dual
 Penrose,
 Atiyah-Hitchin-Singer

$T(M^4)$ Kähler $\iff M^4$ Einstein
 Hitchin

Conj: topological gauge theory on $M^4 \cong$ topological gravity ("M-theory") on $X \setminus M^4$
 SI
 line bundle over $T(M^4)$

$$\boxed{Z_{\text{YM}}(M) \stackrel{?}{=} Z_A^{\text{top}}(T(M^4))}$$