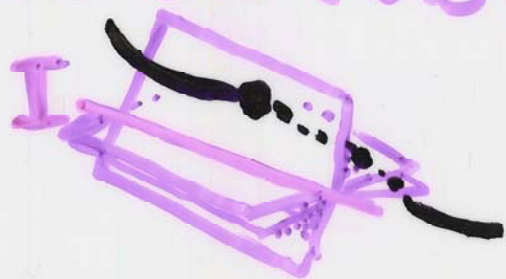


Questions for String Theorists

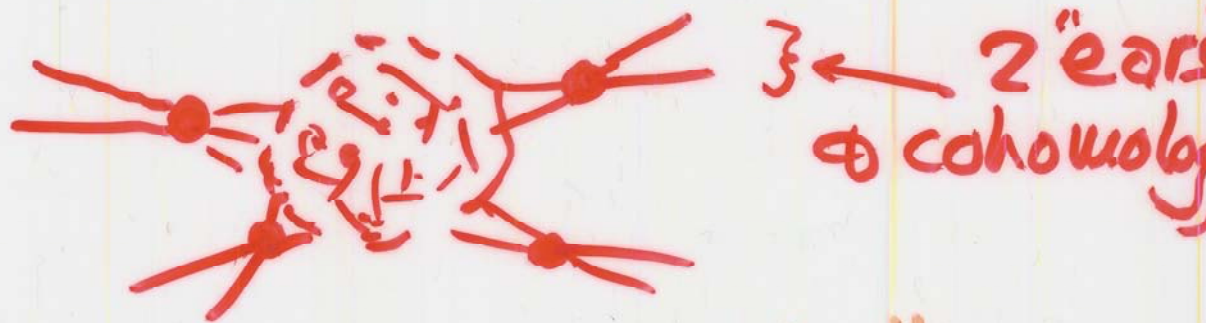
- Do we need the extra dims?
can PT take the place of Calabi-Yau?

- Do we need the supersymm?
Is parametrization invar. needed in twistor theory?



(conformal breaking)

- Do we really want $(++--)$ & a (pseudo-)Wick rotation?
Elemental states?



- No helicity for off-shell
"gluing MHV's"
holomorphic

- Theory?? "Topological" QFT?

Some Remarks on Twistor/Spinor convention

- In $++++$ or $++--$ signature we have a complex conjugation $z^\alpha \mapsto \bar{z}^\alpha$ rather than \bar{z}_α

n -twistor Symmetry Group (Lorentzian case) so does not agree with canonical conjugation $n=1$ so wavefunctions are not holomorphic.

- In $A_{\alpha\beta} = 2\bar{z}^\alpha z^\beta$ consider δ fns on real α -planes. In \mathbb{CM} these would be $\bar{\mu}^{AB} + \mu^{AB}$ $A = \begin{pmatrix} 0 & p_A \\ p^{A'} & 2i\mu^{A'B'} \end{pmatrix}$ $\bar{A} = \begin{pmatrix} -2i\bar{\mu}^{AB} & p_B \\ p_{A'} & 0 \end{pmatrix}$ \therefore cohomology elements in space-time

Compare ASD part with SD part Hedges = (elementary states) General elementary

- Up/down nature of spinor parts $(\omega^A, \pi_{A'})$ fits in with conformal trans group

$$\hat{g}_{ab} = \Omega^2 \begin{pmatrix} g_{ab} \\ z^\alpha \\ \bar{z}^\alpha \end{pmatrix} \quad \hat{E}_{AB} = \frac{1}{\Omega^2} \begin{pmatrix} \epsilon_{AB} \\ z^\alpha \\ \bar{z}^\alpha \end{pmatrix} \quad \hat{\pi}_{A'} = \pi_{A'} + i \tau_{AA'} \omega^A$$

Quantum Gravity

The application of standard procedures of quantum(field) theory to general relativity (or other gravitational theory)?

OR

Some more even-handed marriage, with give on both sides?

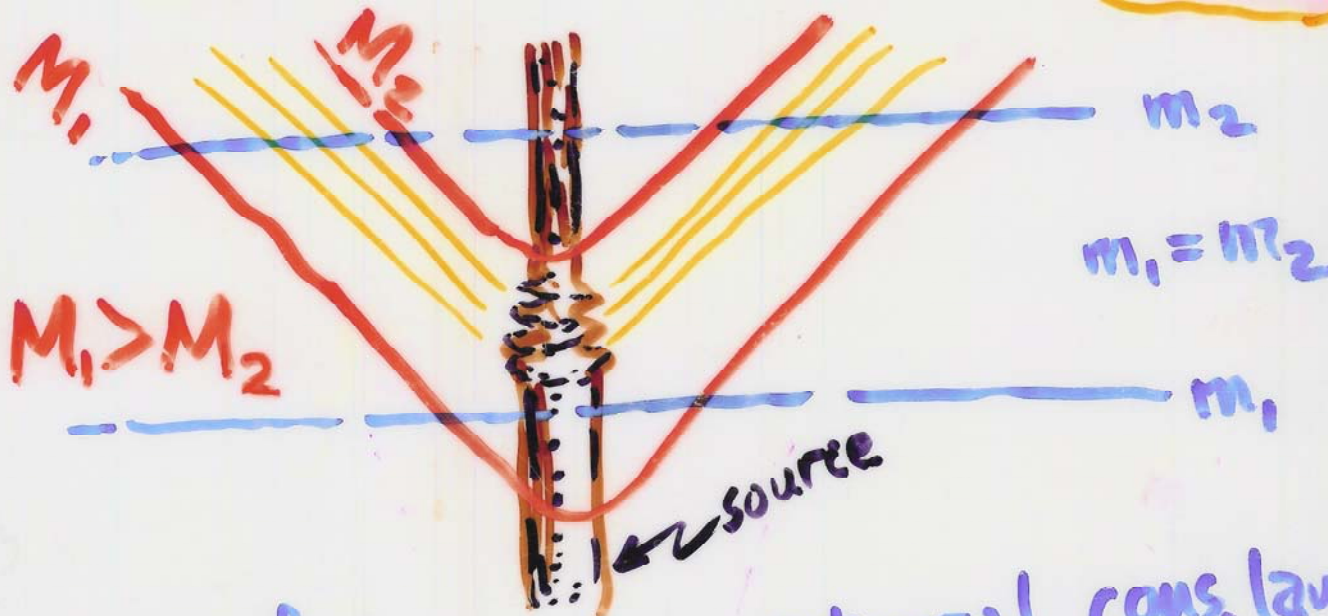
⇒ modified quantum mechanics

Gravitational Radn.

Theory: Bondi 1960
use NULL ref. system

And earlier work by Trautman

time ↑



NB. $\nabla_a T^{ab} = 0$ not an integral cons. law

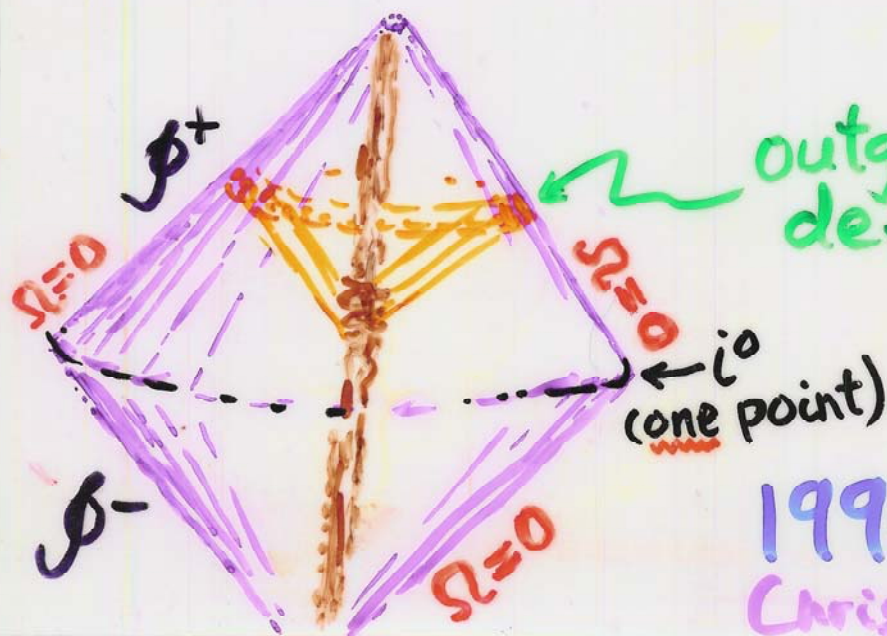
~1960 Trautman

1962 Bondi, van der Burg, Metzner

1962 Sachs

1962 Newman, Unti; Newman, RP

1964/5 RP conformal approach



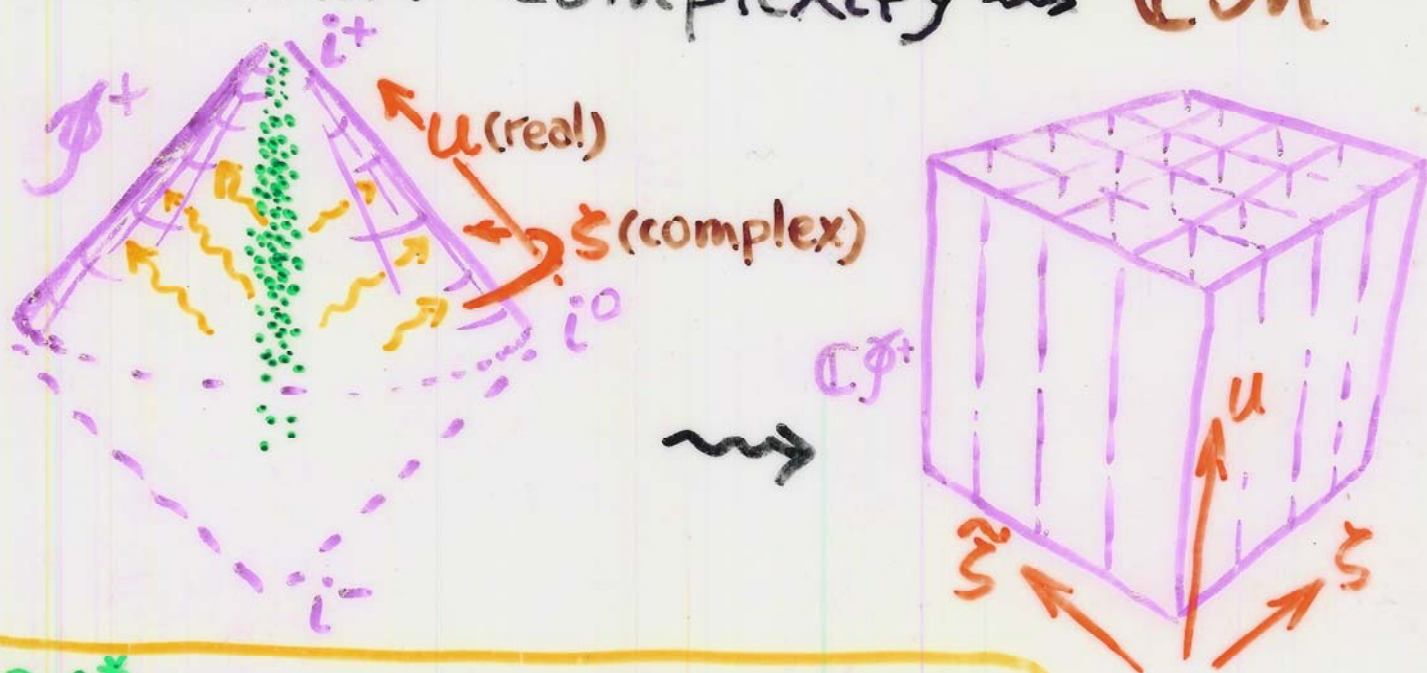
outgoing radiation defined here

$$\hat{g} = \Omega^2 g$$

1990s Friedrich
Christodoulou-Kleinerman

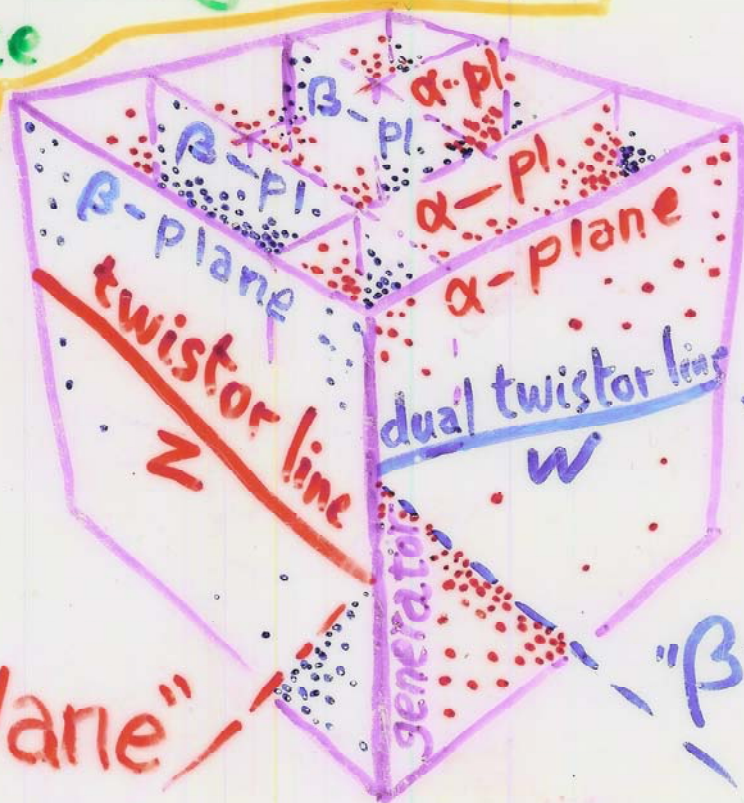
Asymptotic twistors

\mathcal{M} asymptotically flat, radiating, analytic vacuum: complexify $\leadsto \mathbb{C}\mathcal{M}$



\mathcal{H}^* -space = ASD "space-time" whose ("non-linear graviton") twistor space is the space

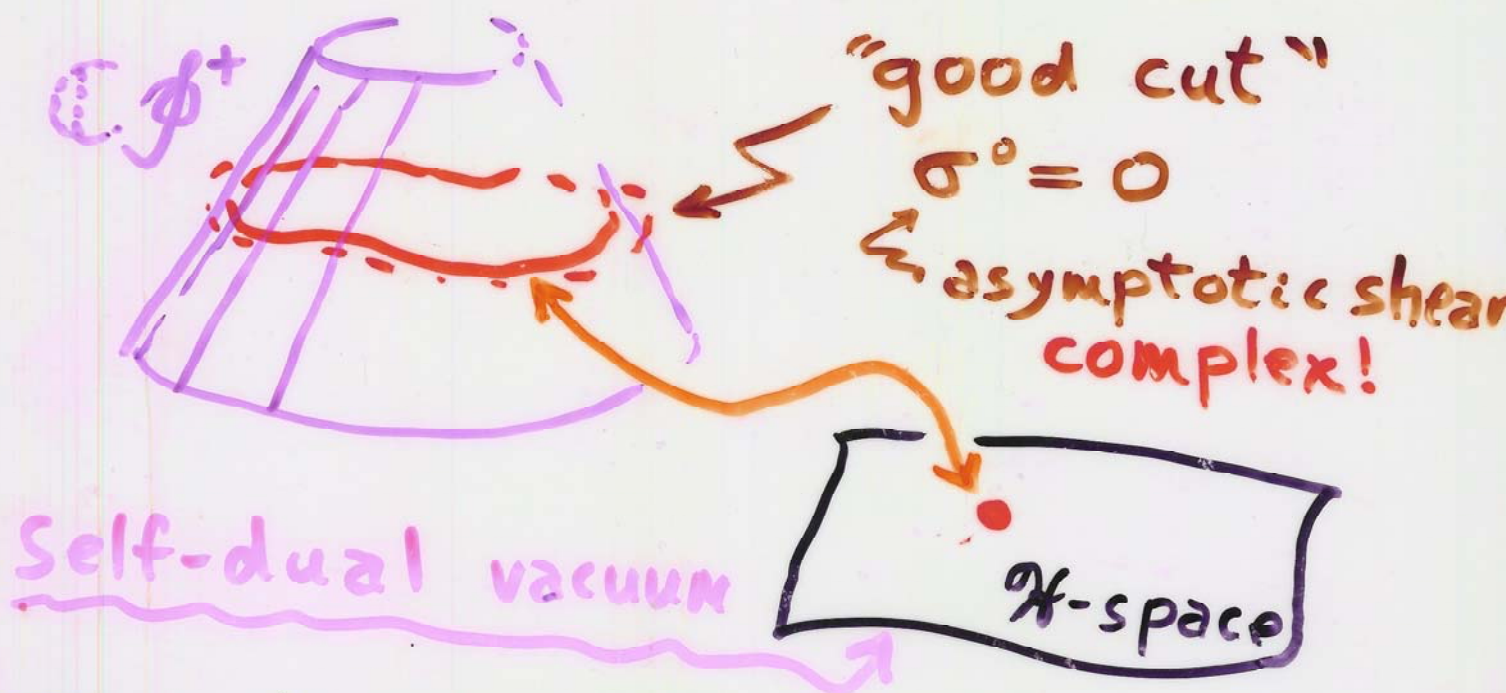
\mathcal{T} of twistor lines \mathbb{Z}



" α -plane"

" β -plane"

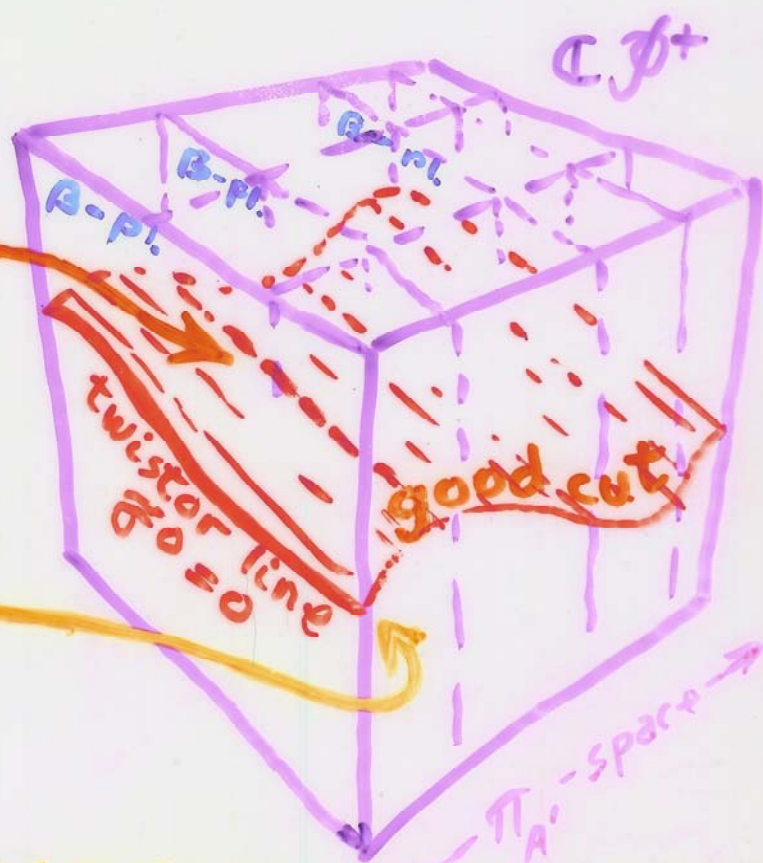
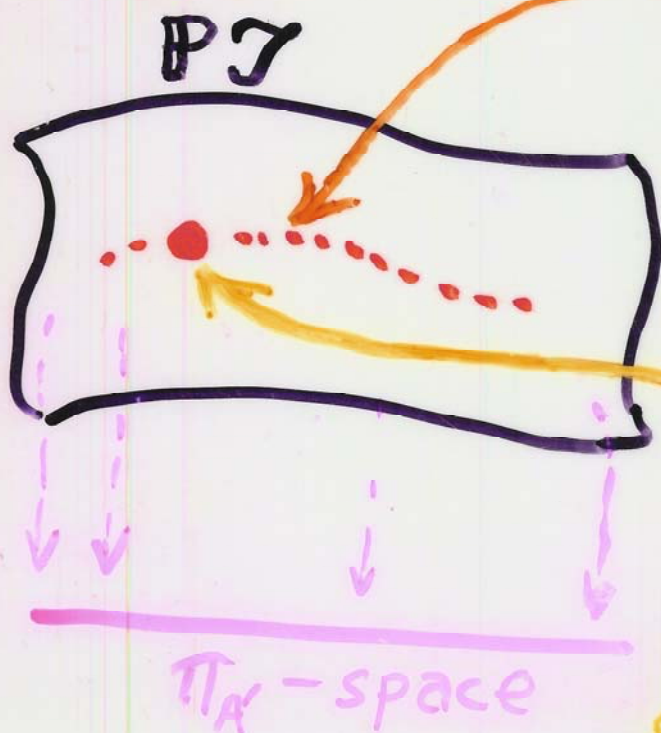
Newman's \mathcal{H} -space



Twistor view (take \mathcal{H}^* -space:)

$\bar{\sigma}^0 = 0$

Holomorphic curve in \mathbb{P}^7 corr. (dual) good cut of $\mathbb{C}\phi^+$



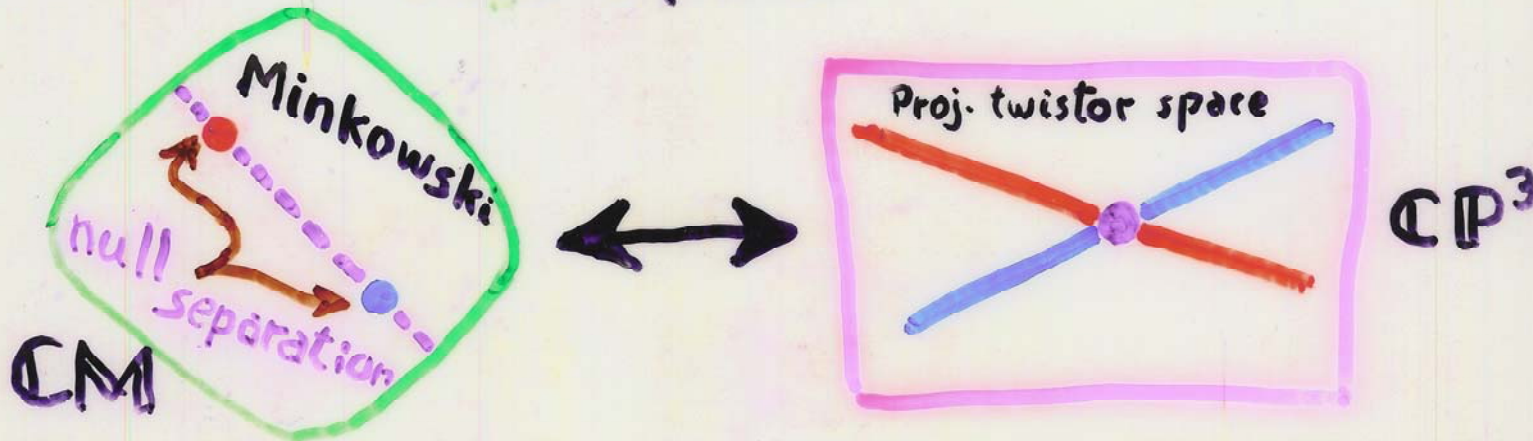
NB. A twistor line (and also a dual twistor line and a generator) is a **NULL GEODESIC** on $\mathbb{C}\phi^+$

General Relativity

Numerous special applications
(e.g. Woodhouse-Mason: stationary axi-symm

As part of general programme:
"non-linear graviton construction" [R.P. '77]

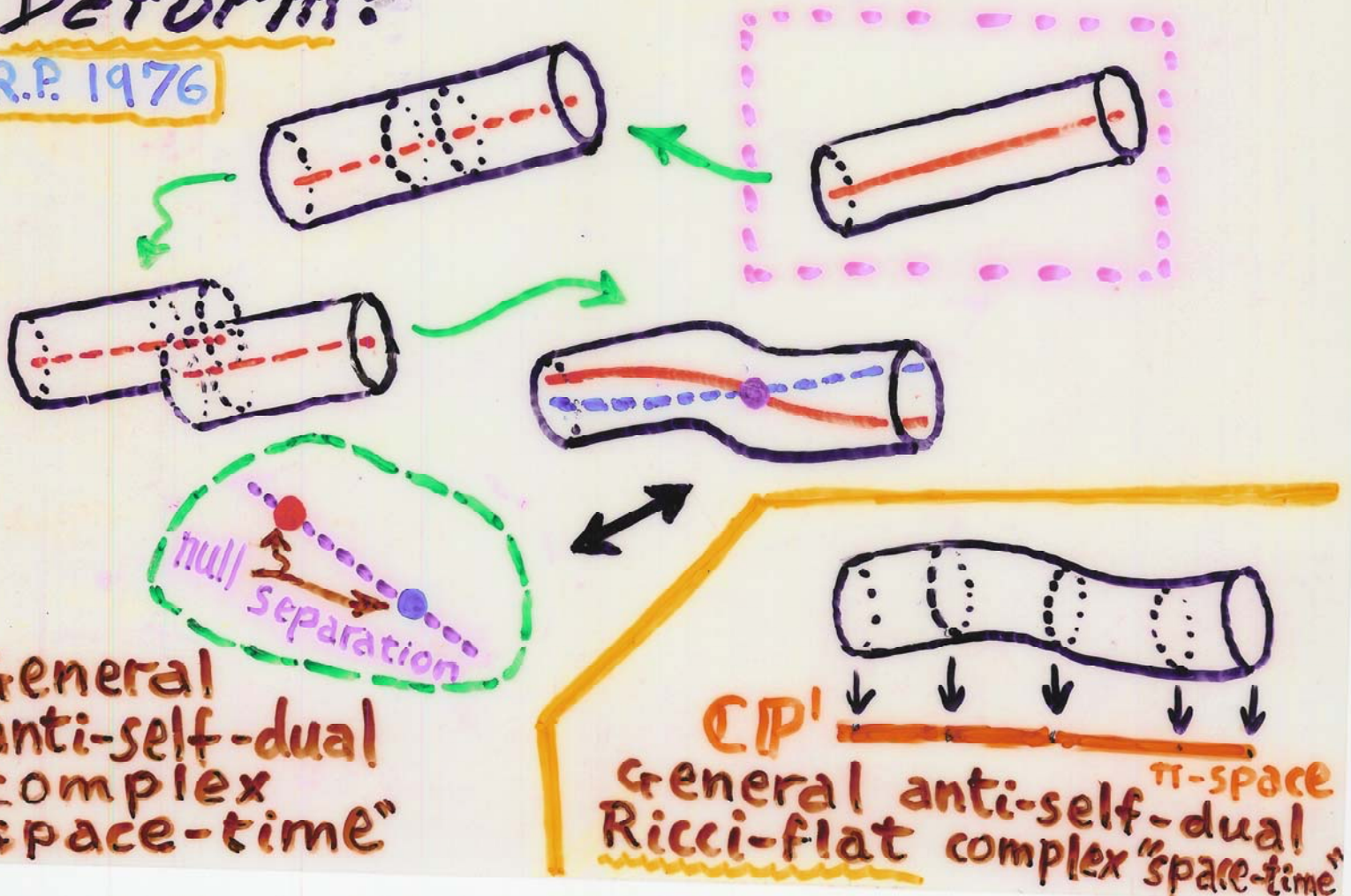
N.B. for flat space:



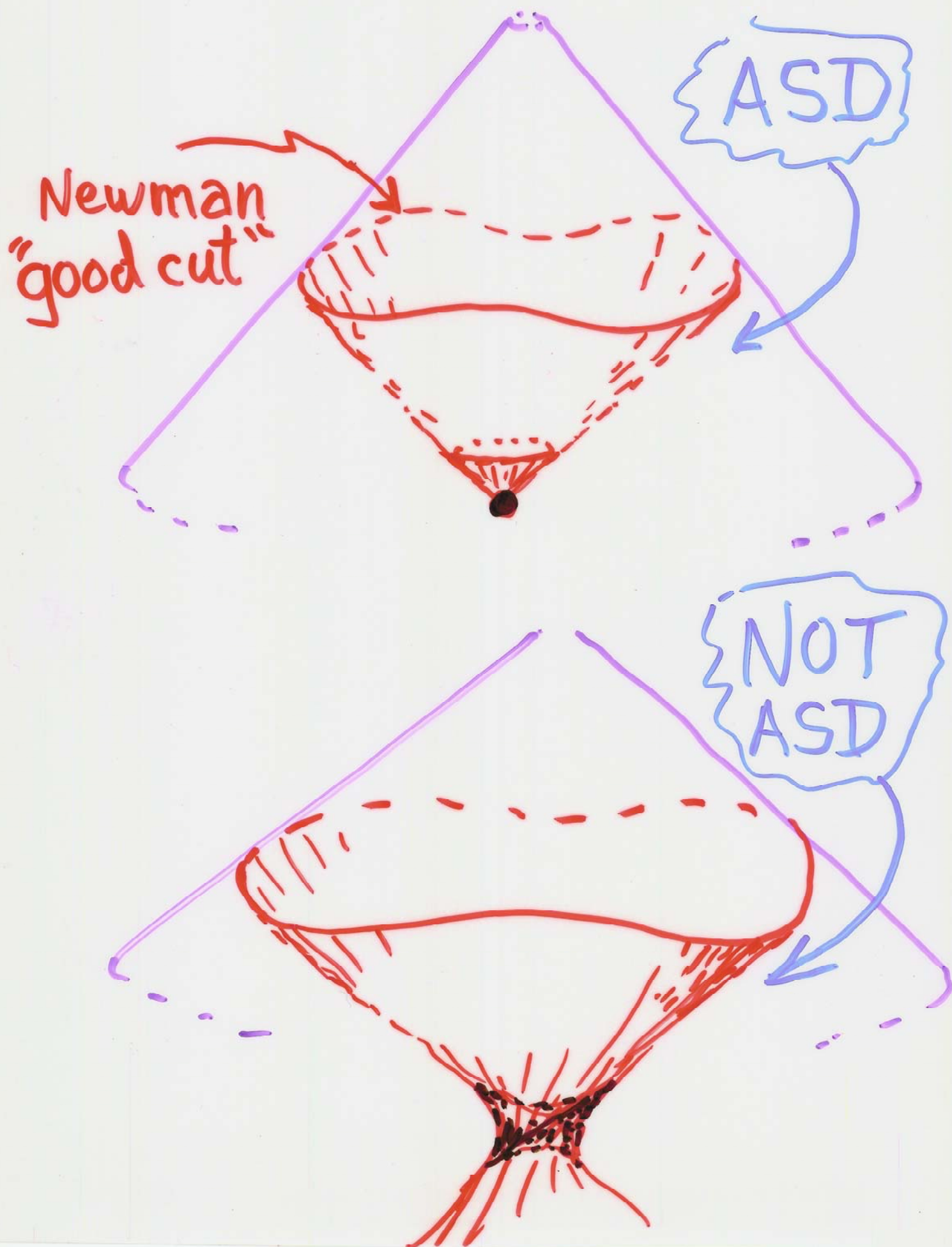
null separation \longleftrightarrow meeting lines

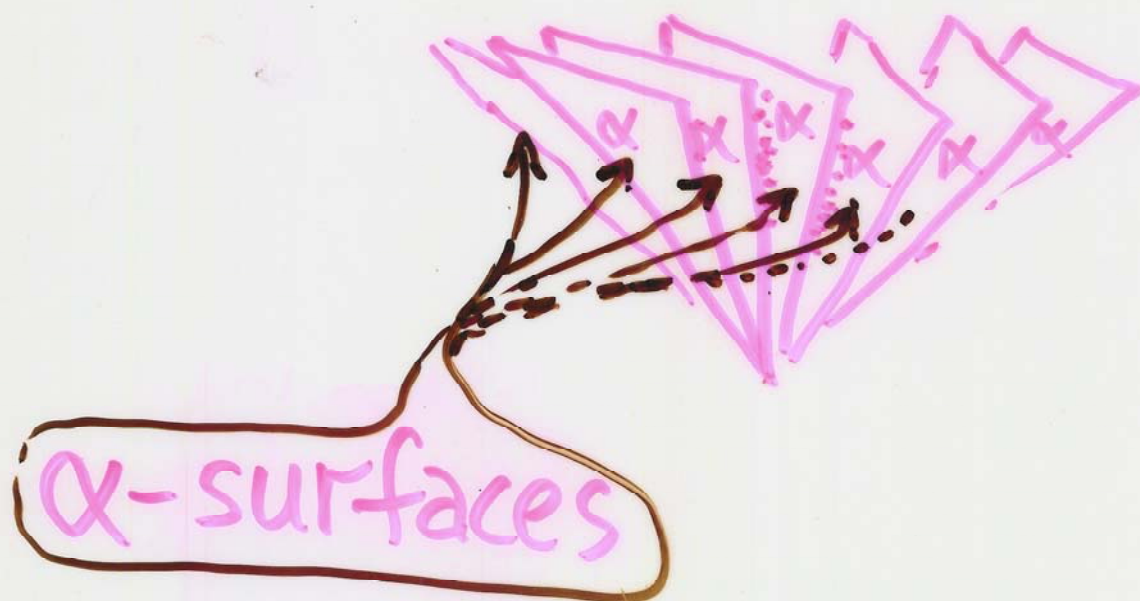
Deform:

R.P. 1976



Light-cone cuts



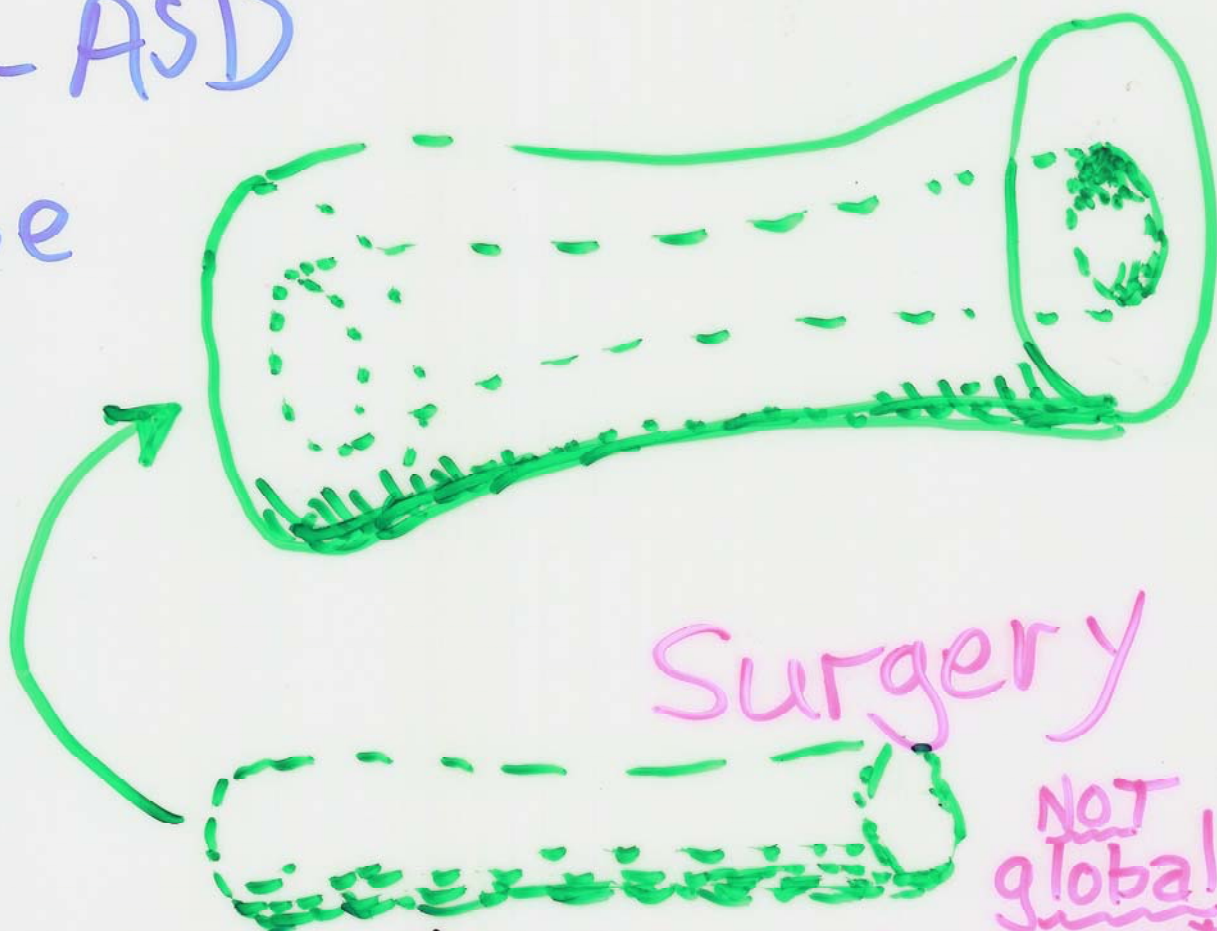


Shearing
comes about
from SD curv.

ASD
case



non-ASD
case



Surgery

NOT
global
in a partic.
way.

Can we understand
these surgeries in terms
of a power series in strings
(holomorphic curves) of higher
and higher order?

Linearized gravity: (can be complex)

$$K_{abcd} = \underbrace{\psi_{ABCD} \epsilon_{A'B'} \epsilon_{C'D'}}_{\text{anti-self-dual}} + \underbrace{\epsilon_{AB} \epsilon_{CD} \tilde{\psi}_{A'B'C'D'}}_{\text{self-dual}}$$

ψ_{ABCD} , $\tilde{\psi}_{A'B'C'D'}$ both symmetric; if K_{abcd} is real, they are complex conjugates

$$\nabla^{AA'} \psi_{ABCD} = 0, \quad \nabla^{AA'} \tilde{\psi}_{A'B'C'D'} = 0$$

Twistor contour integrals:

$$\psi_{ABCD}(x) = \oint_{\omega=i\pi} \frac{\partial}{\partial \omega^A} \frac{\partial}{\partial \omega^B} \frac{\partial}{\partial \omega^C} \frac{\partial}{\partial \omega^D} f(z) \delta z$$

hom. deg. -6 \leftarrow hom. deg. $+2$

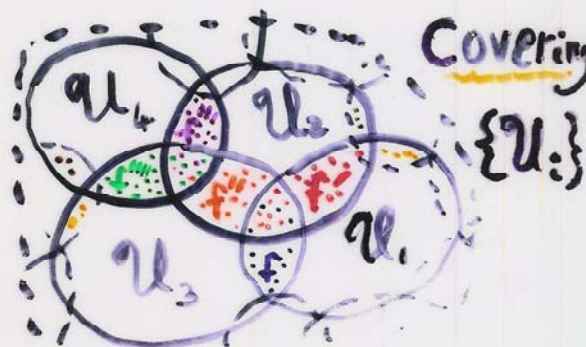
$$\tilde{\psi}_{A'B'C'D'}(x) = \oint_{\omega=i\pi} \pi_{A'} \pi_{B'} \pi_{C'} \pi_{D'} \tilde{f}(z) \delta z$$

$$\delta z = \epsilon^{A'B'} \pi_{A'} d\pi_{B'}$$

Left-handed graviton: f hom. deg. $+2$
 Right-handed graviton: \tilde{f} hom. deg. -6

Really, f and \tilde{f} are representatives of

COHOMOLOGY



Encoding $\mathcal{O}(-6)$

$$Z^{\alpha}_i = (1 + f_{(-6)}{}^{ij})^{1/6} Z^{\alpha}_j$$



$$\mu^{B'} \nabla_{B'} \mu_{A'} = K \mu_{A'} (\mu_i)^{-5} E_6 \tilde{\psi}_{i' i' i' i'}$$

"i" means
contract with $\tilde{l}^{A'}$
or l^A

↗ conf. invar.
version of $\nabla_{H'}$

Twistor Space (New Def.)

Complex 4-manifold \mathcal{T}

Global structure encodes asymp. flat vacuums

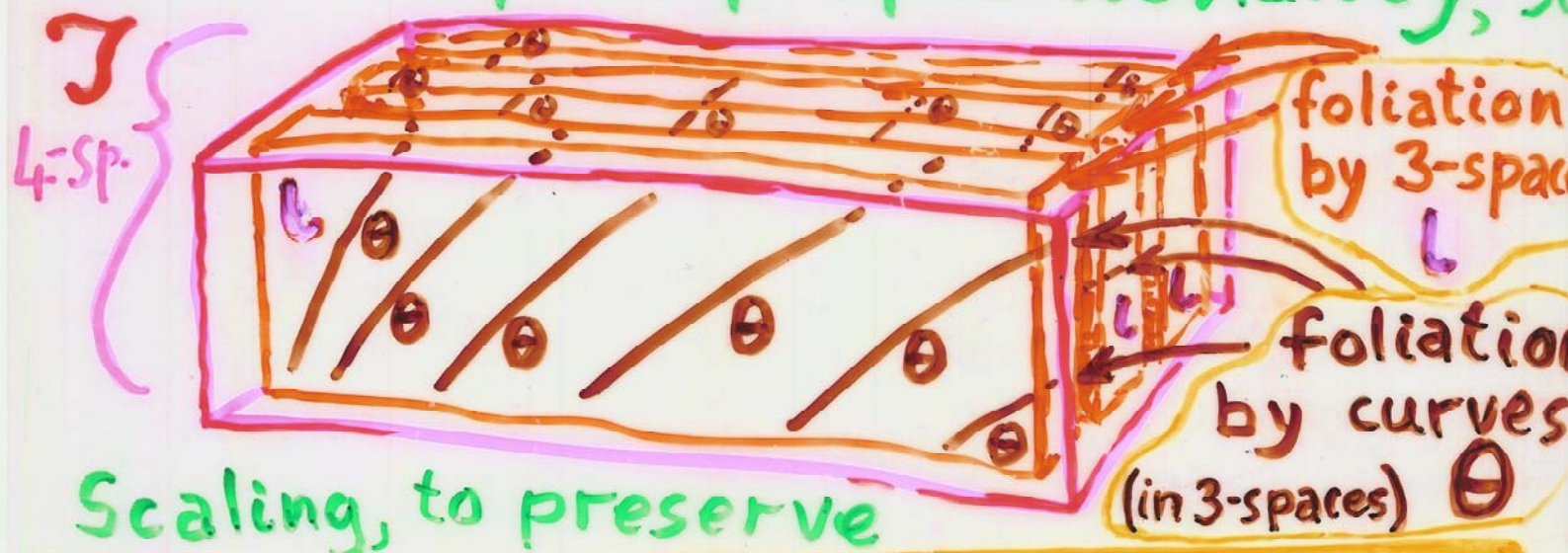
Local structure (no continuous information):

1-form L , 3-form Θ ,
satisfying:

$$L \wedge dL = 0,$$

$$L \wedge \Theta = 0,$$

known up to proportionality, so



$$\Pi = d\Theta \otimes d\Theta \otimes \Theta, \quad \Sigma = d\Theta \otimes L = -2\Theta \otimes dL$$

(where $\eta \otimes (dp \wedge dq) = \eta_\lambda dp \otimes dq - \eta_\lambda dq \otimes dp$)

Twistor-space forms M

1-form: $L (= \delta z)$

$$= I_{\alpha\beta} Z^\alpha dZ^\beta = \varepsilon^{A'B'} \pi_{A'} d\pi_{B'}$$

2-form: $\tau = \frac{1}{2} dL$

$$= \frac{1}{2} I_{\alpha\beta} dZ^\alpha \wedge dZ^\beta = d\pi_{0'} \wedge d\pi_{1'}$$

3-form: θ

$$\begin{aligned} &= \frac{1}{6} \varepsilon_{\alpha\beta\gamma\delta} Z^\alpha dZ^\beta \wedge dZ^\gamma \wedge dZ^\delta \\ &= Z^0 dZ^1 \wedge dZ^2 \wedge dZ^3 \\ &\quad - Z^1 dZ^0 \wedge dZ^2 \wedge dZ^3 \\ &\quad + Z^2 dZ^0 \wedge dZ^1 \wedge dZ^3 \\ &\quad - Z^3 dZ^0 \wedge dZ^1 \wedge dZ^2 \end{aligned}$$

4-form: $\phi = \frac{1}{4} d\theta$

$$\begin{aligned} &= \frac{1}{24} \varepsilon_{\alpha\beta\gamma\delta} dZ^\alpha \wedge dZ^\beta \wedge dZ^\gamma \wedge dZ^\delta \\ &= dZ^0 \wedge dZ^1 \wedge dZ^2 \wedge dZ^3 \end{aligned}$$

Euler: $\gamma = \theta \div \phi = Z^\alpha \partial / \partial Z^\alpha$

$da_\wedge \theta = \gamma(a) \phi$

$L \wedge \tau = 0, L \wedge \theta = 0$

Homogeneity degrees

$$\mathbb{L} \begin{bmatrix} L \\ \tau \\ \theta \\ \phi \end{bmatrix} = \begin{bmatrix} 2L \\ 2\tau \\ 4\theta \\ 4\phi \end{bmatrix}$$

$\} \Leftrightarrow \theta \otimes \tau = -\phi \otimes L$
 $\} \text{ automatic } (\theta \otimes \phi \equiv -\phi \otimes \theta)$

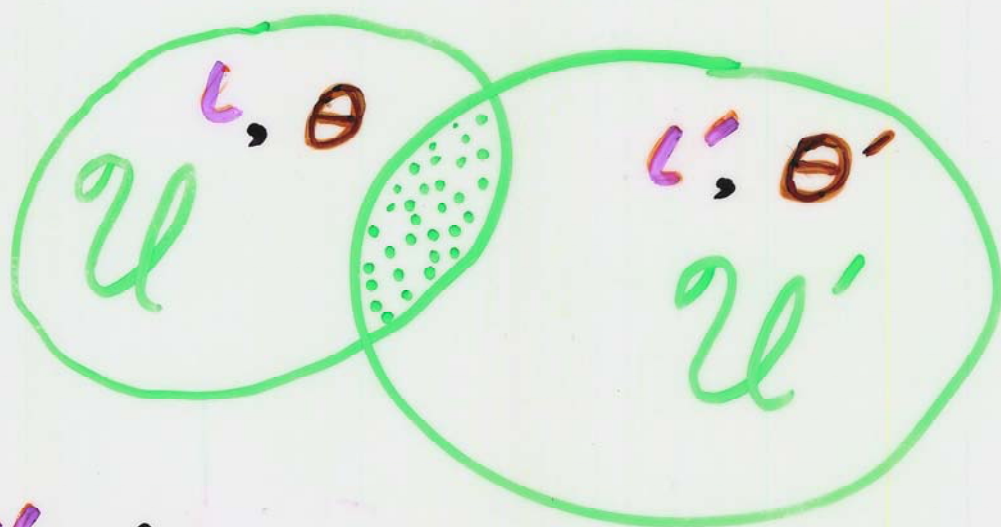
where

$\alpha \otimes \beta = r \alpha [\dots \beta] \dots$
← r-form

e.g.

$\alpha \otimes dp \wedge dq = \alpha \wedge dp \otimes dq - \alpha \wedge dq \otimes dp$

Googly scalings



$$\left. \begin{aligned} L' &= kL \\ \theta' &= k^2 \theta \\ d\theta' &= k^{-1} d\theta \end{aligned} \right\}$$

preserves Π and Σ

$$\text{and } d\theta \otimes L = -2\theta \otimes dL$$

gives

$$r(k) = 2k^{-2} - 2k$$

equiv.: $r'(k^{-1}) = 2k^2 - 2k^{-1}$ since $r' = k^3 r$
 recall: $r = \theta \div (\frac{1}{4} d\theta)$

Standard param. up Euler curves

$$r(z) = z$$

Find: $k^3 = 1 - F z^{-6}$ where F const. on curve

$$\therefore k^3 = 1 - f_{-6}(z^a)$$

References (for googly construction) [First three not very detailed]

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