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A Simple Thousand-Year Prognosis for Oceanic and Atmospheric Carbon Change

A. C. FOWLER^{1,2}

Abstract—A simple carbon-cycle box model allows for very simple quantitative insight into the evolution of climate over the next millennium. While melting ice sheets, rising sea levels, and ocean acidification are well recognised, we show that a further consequence of the repartitioning of ocean carbon is a dramatic rise in atmospheric carbon dioxide on a millennial time scale.

Key words: Climate change, carbon cycle, ocean acidification.

1. Introduction

Climate change is prevalent in the news, and there is a growing consensus on some of the more pertinent facts and hypotheses. It is a simple consequence of radiative transfer theory that increasing CO₂ levels in the atmosphere cause atmospheric warming (PIERRE-HUMBERT 2010), with an approximate logarithmic dependence (HOUGHTON 2009). It is well-known that CO₂ levels oscillated between 180 and 280 ppmv during the Pleistocene ice ages (PETIT et al. 1999) but have risen since 1750 to a level near 400 ppmv. Here, we will preferentially use partial pressures for atmospheric CO₂ levels. Dalton's law of partial pressures states that, for a dilute perfect gas, the ratio of partial CO₂ pressure to air pressure is approximately the molar volume fraction. For an atmospheric pressure of 10⁵ Pa and a partial CO₂ pressure of p_{CO_2} , we get a volume fraction $p_{CO_2}/10^5 = 10 p_{CO_2}$ ppmv; thus, roughly, 400 ppmv = 40 Pa.

The Earth has been cooling since the Eocene, some 50 My (million years) ago (BuL *et al.* 2013). This cooling is generally associated with declining levels of atmospheric carbon (PEARSON and PALMER 2000) that are likely due to the increased weathering associated with Himalayan uplift following the collision of India with Asia, as well as the associated subduction of pelagic carbonates as the Tethys Ocean closed (CALDEIRA 1992).

The Antarctic continent was sub-tropical on its coasts (PROSS et al. 2012), and the Antarctic Ice Sheet grew slowly thereafter, possibly originating as three separate ice caps (POLLARD and DECONTO 2005) before these coalesced at the Eocene-Oligocene transition at 34 My (ZACHOS et al. 1992). In model experiments, the hysteretic decay of the Antarctic Ice Sheet (POLLARD and DECONTO 2005) is associated with the increase of atmospheric CO₂ levels above about 900 ppmv (90 Pa), but one would expect melting to be initiated at levels well below that in conditions of rapid warming. And indeed, the great ice sheets do appear to be melting: outlet glaciers on Greenland and Antarctica have all shown dramatic signs of increased discharge and collapse in recent decades (HOLLAND et al. 2008; ROTT et al. 1996; RIGNOT et al. 2002).

Apart from the atmospheric warming caused by increased atmospheric CO₂ levels, a secondary effect is that of ocean acidification (ZEEBE 2012a). The increased CO₂ in the atmosphere causes an increase in dissolved CO₂ in the ocean, and the effect of the bicarbonate buffering system is to cause a decrease in carbonate ion, and hence a decrease in pH (ARCHER *et al.* 1997; ARCHER 1999). Emphasis has been placed on the possible consequences for calcifiers in the ocean (RIDGEWELL and ZEEBE 2005; ZEEBE *et al.* 2008), with the suggestion that decreased calcium carbonate formation leads to lower oceanic burial rates, and thus, further enhancement of ocean carbon stocks. In addition, simulation models (ZEEBE 2012b) indicate an associated increase in temperature of the

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order of ten degrees on a millennial time scale (ZEEBE and ZACHOS 2013).

Much of the discussion concerning climate prediction focusses on details of recent warming over the last century, and restricts itself to a range of forecasts over the coming century based on ever more detailed numerical weather prediction models, as described in the 2007 Intergovernmental Panel on Climate Change (IPCC) report (SOLOMON *et al.* 2007), for example. Together with the increasing complexity of these models comes an inevitable uncertainty, due both to chaotic dynamics and model uncertainty (PALMER 2001), with much of the latter being attributable to uncertainty in cloud parameterisation (TIEDTKE 1993).

Arguably, this 'kitchen sink' approach (include every process possible) misses the elephant in the room, which becomes visible when one adopts a 'topdown' approach (include only the essentials of the problem); indeed, this is partly the view of Archer, and of Zeebe and Zachos, who put an emphasis on ocean acidification beyond the century time scale.

There have been a number of recent studies which emphasise the long-term (millennial) evolution of atmospheric and oceanic carbon (GILLETT *et al.* 2011; SOLOMON *et al.* 2009, 2010; ARCHER and BROVKIN 2008; EBY *et al.* 2009). The studies use models of intermediate complexity, suitable for such longer time-scale studies, but still necessary to be solved numerically, so that the mechanisms of the processes are somewhat opaque.

The purpose of the present paper is to illustrate the way in which atmospheric carbon evolves over long time scales using the simplest model that still represents the essential exchange processes in the system, and particularly, to point out a third consequence of carbon injection (beyond warming and acidification). The criticism of over-simplicity must be balanced against the clarity of the insight that emerges.

2. A Simple Model

Our simplest model for climate prediction is based on a box model introduced by FowLER *et al.* (2013). Conceptually, the model is a generalised

energy balance model of the type introduced by BUDYKO (1969) and SELLERS (1969), and whose progeny are the models of intermediate complexity (GANOPOLSKI *et al.* 2010), which are applied with some success to climate models, ocean models (JOHNSON *et al.* 2007; TOGGWEILER 2008), and biogeochemical models (ZEEBE 2012b). Our approach here is to choose the very simplest such model that still encapsulates the essential physics of climate control. Because the energy balance part of the model simply reflects an essentially algebraic dependence of global average temperature on atmospheric CO₂, our focus is on the evolution of carbon within the ocean and atmosphere.

The basic ingredient of the model of FOWLER *et al.* (2013) is the dependence of mean atmospheric temperature on CO_2 content. To describe the atmospheric CO_2 , we must balance anthropogenic or volcanic sources with fluxes to the ocean via dissolution or indirectly through weathering of silicate and carbonate rocks. At the ocean surface, the CO_2 concentration is related to the dissolved CO_2 in the water by Henry's law:

$$p_{\rm CO_2}^s = \frac{[{\rm CO_2}]}{K_{\rm H}};$$
 (2.1)

 $K_{\rm H}$ is Henry's constant, and has units M Pa⁻¹, where 1 M = 1 mol kg⁻¹. Similarly, we may take the dissolved CO₂ in raindrops to be $K_{\rm H}p_{\rm CO_2}$, where $p_{\rm CO_2}$ is atmospheric CO₂ partial pressure. If *P* is the global rate of precipitation, then this leads to a net CO₂ precipitation of $K_{\rm H}\rho_{\rm H_2O}Pp_{\rm CO_2}$, with units of mol y⁻¹. This leads us to a form of FowLER *et al.*'s (2013) equation (3.21):

$$\frac{A_{\rm E}}{M_a g} \dot{p}_{\rm CO_2} = I - K_{\rm H} \rho_{\rm H_2O} P p_{\rm CO_2} - h_{oc} (p_{\rm CO_2} - p_{\rm CO_2}^s),$$
(2.2)

which describes conservation of CO_2 in the atmosphere; the units are mol y^{-1} . The constants are as described in Table 1.

In order to determine p_{CO_2} , we need to know the dissolved CO_2 concentration in the ocean, which itself depends on the carbon balance. Particularly, carbon is buffered in the ocean by means of the following reactions:

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Table 1

Values and description of the constants in the model. The preindustrial CO_2 production rate is based on an estimated volcanic production rate of 3×10^{11} kgCO₂ y⁻¹ (WALKER et al. 1981), while the anthropogenic value corresponds to a fossil fuel and cement emissions rate of 6 GtC y⁻¹, a value typical of the 1990s (DENMAN et al. 2007; indicating production rates of 5.4, 6.4 and 7.2 GtC y⁻¹ in the 1980s, 1990s, and early 2000s, respectively)

Symbol	Meaning	Typical value	Comment
A _E	Earth surface area	$5.1 \times 10^{14} m^2$	
A_L	Land surface area	$1.5 \times 10^{14} \text{ m}^2$	
g	Acceleration due to gravity	9.81 m s ⁻²	
h _{oc}	Interfacial transport coefficient	$1.8 \times 10^{13} \text{ mol}$ $Pa^{-1} \text{ y}^{-1}$	
Ι	CO ₂ production rate	$0.7 \times 10^{13} \text{ mol y}^{-1}$	Pre-industrial
		$5 \times 10^{14} \text{ mol y}^{-1}$	Anthropogenic
$K_{\rm H}$	Henry's law coefficient	$4.5 \times 10^{-7} \text{ M Pa}^{-1}$	
M_a	Molecular weight of air	$\frac{2.88 \times 10^{-2}}{\text{kg mol}^{-1}}$	
moc	Ocean mass	$1.38 \times 10^{21} \text{ kg}$	
Р	Global precipitation	$5 \times 10^{14} m^3 y^{-1}$	
w	Net carbonate loss	$0.5 \times 10^{13} \text{ mol y}^{-1}$	
$ ho H_2O$	Density of sea water	1.025×10^{3} kg m ⁻³	

$$\begin{aligned} \mathrm{H}_{2}\mathrm{O}) + \mathrm{CO}_{2} & \underset{k_{-1}}{\overset{k_{1}}{\underset{k_{-1}}{\longrightarrow}}} \mathrm{HCO}_{3}^{-} + \mathrm{H}^{+}, \\ \mathrm{HCO}_{3}^{-} & \underset{k_{-2}}{\overset{k_{2}}{\underset{k_{-2}}{\longrightarrow}}} \mathrm{CO}_{3}^{2-} + \mathrm{H}^{+}. \end{aligned}$$

$$(2.3)$$

These reactions are very fast (ZEEBE and WOLF-GLADROW 2001), and thus can be taken to be in equilibrium, whence we have

$$X = \frac{K_2 Q}{S}, \quad Y = \frac{K_2 Q^2}{K_1 S},$$
 (2.4)

where we write

(

$$X = [\mathrm{H}^+], \quad Y = [\mathrm{CO}_2], \quad Q = [\mathrm{HCO}_3^-], \quad S = [\mathrm{CO}_3^{2-}],$$
(2.5)

and the equilibrium constants K_i are defined by

$$K_1 = \frac{k_1}{k_{-1}}, \quad K_2 = \frac{k_2}{k_{-2}}.$$
 (2.6)

The total dissolved inorganic carbon concentration in the ocean is

$$C = Q + S + Y, \tag{2.7}$$

and a conservation law for total carbon is

$$m_{oc}\dot{C} = h_{oc}(p_{\rm CO_2} - p_{\rm CO_2}^s) + \left(\frac{A_{\rm E} - A_{\rm L}}{A_{\rm E}}\right) K_{\rm H} \rho_{\rm H_2O} P p_{\rm CO_2} - r_c - r_p + W,$$
(2.8)

in which m_{oc} is ocean mass, A_L is land surface area, r_c is the removal rate of bicarbonate by calcifiers, r_p is the removal rate of carbonate by calcium carbonate precipitation, and W represents the supply of bicarbonate to the ocean through weathering and riverine discharge. FowLER *et al.* (2013) provide further individual equations for the separate components, but in the present situation a more succinct discussion is appropriate.

The source and sink terms in (2.8) are relatively small; they cause adjustments to oceanic concentrations on time scales upwards of forty thousand years, which represents the normal adjustment time for carbonate. Because these time scales are much longer than the ocean mixing time, the box model makes sense on such long time scales. On the much shorter time scale associated with anthropogenic input to the atmosphere and thence, the ocean, the dissolved CO_2 adds no charge to the charge-neutral ocean. The concentration of the conservative ions chloride, sodium, etc., do not change at all, and the calcium also will not change, since its time scale for adjustment is so long, even if the calcifier populations were wiped out. Therefore, the net-negative charge Q + 2S - X of the active ions H^+ , HCO_3^- , and CO_3^{2-} must remain constant under conditions of rapid change. Since from Table 2 we see that $X \ll Q, S$, this implies

$$Q + 2S \approx \text{constant}$$
 (2.9)

over relatively short (millennial) time scales. Since also $Y \ll S$, we see that (2.8) takes the form

$$-m_{oc}\dot{S} = h_{oc}(p_{\rm CO_2} - p_{\rm CO_2}^s) + \left(\frac{A_{\rm E} - A_{\rm L}}{A_{\rm E}}\right) K_{\rm H}\rho_{\rm H_2O} P p_{\rm CO_2}$$
$$-r_c - r_p + W.$$
(2.10)

This completes our millennial-scale model. For such short time scales, we will eventually ignore the small A. C. Fowler

Table 2

Present estimates of the equilibrium constants K_i and the ionic species concentrations in units of M (1 M = 1 mol kg⁻¹) (EMERSON and HEDGES 2008). C is dissolved inorganic carbon

Species/constant	Typical value (M)
[Ca ²⁺]	1.03×10^{-2}
$[CO_2] = Y$	0.8×10^{-5}
$[CO_3^{2-}] = S$	0.24×10^{-3}
$[\mathrm{H}^+] = X$	0.63×10^{-8}
$[\text{HCO}_3^-] = Q$	1.7×10^{-3}
<i>K</i> ₁	1.4×10^{-6}
<i>K</i> ₂	1.1×10^{-9}
С	2.0×10^{-3}

source and sink terms, but we retain them initially so that we can sensibly discuss the pre-industrial situation.

3. The Past and the Future

We take our model in the form (2.2) and (2.10), as follows:

$$\begin{aligned} \frac{A_{\rm E}}{M_{ag}} \dot{p}_{\rm CO_2} &= I - K_{\rm H} \rho_{\rm H_2O} P p_{\rm CO_2} - h_{oc} (p_{\rm CO_2} - p^s_{\rm CO_2}), \\ m_{oc} \dot{S} &= -h_{oc} (p_{\rm CO_2} - p^s_{\rm CO_2}) \\ &- \left(\frac{A_{\rm E} - A_{\rm L}}{A_{\rm E}}\right) K_{\rm H} \rho_{\rm H_2O} P p_{\rm CO_2} + w, \end{aligned}$$
(3.1)

where w represents the net loss of dissolved inorganic carbon (DIC) due to precipitation and calcification minus the supply from weathering. Our estimate of its value in Table 1 is based on a river efflux to the ocean of $3.5 \times 10^{13} \text{ m}^3 \text{ y}^{-1}$ and a river bicarbonate concentration of 0.96 mM (EMERSON and HEDGES 2008). This gives a net molar flux of carbon of 3.4×10^{13} mol y⁻¹, and we suppose the loss in the ocean is comparable to this. The value we choose for w is determined by assuming the pre-industrial carbon reservoir in the ocean is approximately in a state of balance between input and output (the balance will not be exact due to transient effects associated with recovery from the last ice age, but is not expected to be wildly out of balance, because the variation of atmospheric CO₂ during the ice ages is not enormous). In addition, we have Henry's law and charge conservation from (2.1), (2.4), and (2.9):

$$p_s = \frac{K_2 Q^2}{K_1 K_H S}, \quad Q + 2S = A,$$
 (3.2)

where A is, approximately, the carbonate alkalinity.

It is helpful, indeed essential, to non-dimensionalise the equations. We do this by writing

$$p_{\text{CO}_2} = p_0 p, \quad p_{\text{CO}_2}^s = p_0 p_s, \quad Q = Aq, \quad S = S_0 s,$$

 $t \sim t_0,$ (3.3)

where

$$t_0 = \frac{A_{\rm E}}{M_a g h_{oc}}, \quad p_0 = \frac{K_2 A^2}{K_1 K_H S_0},$$
 (3.4)

and

$$A = 2.2 \text{ mM}, \quad S_0 = 0.24 \text{ mM}$$
 (3.5)

are the current values in Table 2. This leads us to the dimensionless model

$$\dot{p} = \varepsilon(v - p) - (p - p_s),$$

$$\dot{s} = \delta \left[-\frac{1}{\varepsilon} (p - p_s) - \{(1 - \alpha)p - \Omega\} \right],$$
 (3.6)

$$p_s = \frac{q^2}{s}, \quad q = 1 - vs,$$

where

$$\varepsilon = \frac{w_0}{h_{oc}p_0}, \quad v = \frac{I}{w_0}, \quad \delta = \frac{w_0 t_0}{m_{oc} S_0}, \quad \alpha = \frac{A_{\rm L}}{A_{\rm E}},$$
$$\Omega = \frac{w}{w_0}, \quad v = \frac{2S_0}{A}, \tag{3.7}$$

and where we also define

$$w_0 = K_{\rm H} \rho_{\rm H_{2}O} P p_0.$$
 (3.8)

Based on Table 1, typical values of the parameters are given in Table 3.

We are expecting pre-industrial steady values for p and S to be of O(1). It is easy to ascertain steady states of these equations. These are

$$p = \frac{v - \Omega}{\alpha}, \tag{3.9}$$

which with the values in Table 3, give p = 0.86, corresponding to 30.3 Pa (but this motivated the choice of w), and then from $p_s \approx p$, we have $s \approx 0.79$, corresponding to 0.19 mM. Obviously, the

Table 3

Values	of	scales	and	parameters
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Symbol	Meaning	Typical value
Α	Carbonate alkalinity	2.2 mM
p_0	CO ₂ scale	35.2 Pa
S_0	Carbonate scale	0.24 mM
t_0	Time scale	100 y
v		0.87 (pre-industrial)
		60.3 (anthropogenic)
w_0	Weathering scale	$0.81 \times 10^{13} \text{ mol y}^{-1}$
α	e	0.29
δ		2.48×10^{-3}
3		1.28×10^{-2}
v		0.22
Ω		0.62

values do not exactly match, but they indicate selfconsistency in the parameter choice.

3.1. Pre-Industrial Steady State

More interesting is the transient approach to this steady state, which is easy to read, based on the small values of the parameters ε and δ . On a time scale of $t_0 \sim 100$ y (dimensionlessly $t \sim O(1)$), p approaches an equilibrium in which $p \approx p_s$, but more accurately

$$p - p_s = \varepsilon(v - p) \approx \varepsilon(v - p_s).$$
 (3.10)

For longer time scales we substitute this into $(3.6)_2$ to obtain (also using $p \approx p_s$)

$$\dot{s} \approx \delta \left[\Omega - \nu + \frac{\alpha (1 - \nu s)^2}{s} \right],$$
 (3.11)

from which we can see that the carbonate ion approaches equilibrium on the longer time scale $\frac{t_0}{\delta} = \frac{m_{oc}S_0}{w_0} \sim 40,000 \text{ y (dimensionlessly } t \sim \frac{1}{\delta}).$

3.2. Anthropogenic Future

Now, consider the nature of the solutions if we use the anthropogenic value of v = 60.3 (corresponding as in Table 1 to a 1990s release rate of 6 GtC y⁻¹). We take this initially as a constant, though in practice the value is steadily rising and is unlikely even to level off in the near future. We define

$$v_A = \varepsilon v \approx 0.77, \qquad (3.12)$$

and note that it is of O(1). The Eq. (3.6) takes the form

$$\dot{p} = v_A - \varepsilon p - (p - p_s),$$

$$\dot{s} = \delta \left[-\frac{1}{\varepsilon} (p - p_s) - \{ (1 - \alpha)p - \Omega \} \right],$$
 (3.13)

and again, it is simple to trace their solutions. On the same atmospheric time scale of 100 y, p reaches an approximate equilibrium

$$p \approx v_A + p_s. \tag{3.14}$$

The meaning of this is that the carbon we are pouring into the atmosphere reaches a quasi-steady state in which the influx is balanced by the efflux to the ocean. With the values in Table 3, this quasi-steady state is (taking p_s as a pre-industrial value 0.8, corresponding to 28 Pa) $p \approx 1.57$, corresponding to 55.2 Pa, or 552 ppmv. This is the good news; on a shortish time scale, the ocean can soak up our anthropogenic excess.

Beyond the century time scale, we have $p - p_s \approx v_A - \varepsilon p$, so that s is given by

$$\dot{s} \approx \delta \left[-\frac{v_A}{\varepsilon} + \alpha \left\{ v_A + \frac{(1-vs)^2}{s} \right\} + \Omega \right].$$
 (3.15)

The behaviour of s is altered dramatically. It decreases almost linearly for a dimensional time

$$\frac{\varepsilon t_0}{\delta v_A} = \frac{m_{oc} S_0}{I} \approx 670 \text{ y}, \qquad (3.16)$$

when it fairly abruptly reaches an equilibrium in which

$$s = \frac{\varepsilon \alpha}{v_A} \sigma, \quad p = \frac{v_A}{\varepsilon \alpha} \Pi,$$
 (3.17)

so that $s \sim 4.8 \times 10^{-3}$ and $p \sim 208$. Thereafter, $\sigma \approx \frac{1}{\Pi}$, and Π increases toward an eventual equilibrium of

 $\Pi = 1$ on a time scale of $t \sim \frac{1}{e^{\alpha}}$, corresponding to 27,000 years. Illustrations of the short-term (millennial) response are shown in Fig. 1.

In Fig. 2, we show the results of releasing the estimated reserves of 5,000 GtC (ArcHer *et al.* 1997; we take 4,800 GtC) in different ways. The different

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Evolution of CO_2 , CO_3^{2-} (mM), and pH over the next millennium under conditions of continuing production of CO_2 at present levels



Figure 2

Evolution of p_{CO_2} and pH under different emissions scenarios. The labels have this meaning: 'default' means continued release at 'present' levels (1990s, 6 GtC y⁻¹) indefinitely; 1,800 indicates 6 GtC y⁻¹ for 800 years and then zero emissions (switch-off); 2,400 means 12 GtC y⁻¹ for 400 years and then a switch-off; 4,200 means 24 GtC y⁻¹ for 200 years and then a switch-off. After the switch-off there are a further few decades while the ocean and atmosphere come into quasi-equilibrium, which is then maintained for thousands of years.

release mechanisms yield no ultimate difference, either in atmospheric CO_2 or in ocean pH; the only difference is one of time scale.

The message is simply stated. The carbon that we pour into the atmosphere pours into the ocean, but the maintenance of charge neutrality forces the carbonate ion to decrease dramatically to the level of $\approx 1 \mu M$, a drop of about a factor of 200. Consulting (2.4), we see that the hydrogen ion concentration increases by 200, as does the dissolved CO₂. This causes a drop of

pH of about 2.3 to 6.9, while the atmospheric CO₂, which follows ocean CO₂ via Henry's law in (3.14), jumps by about 200 to a dimensionless value of around 166, corresponding to 5,800 Pa, or 58,000 ppmv: 6 % of the atmosphere! Mankind would be long extinct by then. Of course, this assumes unlimited reserves and their continuing use; the more realistic case of finite reserves results in a saturation at 110 Pa, or 1,100 ppmv, and ocean pH levels at 7.5, as seen in Fig. 2.

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Figure 3 The ocean bicarbonate buffering system

4. Conclusions

The main purpose of this article was to make explicit in simple mathematical terms what the consequences of continued anthropogenic CO2 output will be. Our model is stripped to its essentials, and consists only of two ordinary differential equations, for atmospheric CO₂ and oceanic CO₂. Yet this model yields essentially the same kind of quantitative result as do more complex simulation models, as used for example by Archer et al. (1997) and ZEEBE et al. (2008). Specifically, the same results of rapid CO2 rise and rapid pH decline are seen, with much slower recovery following the termination of carbon emissions. However, unlike the computational models, the simplicity of our approach enables the mechanisms of warming and acidification to be portrayed more clearly, and the time scales of the responses can be explicitly identified.

In its essence, the message is this. Increasing atmospheric CO_2 causes increased flux to the ocean, until a quasi-equilibrium is reached after a century or two when the ocean uptake balances the discharge to the atmosphere. However, when a mole of CO_2 is dissolved, the bicarbonate buffering system (Fig. 3) immediately empties it into the bicarbonate pool; this would cause charge imbalance, and thus at the same time, a carbonate ion is also transferred to the bicarbonate pool. In effect, the reaction is that described by ArcHER *et al.* (1997):

$$CO_2 + CO_3^{2-} + H_2O \rightarrow 2HCO_3^{-}.$$
 (4.1)

Thus, the flux of CO_2 to the ocean provides also an equivalent flux from the carbonate pool.

Because the acidity is controlled through the buffering relation (2.4)

$$[H^{+}] \propto \frac{[HCO_{3}^{-}]}{[CO_{3}^{2-}]},$$
 (4.2)

[H⁺] increases, i. e., pH decreases, resulting in ocean acidification. But at the same time, because

$$[CO_2] \propto \frac{[HCO_3^-]^2}{[CO_3^{2^-}]},$$
 (4.3)

the dissolved CO_2 increases; not directly because of the flux from the atmosphere, but because of the decrease in carbonate, and the rate of increase is faster than linear. This second consequence of carbonate loss is what causes the accelerated increase of atmospheric CO_2 , and its attendant greenhouse effect, following the initial lull.

Our model simulations in Fig. 1 are based on a constant input, while ZEEBE *et al.* (2008) consider the input of a fixed quantity corresponding to the total reserves, perhaps 5,000 GtC. It is straightforward to read from (3.15) what happens in this case. The carbonate declines until the ocean has absorbed most of the atmospheric load, but the subsequent recovery of the ocean carbonate takes place over a much longer time scale, that being tens of thousands of years.

As with all models, and particularly this one, the results must be tempered with caution. In particular, the neglect of ocean mixing requires consideration, most notably as regards the vertical profile. With a mixing time on the order of 1,000 years, carbonate changes on shorter time scales than this will be restricted to the upper parts of the ocean. The practical effect is to reduce the effective value of m_{oc} in (3.1), and thus increase δ in (3.6). Consulting (3.16), this suggests that the drawdown in carbonate ion content will be even faster, and further study of this issue is warranted.

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