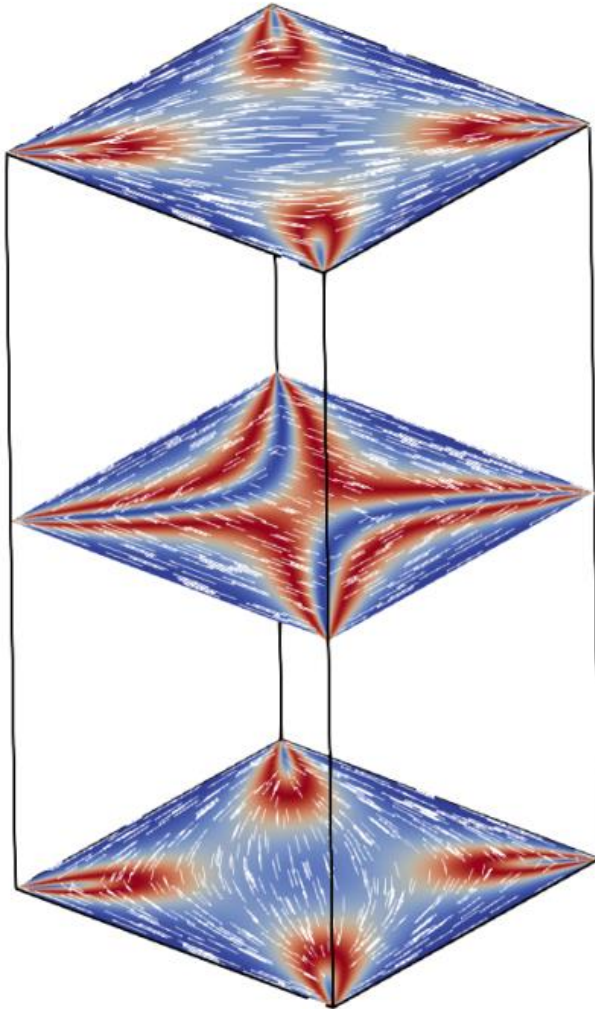




THE WELL ORDER RECONSTRUCTION SOLUTION FOR THREE- DIMENSIONAL WELLS

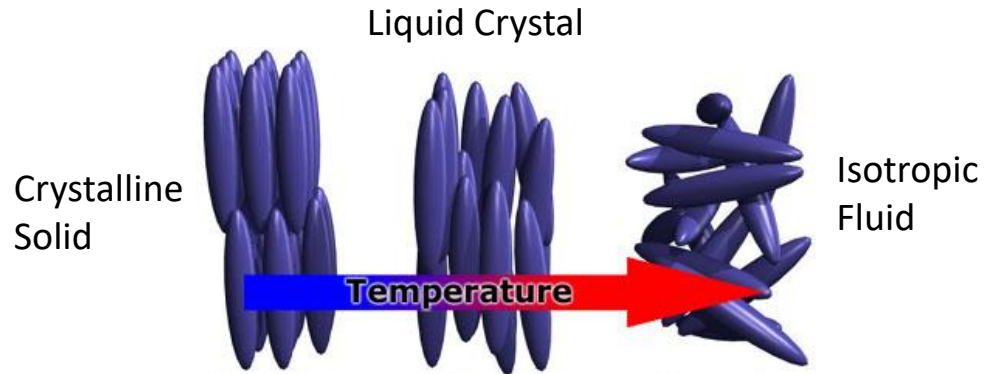
Joseph Harris

Joint work with Apala Majumdar (PhD
Supervisor), Giacomo Canevari and Yiwei
Wang

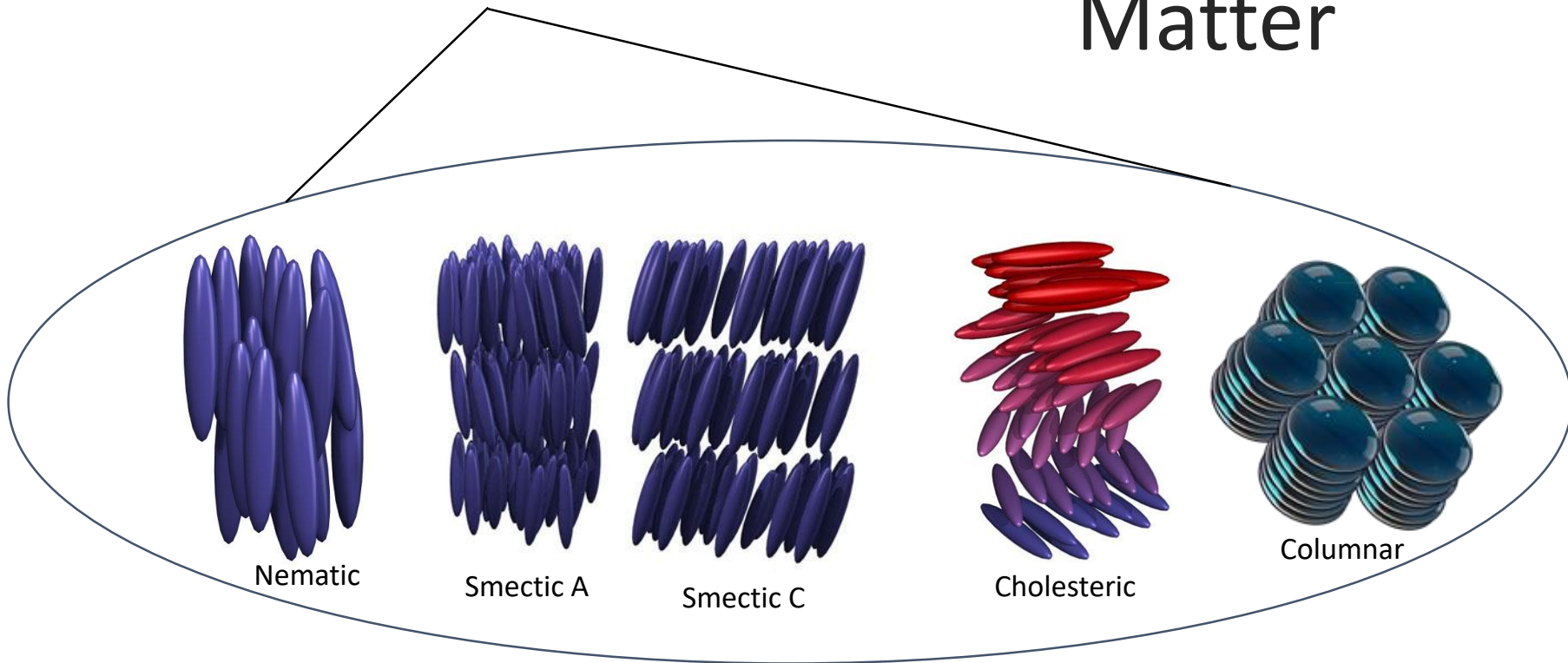


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An Intermediate Phase of Matter



THE LANDAU-DE GENNES MODEL

We describe the nematic state by the macroscopic order parameter \mathbf{Q} (traceless, symmetric 3x3 matrix).

In the absence of surface energies, we minimize the LdG energy given by

$$\mathcal{F}_\lambda[\mathbf{Q}] = \iiint \frac{1}{2} |\nabla \mathbf{Q}|^2 + \frac{\lambda^2}{L} f_b(\mathbf{Q}) dV$$

where λ is a characteristic length scale, $L > 0$ a material-dependent elastic constant.

The thermotropic bulk potential is given by

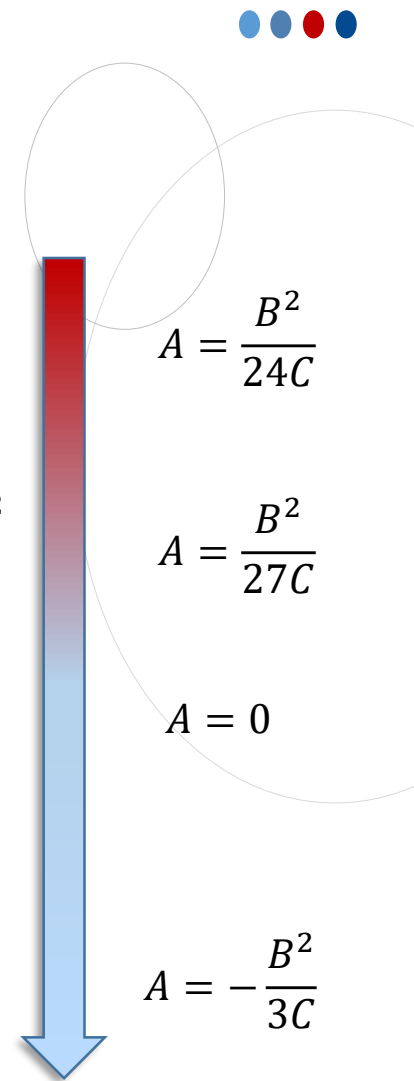
$$f_b(\mathbf{Q}) = \frac{A}{2} \text{tr} \mathbf{Q}^2 - \frac{B}{3} \text{tr} \mathbf{Q}^3 + \frac{C}{4} (\text{tr} \mathbf{Q}^2)^2$$

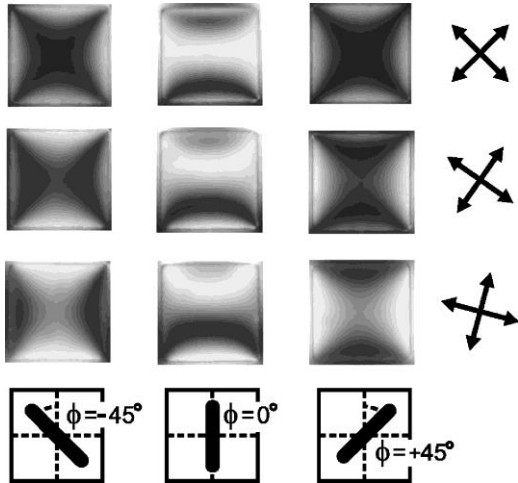
For a given $A < 0$, minimizers of f_b have the form:

$$\mathbf{Q} = s_+ \left(\mathbf{n} \otimes \mathbf{n} - \frac{I}{3} \right), \quad \mathbf{n} \in S^2$$

where

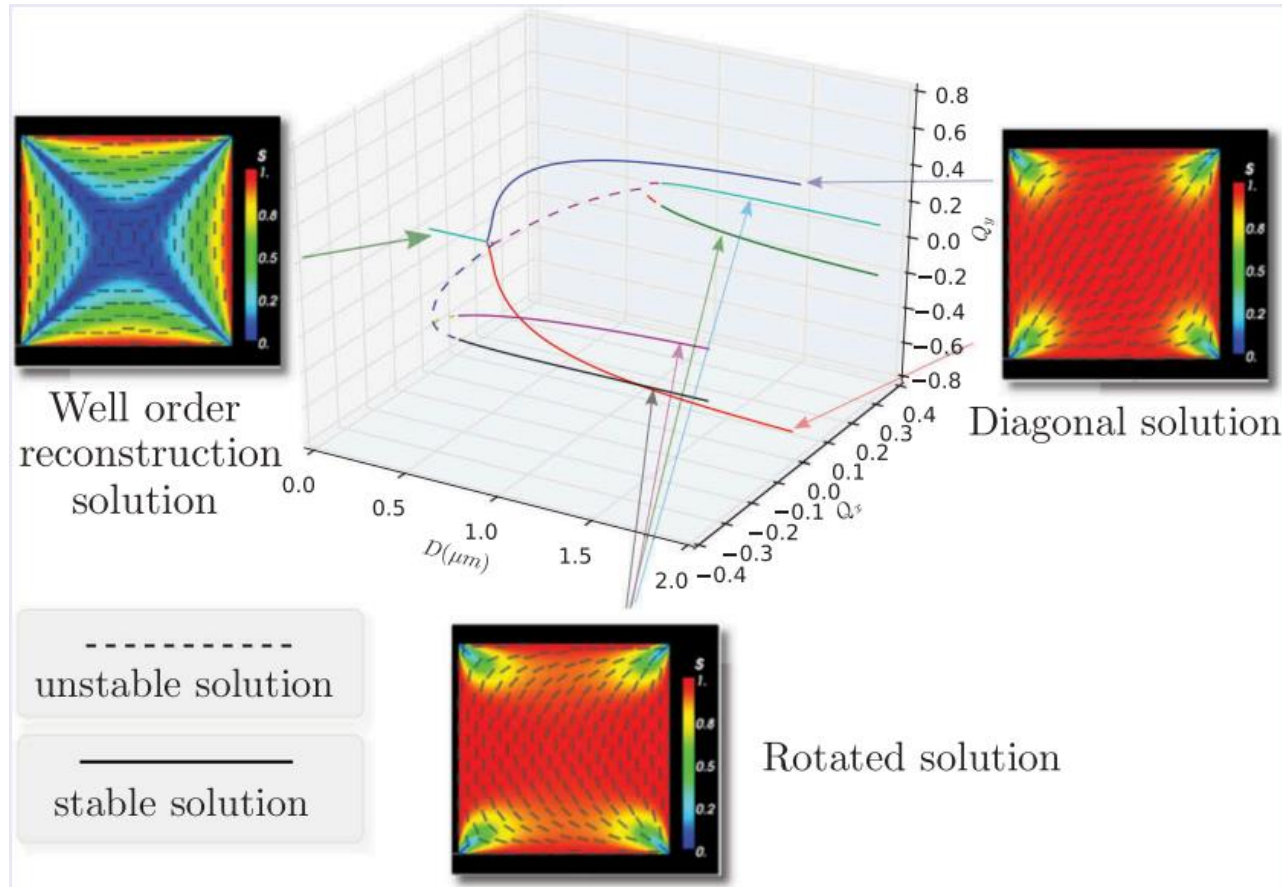
$$s_+ = \frac{B + \sqrt{B^2 + 24|A|C}}{4C}$$





Tsakonas et.al.

(2007). Multistable alignment states in Nematic liquid crystal filled wells. *Appl. Phys. Lett*, 1-18. <https://doi.org/10.1063/1.2713140>



Robinson et.al. (2017). From molecular to continuum modelling of bistable liquid crystal devices. *Liquid Crystals*, 1-18. <https://doi.org/10.1080/02678292.2017.1290284>



THE WELL ORDER RECONSTRUCTION SOLUTION (WORS)

The WORS is an interesting nematic liquid crystal equilibria for two reasons:

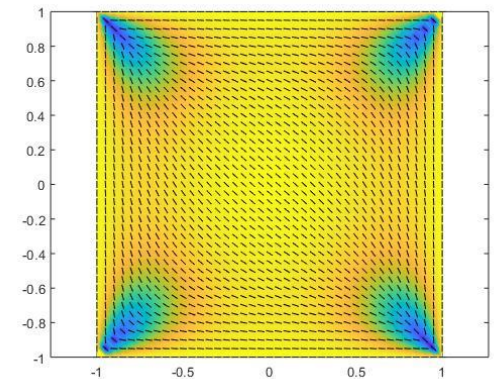
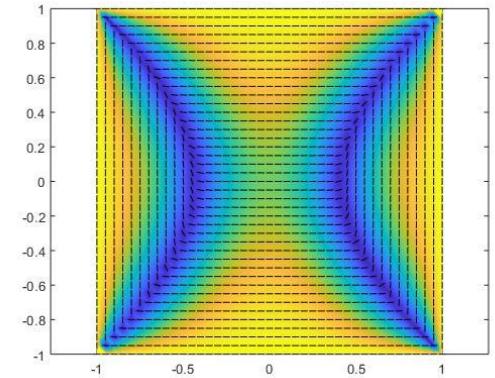
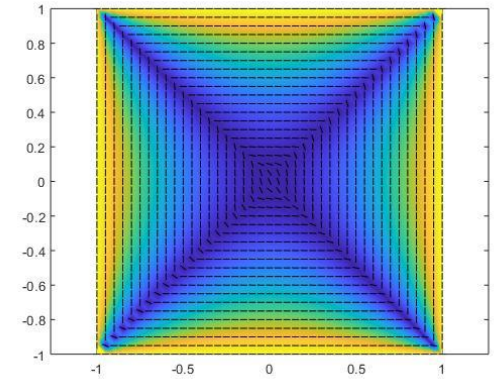
- It partitions the square into four quadrants and the nematic director is approximately constant in each quadrant according to the tangent condition on the corresponding edge
- The WORS has a defect line along each square diagonal and the two intersect at the centre, yielding the quadrant structure

The WORS has a constant eigenframe and is uniaxial along the square diagonals with negative order parameter and can therefore be described by solutions of the form –

$$\mathbf{Q} = q_1(\mathbf{n}_1 \otimes \mathbf{n}_1 - \mathbf{n}_2 \otimes \mathbf{n}_2) + q_3(\hat{\mathbf{z}} \otimes \hat{\mathbf{z}} - \mathbf{I}/3)$$



- The WORS was numerically reported by **Kralj, Majumdar (2014)** in 2D
- This was analysed further in **Majumdar et.al. (2016)**, **Canevari et.al. (2017)** and **Wang et.al. (2018)** at a fixed temperature
- As proven in **Wang et.al. (2018)**, the WORS survives in the *thin film limit*
- It is our aim to observe the WORS for square wells with a finite height, exemplifying the 3D relevance of WORS-type solutions for all temperatures below the nematic supercooling temperature (i.e. the deep nematic regime $A < 0$)



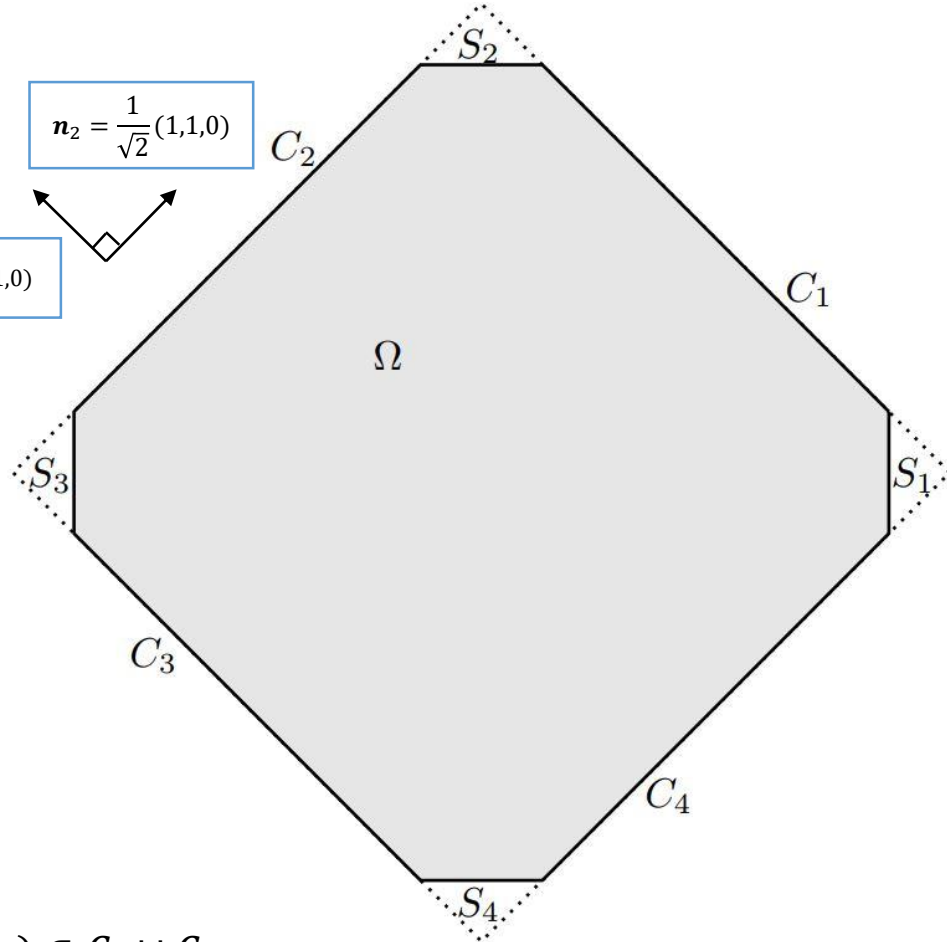
BOUNDARY CONDITIONS ON THE LATERAL SURFACES



Define $V = \Omega \times (0, \epsilon)$,
for ϵ finite.

$$\mathbf{n}_1 = \frac{1}{\sqrt{2}}(-1, 1, 0)$$

$$\mathbf{n}_2 = \frac{1}{\sqrt{2}}(1, 1, 0)$$



We impose tangent uniaxial Dirichlet conditions on the lateral surfaces of the well:

$$\mathbf{Q} = \mathbf{Q}_b \quad \text{on} \quad \partial\Omega \times (0, \epsilon)$$

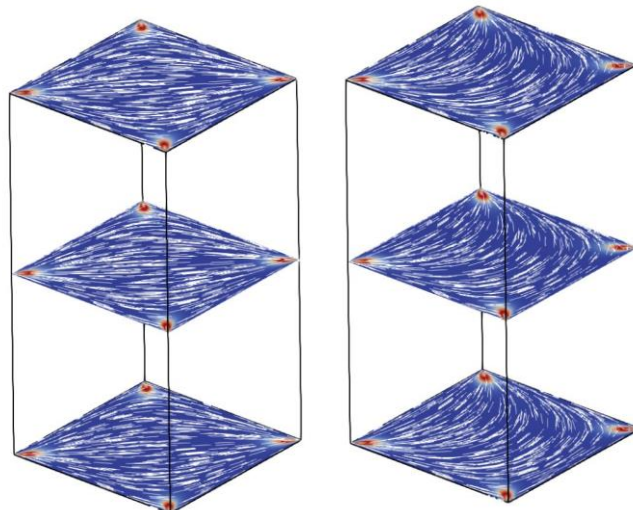
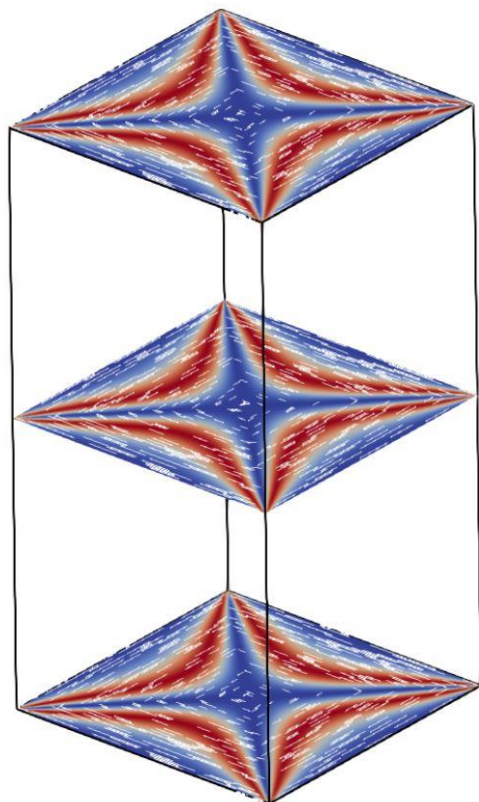
$$\mathbf{Q}_b(x, y, z) := \begin{cases} s_+ \left(\mathbf{n}_1 \otimes \mathbf{n}_1 - \frac{\mathbf{I}}{3} \right) & \text{if } (x, y) \in C_1 \cup C_3 \\ s_+ \left(\mathbf{n}_2 \otimes \mathbf{n}_2 - \frac{\mathbf{I}}{3} \right) & \text{if } (x, y) \in C_2 \cup C_4 \end{cases}$$



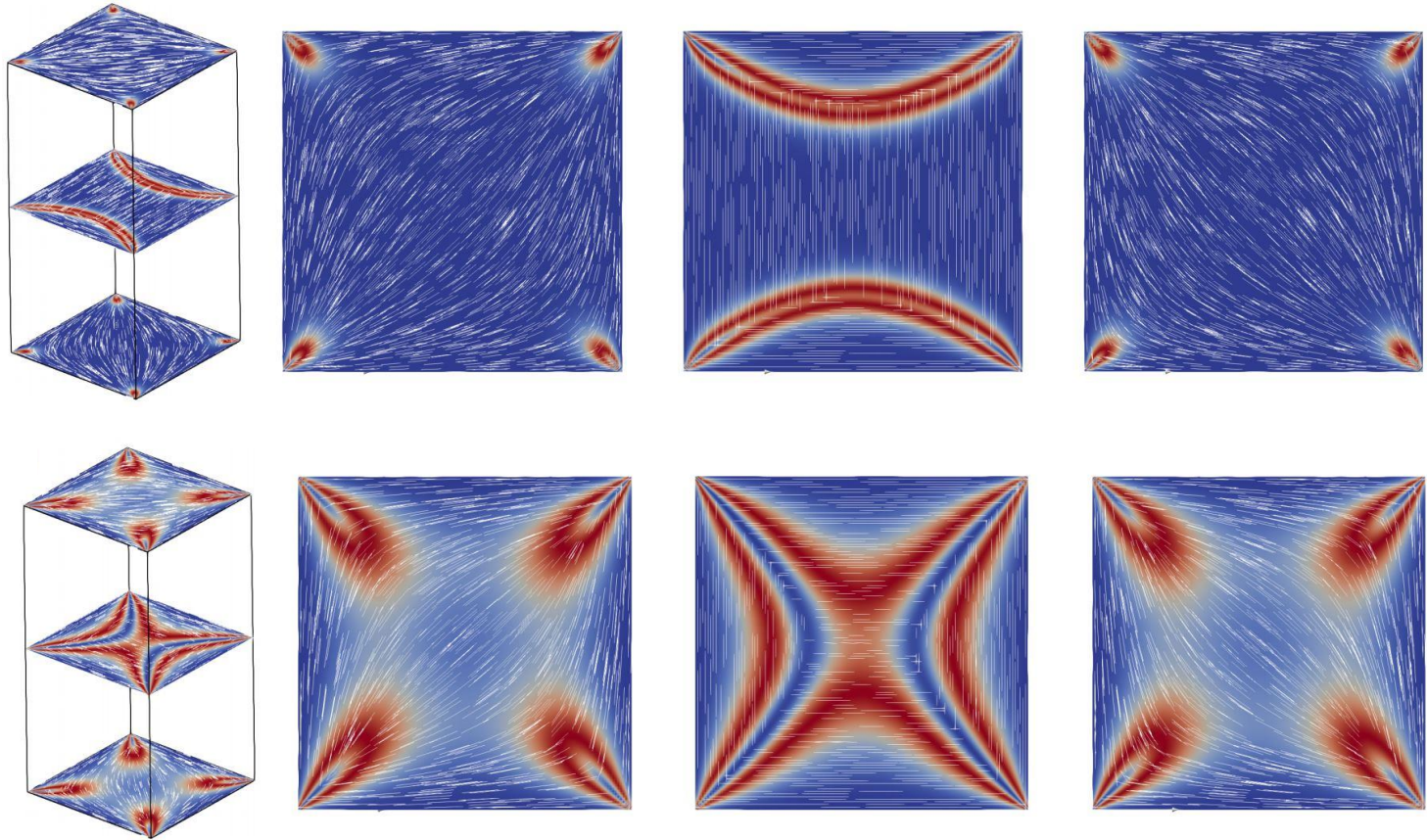
NATURAL BOUNDARY CONDITIONS ON THE TOP AND BOTTOM PLATES

$$\partial_z \mathbf{Q} = 0 \quad \text{on} \quad \Omega \times \{0, \epsilon\}$$

For sufficiently small cross section size, λ , the WORS is the global LdG minimizer for these 3D problems for all temperatures below the nematic supercooling temperature.



WORS, diagonal and rotated solutions with $|\partial_z \mathbf{Q}|^2 \approx 10^{-12}$. With square cross section $\lambda^2 = 5$ (left), and $\lambda^2 = 100$ (right).



For large enough well height, ϵ , and relatively large λ , we demonstrate the existence of stable mixed 3D solutions with two different diagonal profiles on the top and bottom plates.

Mixed 3D solutions for $\lambda^2 = 100$ (top), and $\lambda^2 = 10$ (bottom), with cross sections at $z = 0, 2, 4$.



SURFACE ANCHORING ON THE TOP AND BOTTOM PLATES

$$\mathcal{F}_\lambda[\mathbf{Q}] = \iiint \frac{1}{2} |\nabla \mathbf{Q}|^2 + \frac{\lambda^2}{L} f_b(\mathbf{Q}) dV + \frac{\lambda}{L} \int f_s(\mathbf{Q}) dS$$

$$f_s(\mathbf{Q}) = \alpha_z \left(\mathbf{Q} \hat{\mathbf{z}} \cdot \hat{\mathbf{z}} + \frac{s_+}{3} \right)^2 + \gamma_z |(\mathbf{I} - \hat{\mathbf{z}} \otimes \hat{\mathbf{z}}) \mathbf{Q} \hat{\mathbf{z}}|^2, \quad \alpha_z, \gamma_z > 0$$

f_s favours \mathbf{Q} -tensors that have $\hat{\mathbf{z}}$ as an eigenvector with constant eigenvalue $-\frac{s_+}{3}$ on the top and bottom plates.

For sufficiently small λ , there is a unique WORS-type solution to this 3D problem.

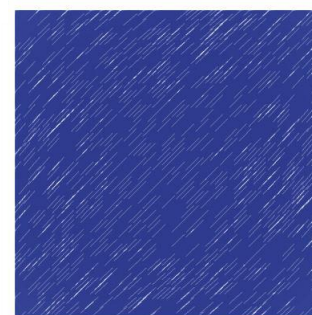
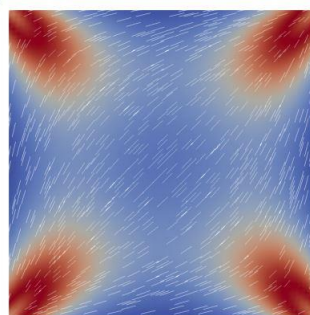
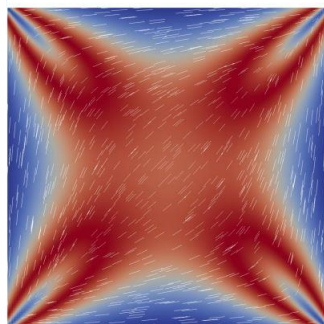
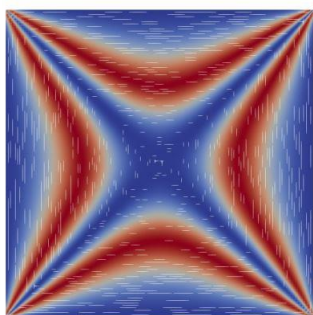


WEAK ANCHORING ON THE LATERAL SURFACES

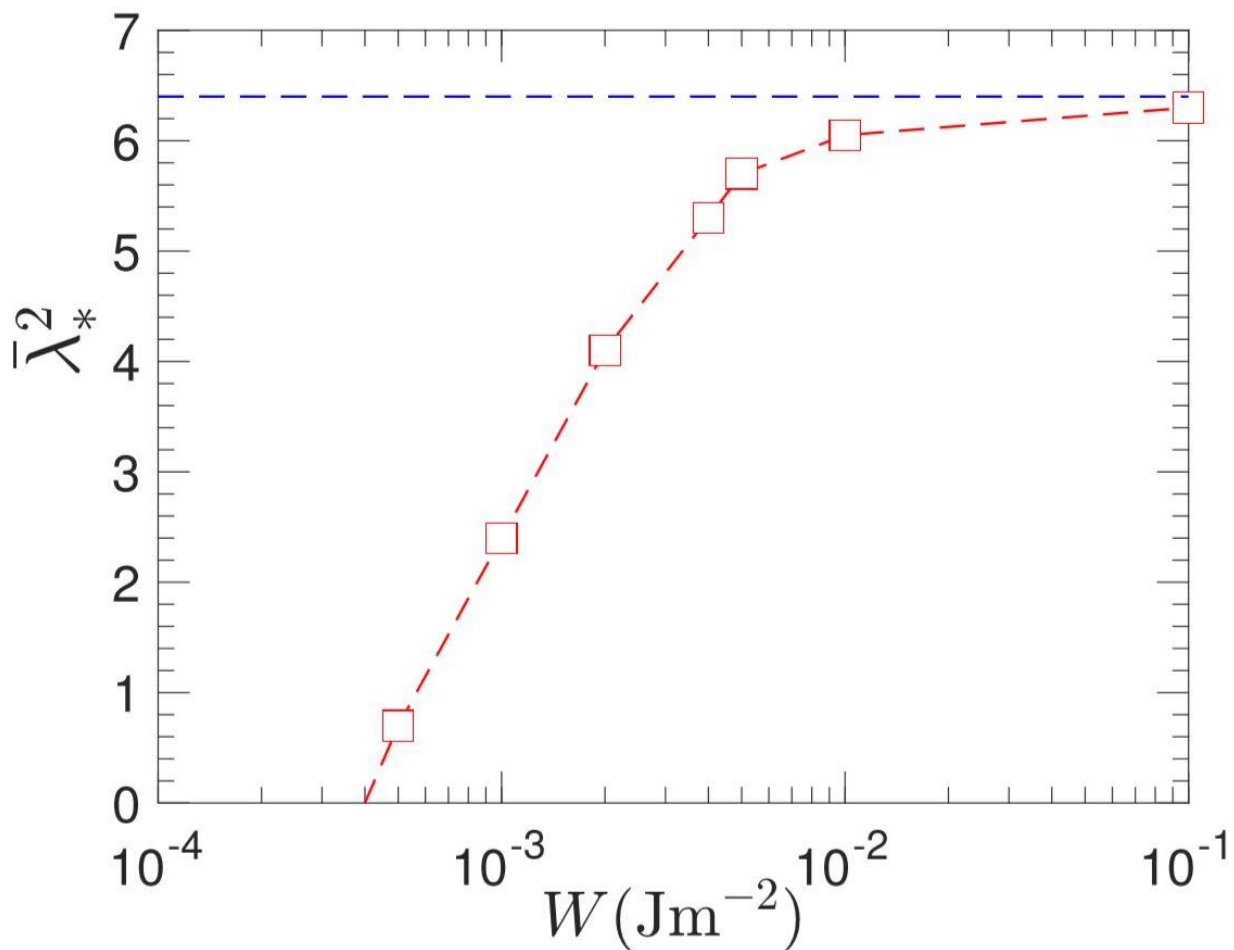
$$f_s(\mathbf{Q}) = \begin{cases} \frac{W_1\lambda}{L} \left(\mathbf{Q} - g(x) \left(\hat{\mathbf{x}} \otimes \hat{\mathbf{x}} - \frac{I}{3} \right) \right)^2, \\ \frac{W_2\lambda}{L} \left(\mathbf{Q} - g(y) \left(\hat{\mathbf{y}} \otimes \hat{\mathbf{y}} - \frac{I}{3} \right) \right)^2, \end{cases}$$

$$g(x) = s_+, \quad \forall x \in [-1 + \delta, 1 - \delta]$$
$$y = 0, 1$$
$$x = 0, 1$$

Eliminates discontinuity at the corners.



Transition from the WORS to a diagonal solution by weakening the anchoring strength on the lateral surfaces for $W_i = 10^{-2} Jm^{-2}$ (left) to $W_i = 10^{-4} Jm^{-2}$ (right).



Bifurcation points, below which the WORS is the unique solution as a function of anchoring strength. The blue dashed line indicates the bifurcation point of the WORS for Dirichlet boundary condition in a 2D square well.

- This talk is based on joint work with collaborators **Giacomo Canevari**, **Apala Majumdar** and **Yiwei Wang**.
- Submitted paper: The Well Order Reconstruction Solution for Three-Dimensional Wells, in the Landau-de Gennes theory (2019).
- Preprint submitted to the **International Journal of Nonlinear Mechanics**.



THANK YOU



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[arXiv:1903.03873](https://arxiv.org/abs/1903.03873)

