# Computational Algebraic Topology Topic B: Lecture II: The Topology of Paradox 

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This point of view is called non-contextuality. It is equivalent to the assumption of a classical source.

However, this view is impossible to sustain in the light of our actual observations of (micro)-physical reality.

## Hidden Variables: The Mermin instruction set picture



## The 'Hardy Paradox': Bell's theorem without inequalities

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 Hardy models: those whose support satisfies|  | $(0,0)$ | $(0,1)$ | $(1,0)$ | $(1,1)$ |
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| $\left(a_{1}, b_{1}\right)$ | 1 |  |  |  |
| $\left(a_{1}, b_{2}\right)$ | 0 |  |  |  |
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Which 'instruction set' $\lambda$ could the outcomes $(0,0)$ for measurements $\left(a_{1}, b_{1}\right)$ could come? Clearly, we must have

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Thus Hardy models are contextual. They cannot be explained by a classical source.

## Bundle Pictures

## Logical Contextuality

- Ignore precise probabilities
- Events are possible or not
- E.g. the Hardy model:

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Note this achieves the algebraic maximum of 4 for our logical Bell inequality.

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Note this achieves the algebraic maximum of 4 for our logical Bell inequality. In terms of the XOR game, it is a winning strategy.

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## Visualizing Contextuality



The Hardy table and the PR box as bundles

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A hierarchy of degrees of contextuality:

$$
\text { Bell }<\text { Hardy }<\text { GHZ }
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## Contextuality, Logic and Paradoxes

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Liar cycles. A Liar cycle of length $N$ is a sequence of statements
$S_{1}: S_{2}$ is true,
$S_{2}: S_{3}$ is true,
$S_{N-1}: S_{N}$ is true,
$S_{N}: S_{1}$ is false.
For $N=1$, this is the classic Liar sentence

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S: S \text { is false. }
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Following Cook, Walicki et al. we can model the situation by boolean equations:

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S_{N-1}: S_{N} \text { is true, } \\
S_{N}: S_{1} \text { is false. }
\end{array}
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The "paradoxical" nature of the original statements is now captured by the inconsistency of these equations.

## Contextuality in the Liar; Liar cycles in the PR Box

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We can regard each of these equations as fibered over the set of variables which occur in it:

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\left\{x_{2}, x_{3}\right\}: & x_{2}=x_{3} \\
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Up to rearrangement, the Liar cycle of length 4 corresponds exactly to the PR box.

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Any subset of up to $n-1$ of these equations is consistent; while the whole set is inconsistent.

Up to rearrangement, the Liar cycle of length 4 corresponds exactly to the PR box.

The usual reasoning to derive a contradiction from the Liar cycle corresponds precisely to the attempt to find a univocal path in the bundle diagram.

## Paths to contradiction



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We have discussed a specific case here, but the analysis can be generalised to a large class of examples.

## Formalizing Contextuality: Measurement Scenarios

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A joint outcome or event in a context $C$ is $s \in O^{C}$, e.g. $s=\{a \mapsto 0, b \mapsto 1\}$.

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| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $A$ | $A$ | $H$ | $H$ | $B$ | $I$ | $P$ | $P$ | $Q$ |
| $B$ | $E$ | $I$ | $K$ | $E$ | $K$ | $Q$ | $R$ | $R$ |
| $C$ | $F$ | $C$ | $G$ | $M$ | $N$ | $D$ | $F$ | $M$ |
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The original K-S construction used 117 variables!

## Empirical Models

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These distributions are the rows of our probability tables.

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Each row of the table specifies a probability distribution on events $O^{C}$ for a given choice of measurements $C$.

## Gluing functional sections



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If $s_{U}\left|u \cap v=s_{v}\right| u \cap V$, they can be glued to form

$$
s: U \cup V \longrightarrow O
$$

such that $\left.s\right|_{U}=s_{U}$ and $\left.s\right|_{V}=s_{V}$.

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We want to do this by saying that the distributions "agree on overlaps". For all $C, C^{\prime} \in \mathcal{M}$ :

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A formula for restriction of distributions: if $C^{\prime} \subseteq C, d \in \operatorname{Prob}\left(O^{C}\right)$,

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So compatibility says that the distributions on different contexts have consistent marginals.

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- Suppose that $C=\{a, b\}$, and $C^{\prime}=\left\{a, b^{\prime}\right\}$, where $a$ is a variable measured by an agent Alice, while $b$ and $b^{\prime}$ are variables measured by Bob, who may be spacelike separated from Alice.


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- This is captured by saying that the distribution on $\{a\}=\{a, b\} \cap\left\{a, b^{\prime}\right\}$ is the same whether we marginalize from the distribution $e_{C}$, or the distribution $e_{C^{\prime}}$.
- This condition is generalized by compatibility - and this general form is satisfied by quantum systems.

No-Signalling for Alice-Bob Tables

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Consider the following schematic representation of an Alice-Bob table:

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where we have labelled the entries with the letters $c, \ldots, r$.
The no-signalling conditions for the non-empty intersections of contexts are given by the following equations:
$c+e=k+m$,
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You can check that these conditions are satisfied by the Bell table.
Moreover, the PR box has a unique family of distributions which satisfy these conditions.

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An empirical model $\left\{d_{c}\right\}_{c \in \mathcal{M}}$ on a measurement scenario $(X, \mathcal{M}, O)$ is non-contextual if there is a distribution $d \in \operatorname{Prob}\left(O^{X}\right)$ such that, for all $C \in \mathcal{M}$ :

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If no such global section exists, the empirical model is contextual.
The import of Bell's theorem and similar results is that there are empirical models arising from quantum mechanics which are contextual.

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Compatibility is in fact the general form of an important physical principle known as No-Signalling, which ensures the consistency of quantum mechanics with Special Relativity.

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It turns out that all quantum realizable models are compatible.
Compatibility is in fact the general form of an important physical principle known as No-Signalling, which ensures the consistency of quantum mechanics with Special Relativity.

However, there are compatible (i.e. No-Signalling) empirical models which are not quantum realizable.

## Classes of Empirical Models

There is a class of empirical models, for each measurement scenario $(X, \mathcal{M}, O)$, which are quantum realizable.

That is, we can find quantum states and local observables which generate the family of distributions $\left\{d_{C}\right\}_{C \in \mathcal{M}}$.

It turns out that all quantum realizable models are compatible.
Compatibility is in fact the general form of an important physical principle known as No-Signalling, which ensures the consistency of quantum mechanics with Special Relativity.

However, there are compatible (i.e. No-Signalling) empirical models which are not quantum realizable.

We thus get a strict hierarchy of empirical models:

$$
\mathrm{NC} \subset \mathrm{QM} \subset \mathrm{NS}
$$

## The PR Box

| A | B | $(0,0)$ | $(1,0)$ | $(0,1)$ | $(1,1)$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $a_{1}$ | $b_{1}$ | 1 | 0 | 0 | 1 |
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This satisfies No-Signalling, so is consistent with SR, but it is not quantum realisable.

## Empirical Models as Vectors

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\mathbf{v}=\left(\mathbf{v}_{C, s}\right)_{C \in \mathcal{M}, s \in O^{C}}, \quad \mathbf{v}_{C, s}:=d_{C}(s)
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Moreover, convex combinations of compatible models are compatible.

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Key question: find compelling principles to explain why Nature picks out the quantum set.

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The support of an empirical model $\left\{d_{C}\right\}_{C \in \mathcal{M}}$ is defined as follows. For each $C \in \mathcal{M}$, we define $\mathcal{S}(C) \subseteq O^{C}$ :

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This is equivalent to saying that, for all $C \subseteq C^{\prime}$, the restriction map

$$
\rho_{C}^{C^{\prime}}: \mathcal{S}\left(C^{\prime}\right) \longrightarrow \mathcal{S}(C) \quad::\left.s \mapsto s\right|_{C}
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is surjective.

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Thus in terms of well-known quantum examples, we have

$$
\text { Bell }<\text { Hardy }<\mathrm{GHZ}
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