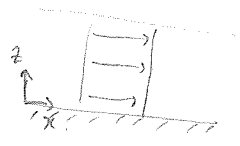


1. (i) $\underline{u} = \begin{pmatrix} U \\ 0 \\ 0 \end{pmatrix}$

$\Rightarrow \nabla \cdot \underline{u} = 0$

$\frac{\partial u_i}{\partial x_j} = \begin{pmatrix} 0 & 0 & 0 \\ U & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \Rightarrow \dot{\Sigma}_{ij} = \begin{pmatrix} 0 & 0 & 0 \\ U & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$

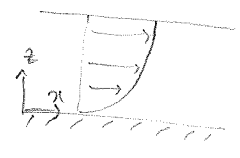


plug flow - found in rapidly sliding regions, eg. ice streams.

(ii) $\underline{u} = \lambda \begin{pmatrix} H^4 - (H-z)^4 \\ 0 \\ 0 \end{pmatrix}$

$\Rightarrow \nabla \cdot \underline{u} = 0$

$\frac{\partial u_i}{\partial x_j} = \begin{pmatrix} 0 & 0 & 4\lambda(H-z)^3 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \Rightarrow \dot{\Sigma}_{ij} = 2\lambda(H-z)^3 \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$

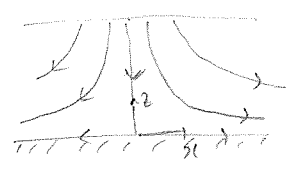


shear flows with no sliding (shallow ice approx.)
- found, eg. in a region with frozen bed.

(iii) $\underline{u} = \frac{a}{H} \begin{pmatrix} x/2 \\ y/2 \\ -z \end{pmatrix}$

$\Rightarrow \nabla \cdot \underline{u} = \frac{1}{2} + \frac{1}{2} - 1 = 0$

$\frac{\partial u_i}{\partial x_j} = \frac{a}{H} \begin{pmatrix} \frac{1}{2} & 0 & 0 \\ 0 & \frac{1}{2} & 0 \\ 0 & 0 & -1 \end{pmatrix} = \dot{\Sigma}_{ij}$

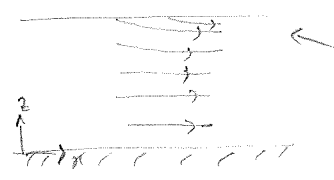


ice dome (centre of an ice sheet).

(iv) $\underline{u} = \begin{pmatrix} U \\ 0 \\ -\kappa e^{\frac{z-H}{d}} \end{pmatrix}$

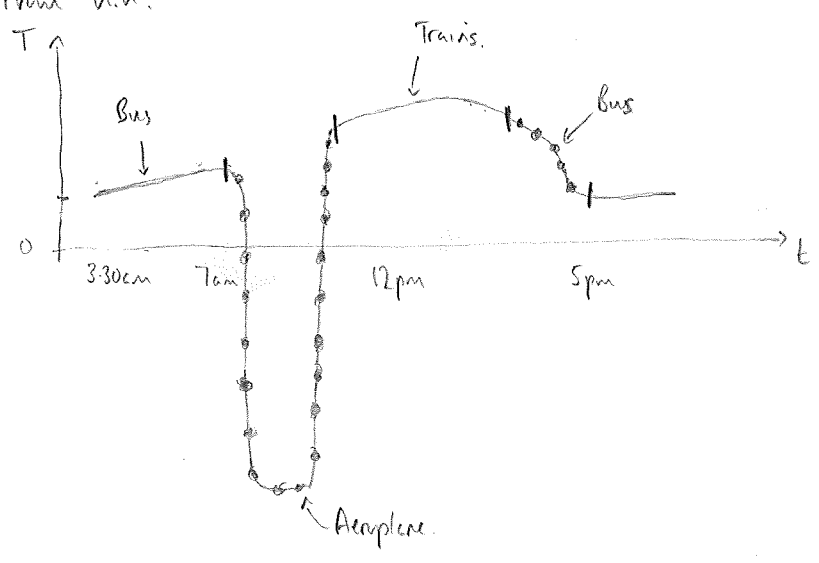
$\Rightarrow \nabla \cdot \underline{u} = -\frac{\kappa}{d} e^{\frac{z-H}{d}}$

$\frac{\partial u_i}{\partial x_j} = -\frac{\kappa}{d} e^{\frac{z-H}{d}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \dot{\Sigma}_{ij}$



compacting firn layer near surface

2. From U.K.



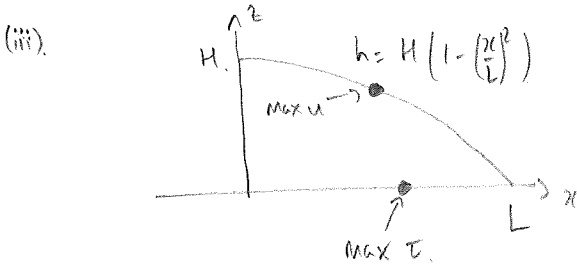
dotted periods dominated by advection $\underline{u} \cdot \nabla T$
other periods dominated by $\frac{\partial T}{\partial t}$

3. (i) $p=0, \tau=0$ at $z=h$ (atmospheric pressure taken as zero)

$$0 = -\frac{\partial p}{\partial z} - \rho g \Rightarrow p = \rho g(h-z)$$

$$0 = -\frac{\partial p}{\partial x} + \frac{\partial \tau}{\partial z} \Rightarrow \frac{\partial \tau}{\partial z} = -\rho g \frac{\partial h}{\partial x} \Rightarrow \tau = \left(-\rho g \frac{\partial h}{\partial x}\right)(h-z)$$

$$(ii) \frac{\partial u}{\partial z} = 2A \tau^\alpha = 2A(\rho g)^\alpha \left(-\frac{\partial h}{\partial x}\right)^\alpha (h-z)^\alpha \Rightarrow u = \frac{2A(\rho g)^\alpha}{\alpha+1} \left(-\frac{\partial h}{\partial x}\right)^\alpha \left[h^{\alpha+1} - (h-z)^{\alpha+1}\right]$$

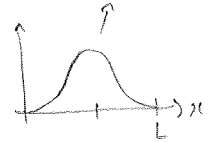


Highest velocity for each x occurs at $z=h$.
 where $u \propto \left(-\frac{\partial h}{\partial x}\right)^\alpha h^{\alpha+1} = \frac{2H^{2\alpha+1}}{L^\alpha} \left(\frac{x}{L}\right)^\alpha \left(1 - \left(\frac{x}{L}\right)^2\right)^{\alpha+1}$

which has max where.

$$2(\alpha+1)\left(\frac{x}{L}\right)^2 = \alpha \left(1 - \left(\frac{x}{L}\right)^2\right)$$

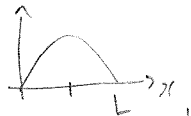
$$\Rightarrow \left(\frac{x}{L}\right)^2 = \frac{\alpha}{3\alpha+2} = \frac{3}{11} \text{ for } \alpha=3$$



Highest strain rate is also where highest shear stress occurs, which is at bed, where

$$\tau \propto \left(-\frac{\partial h}{\partial x}\right) h = \frac{2H^2}{L} \left(\frac{x}{L}\right) \left(1 - \left(\frac{x}{L}\right)^2\right)$$

$$\text{max } \tau \text{ where } 1 = 3\left(\frac{x}{L}\right)^2 \Rightarrow \left(\frac{x}{L}\right)^2 = \frac{1}{3}$$



(iv) $\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0$ with $w=0$ at $z=0$, $w+a = \frac{\partial h}{\partial t} + u \frac{\partial h}{\partial x}$ at $z=h$. (accumulation)

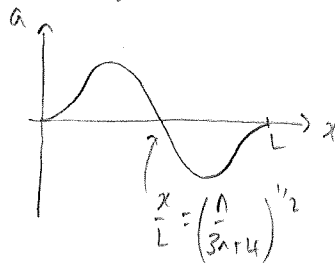
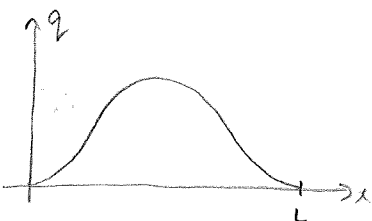
$$\Rightarrow \text{Integrating from } 0 \text{ to } h \Rightarrow \int_0^h \frac{\partial u}{\partial x} dz = -[w]_0^h = -\frac{\partial h}{\partial t} - \frac{u \partial h}{\partial x} + a$$

$$\frac{\partial}{\partial x} \int_0^h u dz = \frac{u \partial h}{\partial x} + a$$

$$\Rightarrow \frac{\partial h}{\partial t} + \frac{\partial q}{\partial x} = a$$

From above $q = \int_0^h u dz = \frac{2A(\rho g)^\alpha}{\alpha+2} \left(-\frac{\partial h}{\partial x}\right)^\alpha h^{\alpha+2} = \frac{2A(\rho g)^\alpha 2^\alpha H^{2\alpha+2}}{\alpha+2 L^\alpha} \left(\frac{x}{L}\right)^\alpha \left(1 - \left(\frac{x}{L}\right)^2\right)^{\alpha+2}$

$$\Rightarrow a = \frac{\partial q}{\partial x} = \frac{2A(\rho g)^\alpha 2^\alpha H^{2\alpha+2}}{\alpha+2 L^{\alpha+1}} \left(1 - \left(\frac{x}{L}\right)^2\right)^{\alpha+1} \left(\frac{x}{L}\right)^{\alpha-1} \left[\alpha \left(1 - \left(\frac{x}{L}\right)^2\right) - 2(\alpha+2)\left(\frac{x}{L}\right)^2\right] = 0 \text{ at } \left(\frac{x}{L}\right)^2 = \frac{\alpha}{3\alpha+4} = \frac{3}{13}$$



Height at ELA

$$H \left(1 - \frac{\alpha}{3\alpha+4}\right) = \frac{2\alpha+4}{3\alpha+4} H$$