# Glacier Hydrology



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#### Water sources

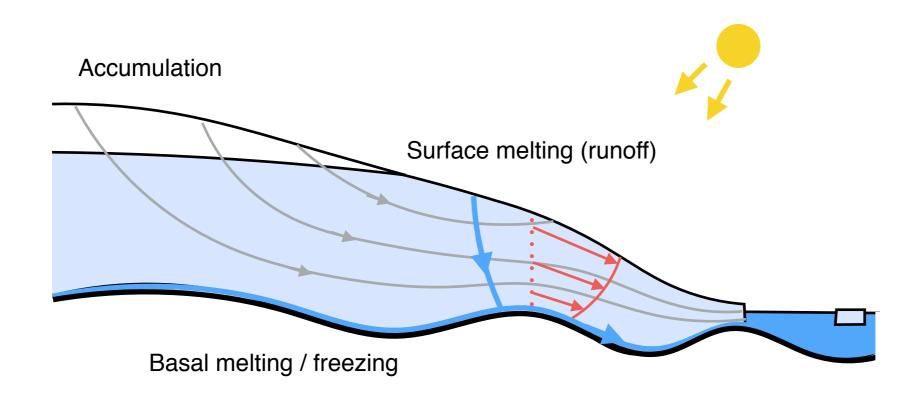
- Basal melting
- Surface melting, precipitation

#### Subglacial hydrology

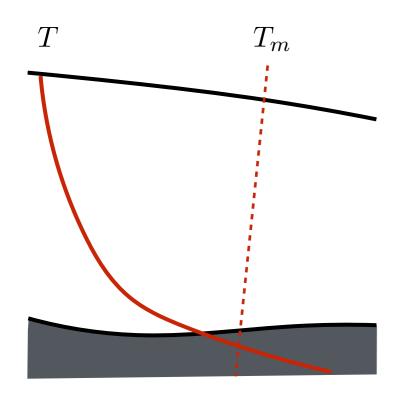
- Sheet flow
- Tunnels / channels
- Cavities
- Canals
- Lakes

Large-scale models

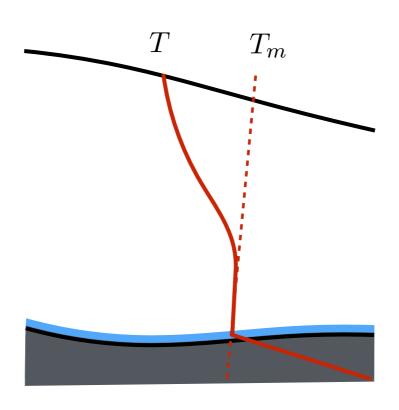
# Glacier hydrology



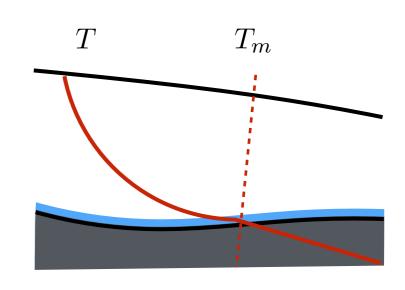
## Thermal setting



Frozen bed (no sliding?)



Basal melting



Basal freezing

Geothermal heating

$$G \sim 0.06~\mathrm{W~m^{-2}}$$

Frictional heating

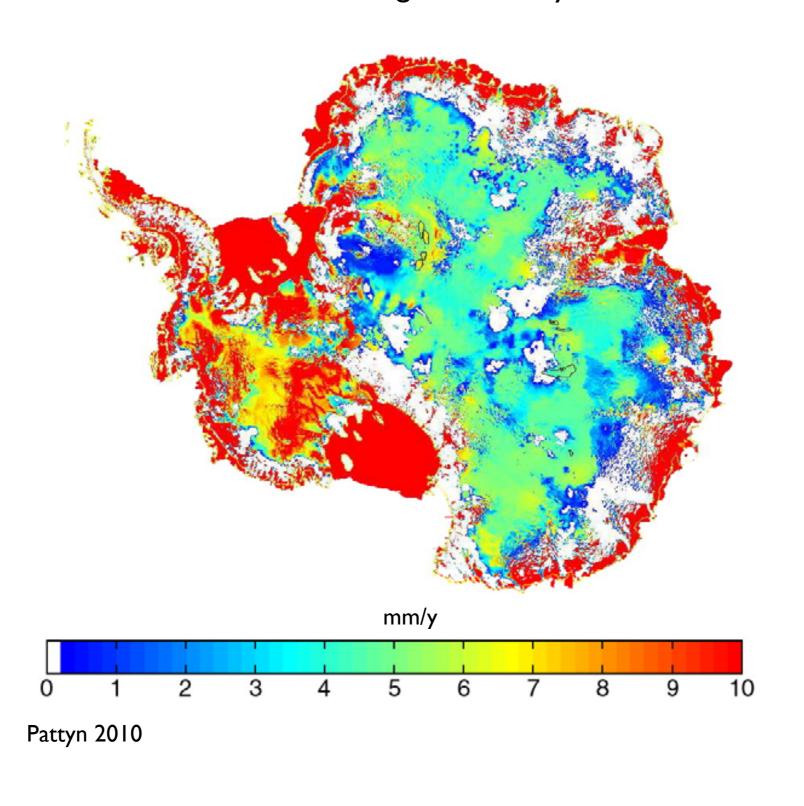
$$\tau_b \sim 100 \text{ kPa}$$
 $u_b \sim 30 \text{ m y}^{-1}$ 

$$\tau_b u_b \sim 0.1 \ \mathrm{W \ m^{-2}}$$

Basal melting  $m \sim 10 \text{ mm y}^{-1}$ 

## Water sources in Antarctica

Basal melting ~ 10 mm/y

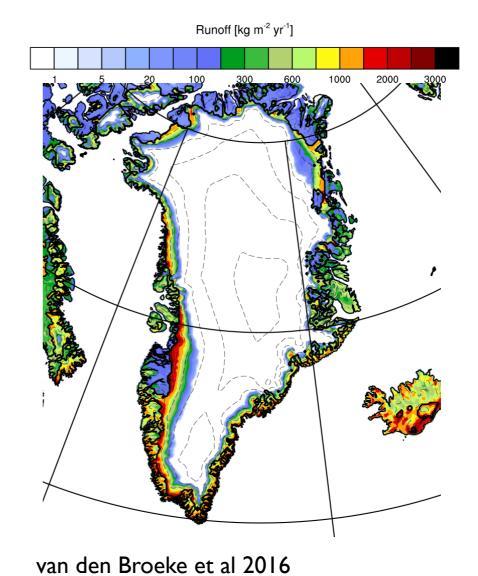


## Water sources in Greenland

#### Basal melting ~ 10 mm/y

# mm yr-I 60°W 30°W \_ 10°E 10°W W°08 20°W 70°W -60°W 55°W 50°W 45°W 40°W 35°W 30°W Aschwanden et al 2012

#### Surface runoff ~ I m/y



## Direction of subglacial water flow

#### **Hydraulic potential**

$$\phi=
ho_w g Z_b+p_w$$
 
$$=
ho_w g Z_b+
ho_i g (Z_s-Z_b)-N \qquad ext{in terms of effective pressure} \quad N=p_i-p_w$$

Potential gradient 
$$-\frac{\partial \phi}{\partial x} = \Psi + \frac{\partial N}{\partial x}$$
  $\Psi = \rho_i g \tan \alpha + (\rho_w - \rho_i) g \tan \theta$ 

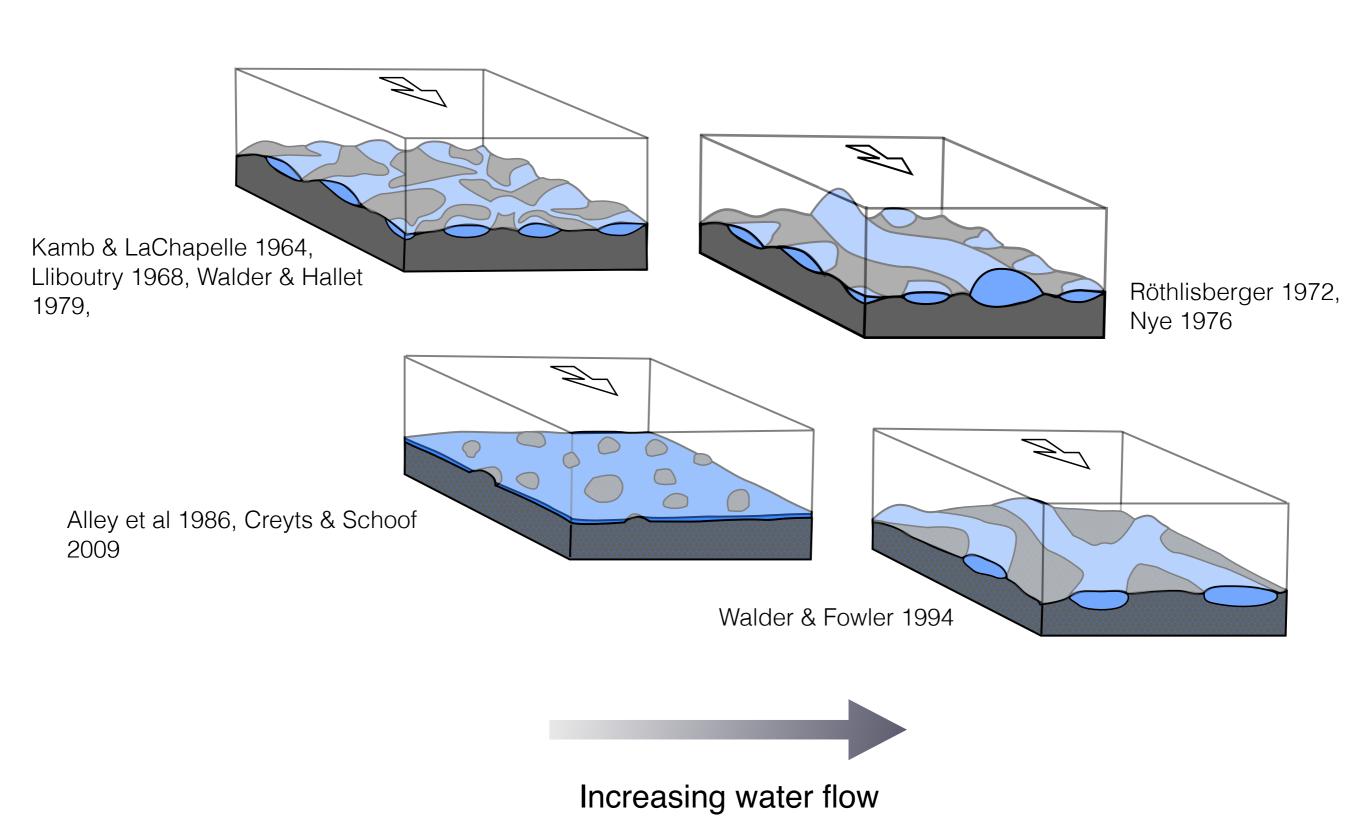
$$\Psi = \rho_i g \tan \alpha + (\rho_w - \rho_i) g \tan \theta$$



Potential gradient if basal water pressure were equal to ice pressure N=0

Predominant control on water flow direction from the ice surface slope

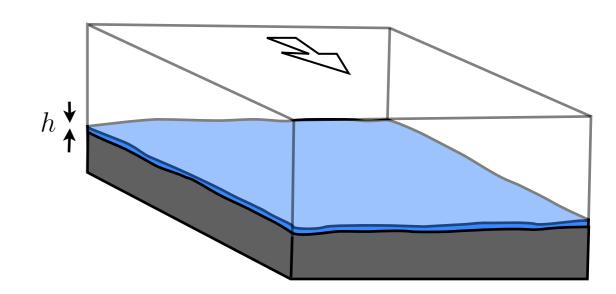
## Subglacial drainage systems



#### Weertman film Weertman 1972, Walder 1982

Weertman suggested water could flow as a film

Poiseuille flux 
$$Q = \frac{h^3}{12\mu} \left( \Psi + \frac{\partial N}{\partial x} \right)$$



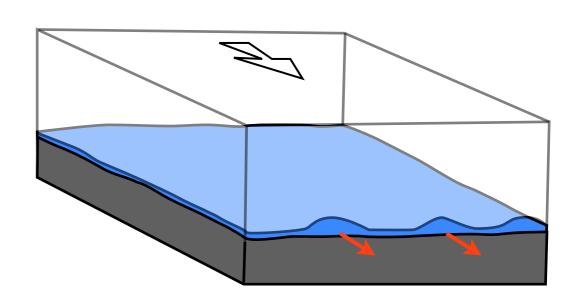
Water flow dissipates energy through heating

Leads to an instability

Larger h

→ Larger flux

→ Melting of ice roof



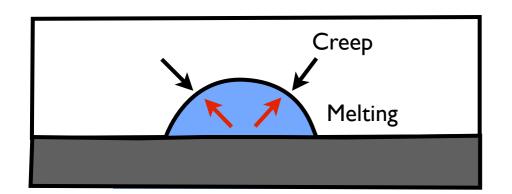
Flow wants to concentrate in localized channels / tunnels

However, a patchy film may still exist eg. Alley 1989, Creyts & Schoof 2009

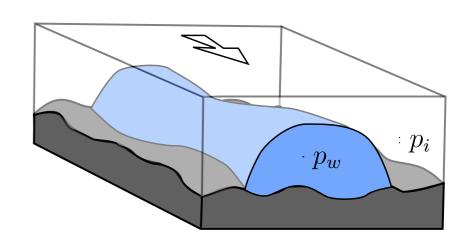


## Röthlisberger channels Röthlisberger 1972, Nye 1976

Ice wall melting is counteracted by viscous creep



$$N = p_i - p_w$$



Röthlisberger/Nye theory (ignoring pressure dependence of melting temperature)

$$\frac{\partial S}{\partial t} + \frac{\partial Q}{\partial x} = \frac{m}{\rho_w} + M$$

$$\frac{\partial S}{\partial t} = \frac{m}{\rho_i} - \tilde{A}SN^n$$

$$mL = Q\left(\Psi + \frac{\partial N}{\partial x}\right)$$

$$Q = K_c S^{4/3} \left( \Psi + \frac{\partial N}{\partial x} \right)^{1/2}$$

water mass conservation

wall evolution

local energy conservation

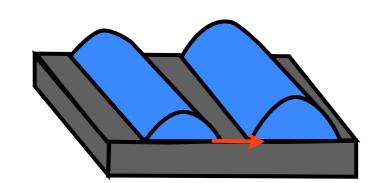
 $Q=K_cS^{4/3}\left(\Psi+\frac{\partial N}{\partial x}\right)^{1/2} \quad \text{momentum conservation}$  (turbulent flow parameterization)

Steady state ⇒

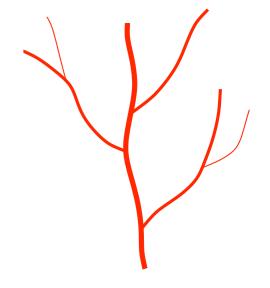
$$N \approx \left(\frac{K_c^{3/4}}{\rho_i L\tilde{A}}\right)^{1/n} \Psi^{11/8n} Q^{1/4n}$$

Effective pressure INCREASES with discharge

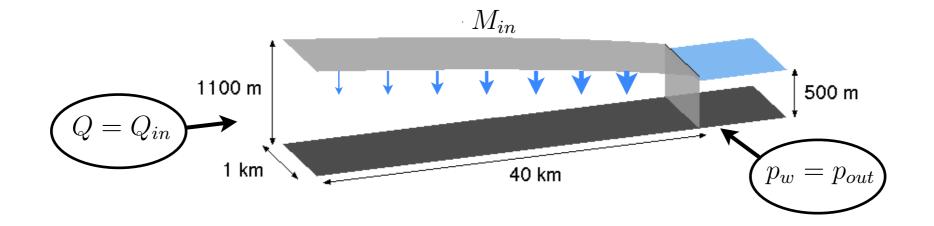
Neighbouring channels compete with one another

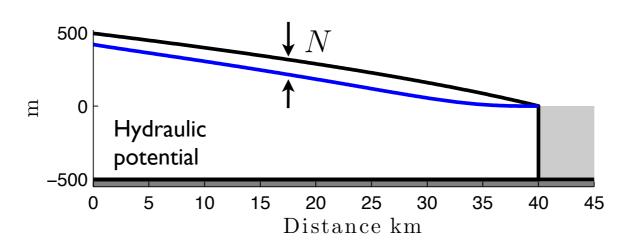


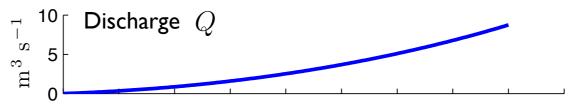
| leads to an arterial network



## Röthlisberger channels







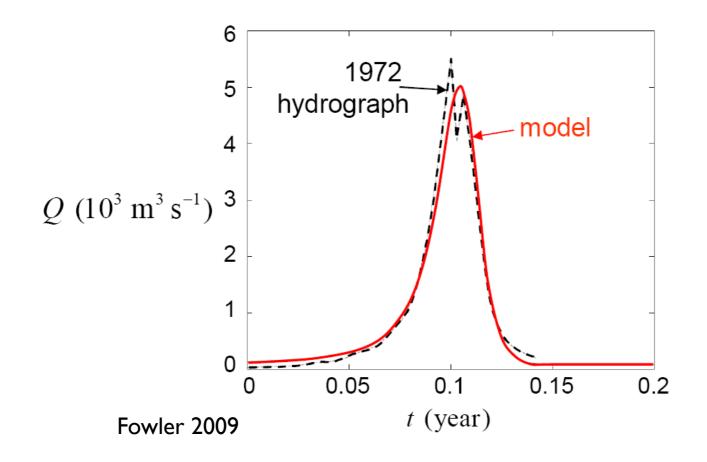
## Jökulhlaups (GLOFs) Nye 1976, Spring & Hutter 1981, Clarke 2003

A success of the channel theory is the application to floods from ice-dammed lakes

Combine channel evolution equation

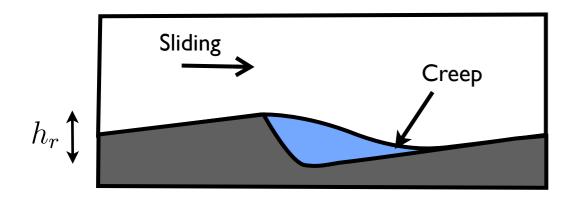
$$\frac{\partial S}{\partial t} = \frac{S^{4/3} \Psi^{3/2}}{\rho_i L} - \tilde{A} S N^n$$

with a **lake filling** equation 
$$-\frac{A_L}{\rho_w g}\frac{\partial N}{\partial t}=m_L-Q \qquad \text{at} \qquad x=0$$



## Linked cavities Walder 1986, Kamb 1987

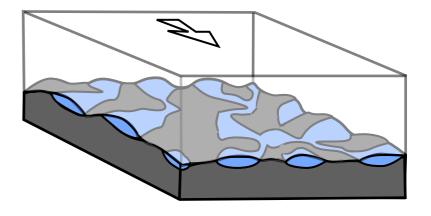
#### Cavities grow through sliding over bedrock

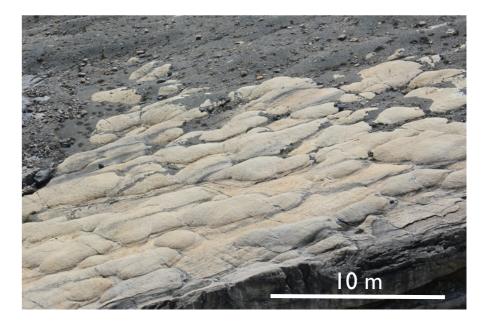


#### Model

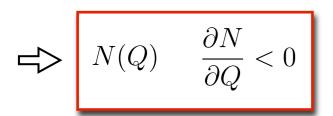
$$\frac{\partial \hat{S}}{\partial t} = U_b h_r - \tilde{A} \hat{S} N^n$$

Smaller 'orifices' control the flow





Approximate steady-state relationship





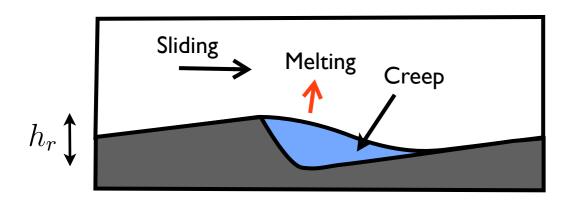
N(Q)  $\frac{\partial N}{\partial Q} < 0$   $\Longrightarrow$  Effective pressure DECREASES with discharge

Flow is distributed

Cavity size is controlled by parameter  $\Lambda = \frac{U_b}{N^n}$  i.e. depends on effective pressure **and** sliding speed

## Drainage system stability Walder 1986, Kamb 1987, Schoof 2010, Hewitt 2011

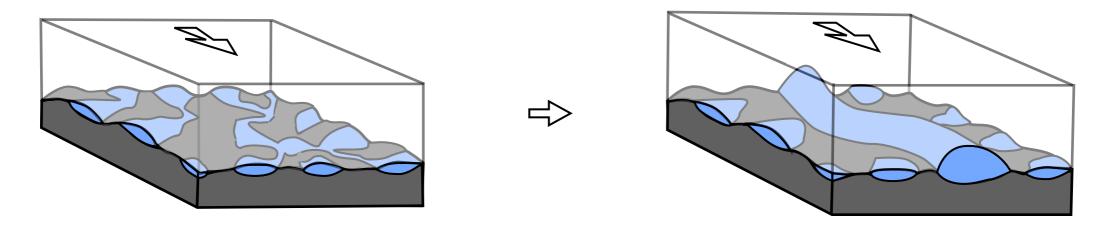
Energy is still dissipated by water flow



$$\frac{\partial S}{\partial t} = \frac{m}{\rho_i} + U_b h_r - \tilde{A}SN^n$$

A linked cavity system can become unstable to produce channels

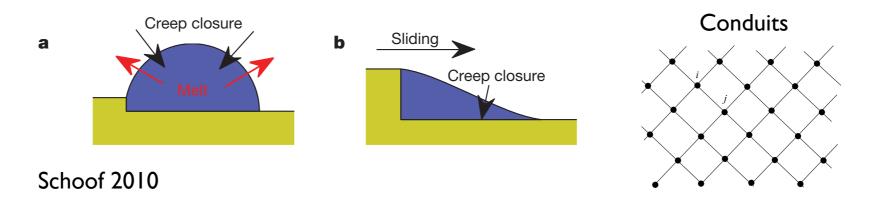
eg. if discharge becomes sufficiently large, or sliding speed sufficiently low

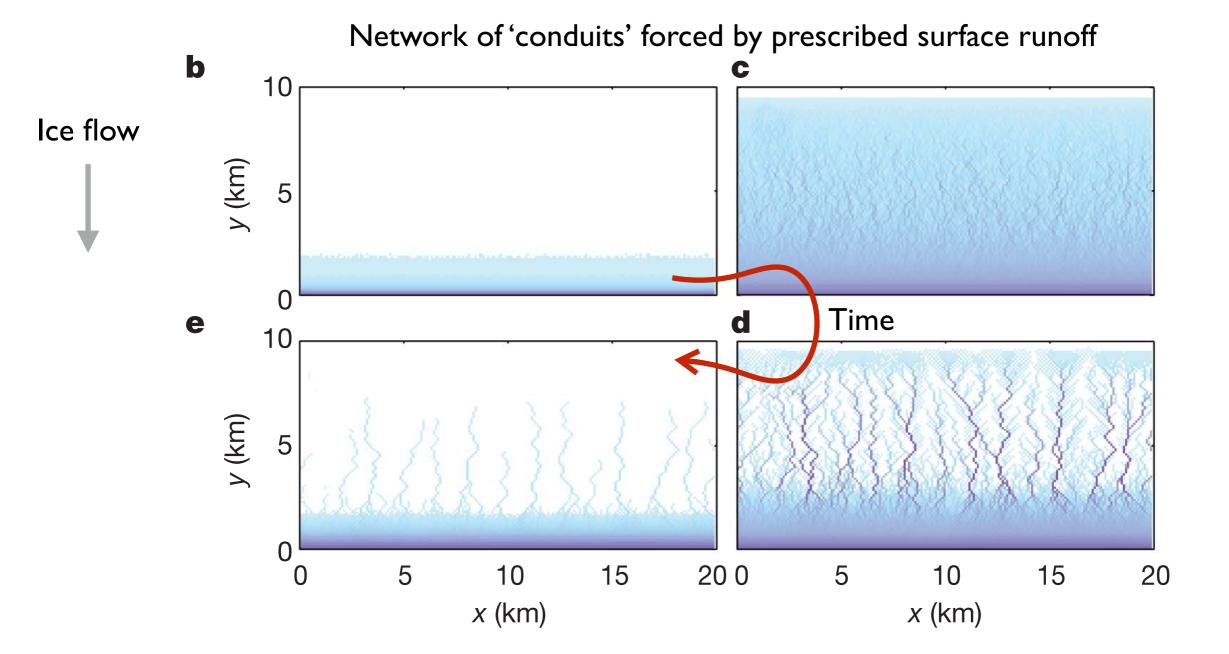


Conversely, a channel can become unstable to cavities

eg. if discharge low, or sliding speed sufficiently high

## Seasonal evolution of drainage system

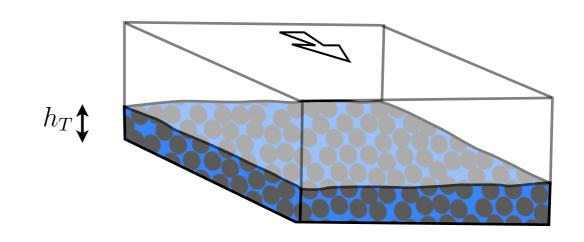


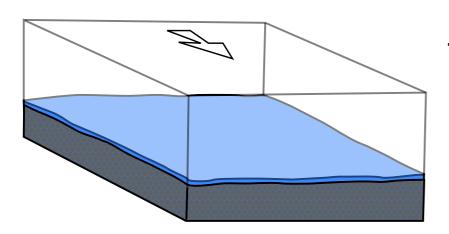


## Drainage through sediments

Water can infiltrate sediment layers and moves according to a diffusion equation

$$\beta_T h_T \frac{\partial \phi}{\partial t} = \nabla \cdot (K_T h_T \nabla \phi) + m$$
 
$$\uparrow$$
 Compressibility 
$$\uparrow$$
 Melting / freezing

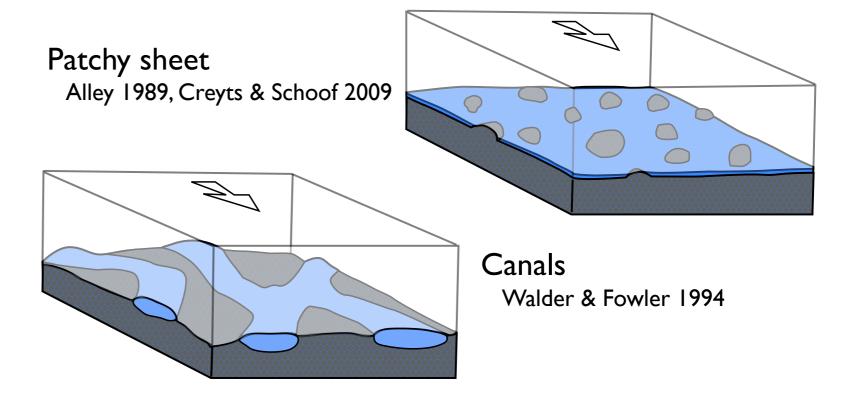




Till is relatively impermeable, so easily saturates

Film flow is unstable (Walder 1982)

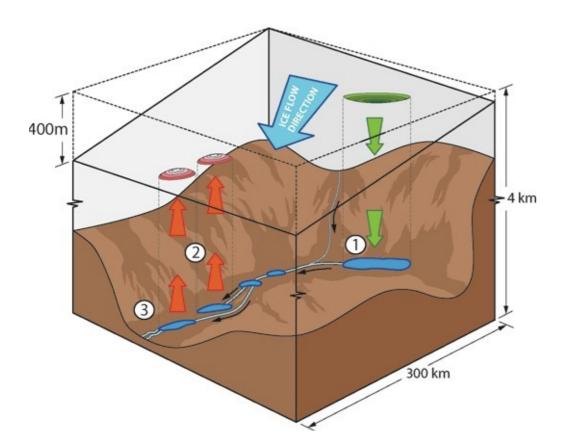
Locally deep water film leads to locally more melting

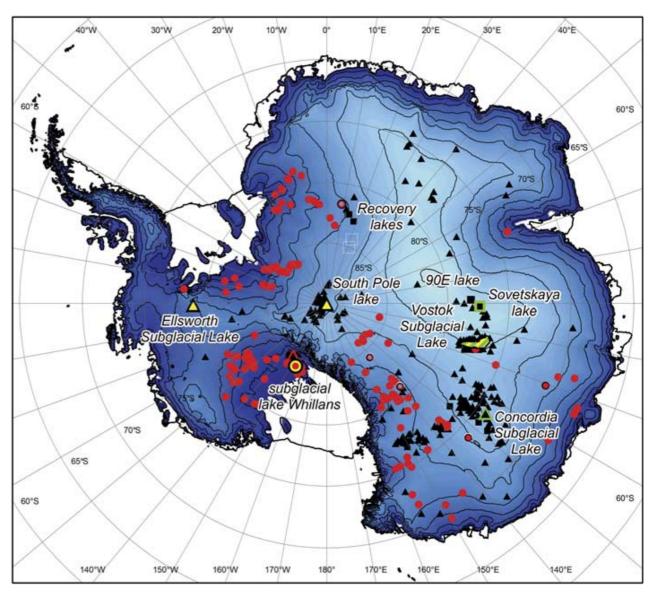


### Subglacial lakes Siegert 2005, Wingham et al 2006, Fricker et al 2007, Stearns et al 2008

Hundreds of lakes have been detected using radar and satellite observations.

'Active' lakes grow and drain quite frequently - through a jokulhlaup-like instability?





Wright & Siegert 2012

The formation and drainage of lakes may be important for ice-stream dynamics.

## Ice-sheet modelling

On a large scale, distributed systems are described as a 'sheet' flow

Average water depth h Average water pressure  $p_w$ 

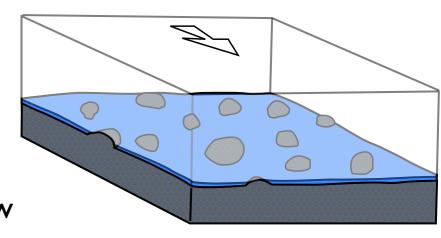
Average water flux

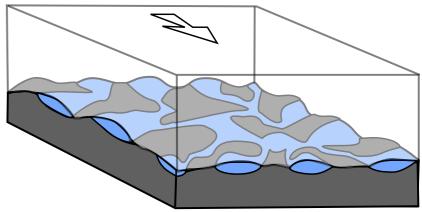
$$\mathbf{q} = -Kh^{\alpha}\nabla\phi$$

Mass conservation

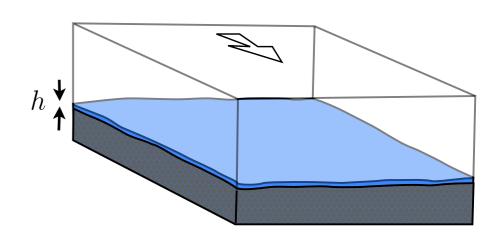
$$\frac{\partial h}{\partial t} + \nabla \cdot \mathbf{q} = \frac{m}{\rho_w} + M$$
 Basal melting Englacial/supraglacial source

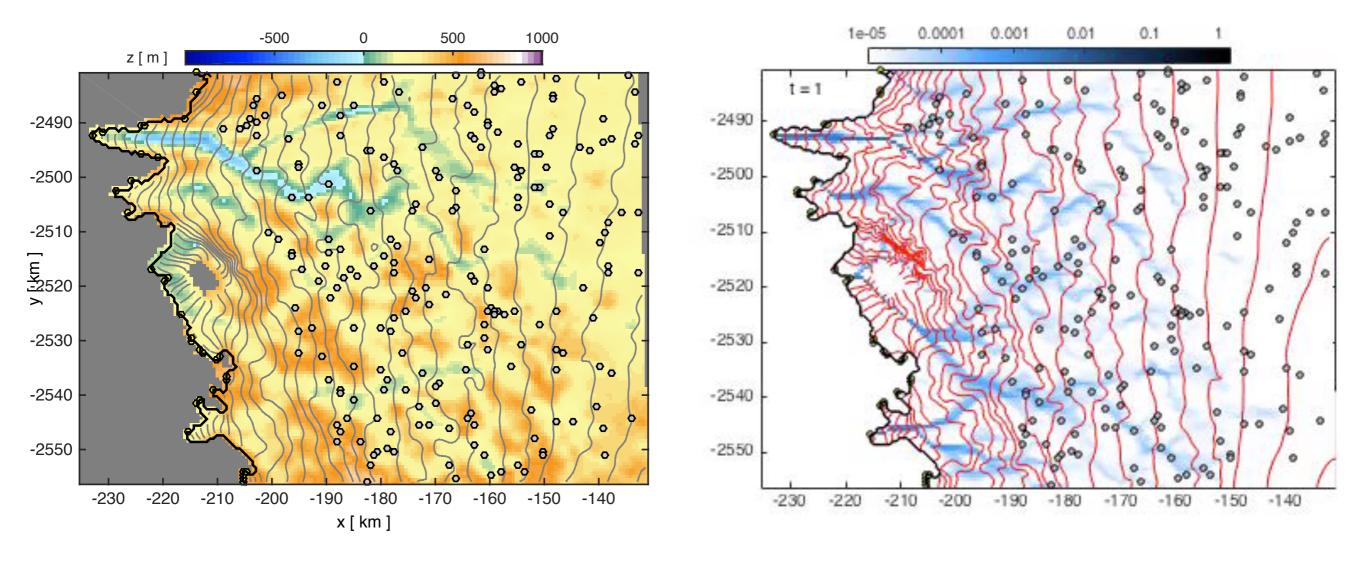
+ some additional ingredients to determine water pressure





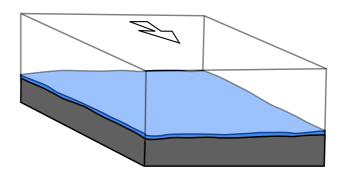




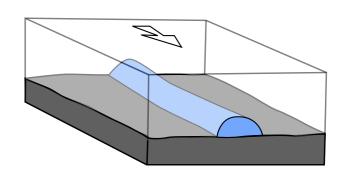


## Summary

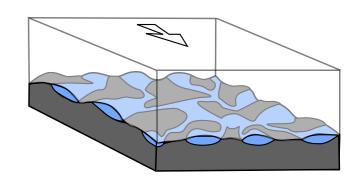
Uniform water film is unstable.



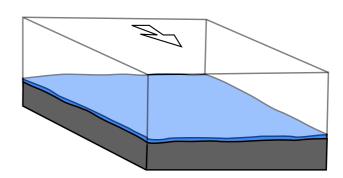
Röthlisberger channels form arterial networks.



Distributed flow in **linked cavities** or **patchy films** is possible.



On a large scale, the drainage system can be modelled as a **water layer** with variable thickness and permeability.



**Evolution of the drainage system** has important consequences for ice dynamics (surges, ice streams, seasonal/diurnal velocity changes)