

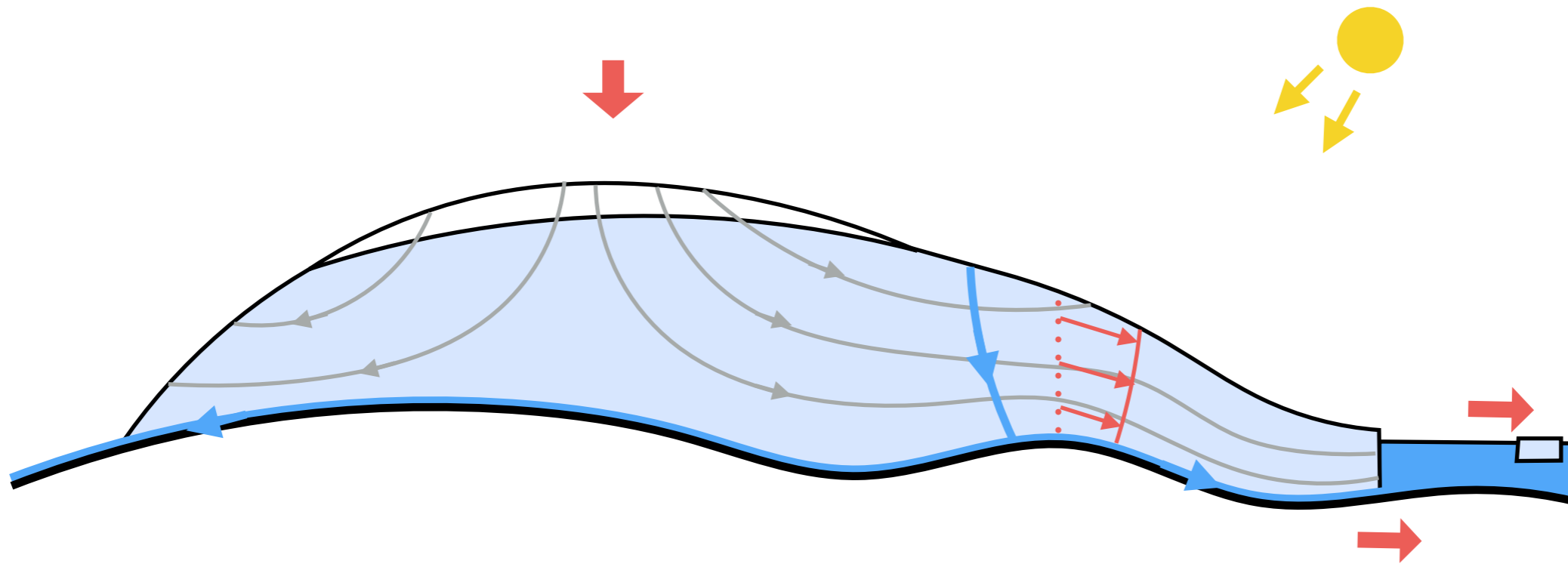
# Ice-sheet dynamics: the influence of glacier sliding on ice loss and sea level

Ian Hewitt, Mathematical Institute, University of Oxford





Greenland

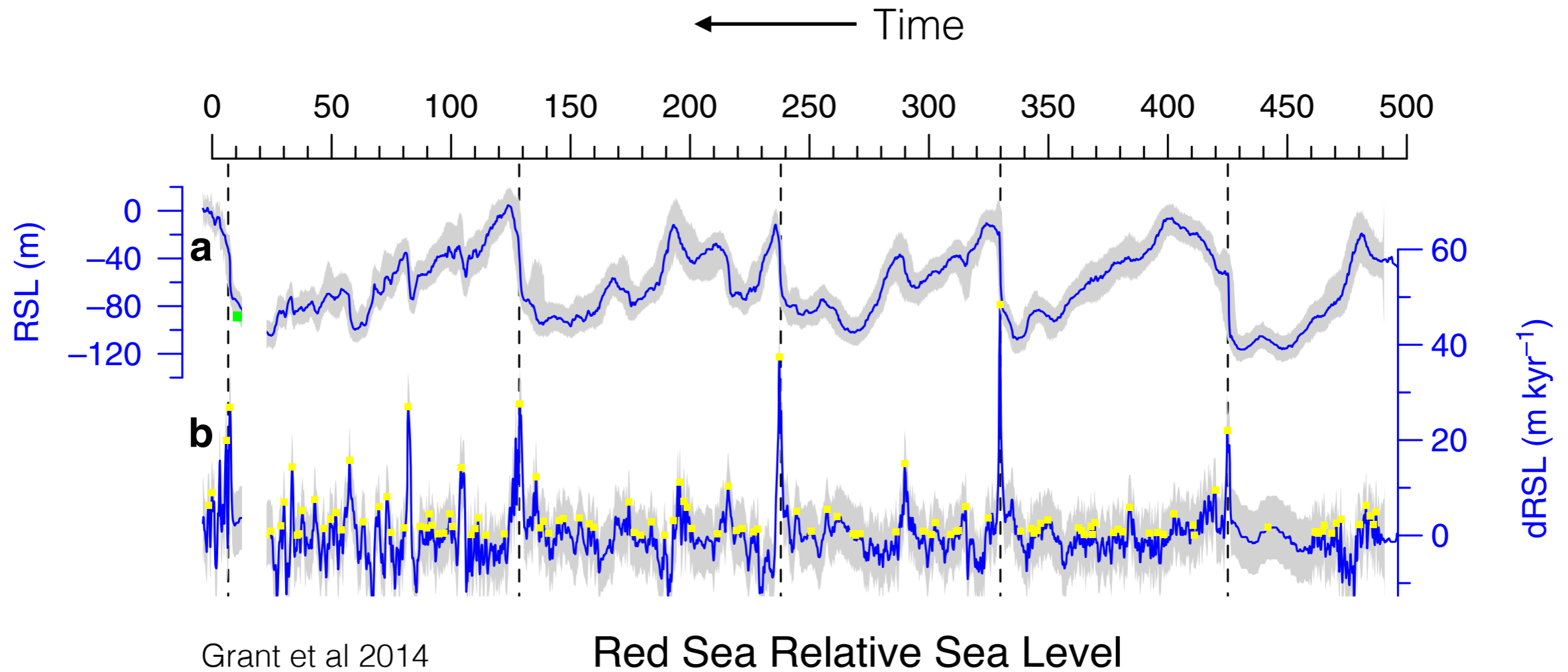


**How does meltwater penetrating to the bed affect ice-sheet motion?**

**What implications does this have for ice loss (sea level)?**

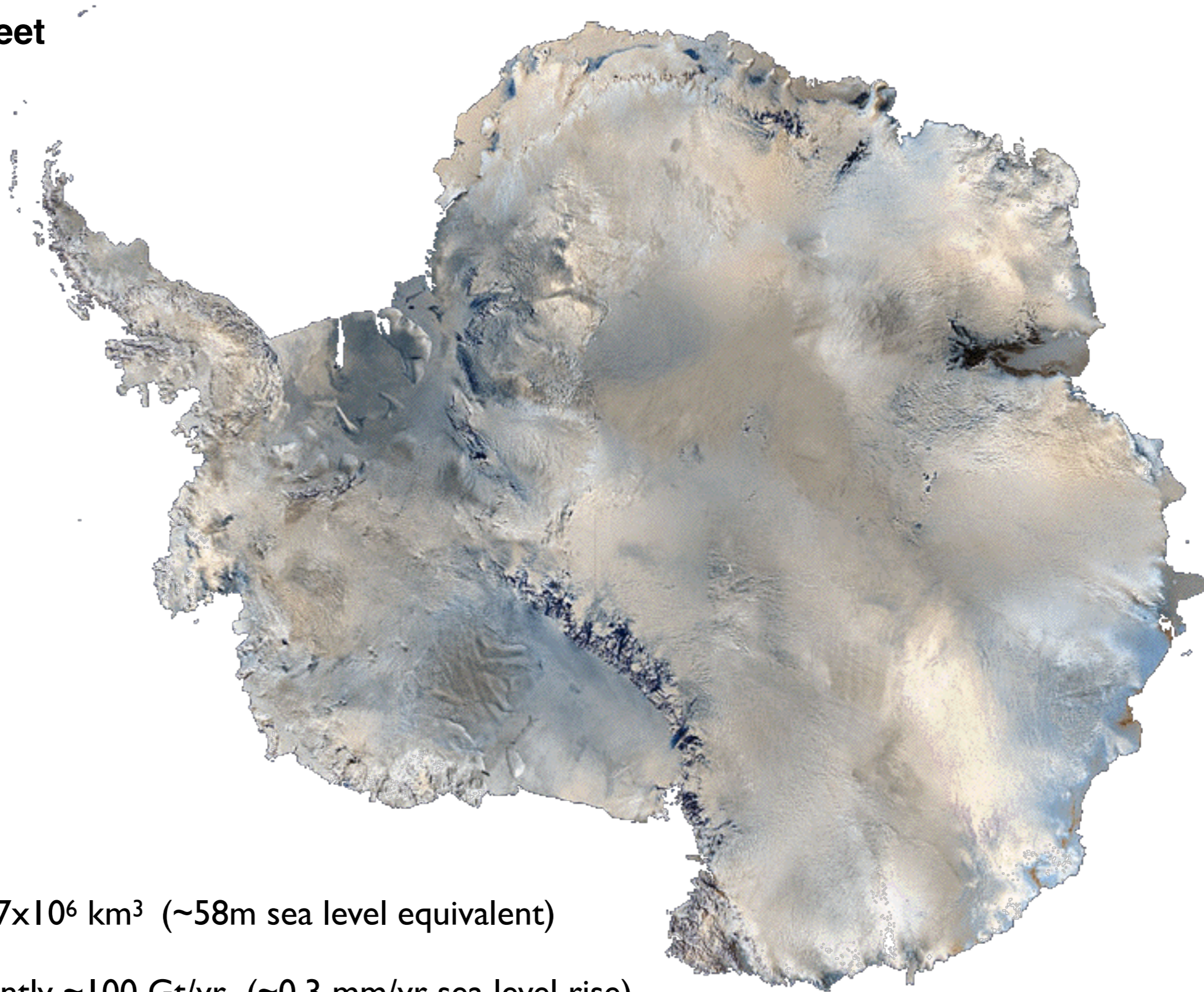
## Sea level history

The glacial period is punctuated by several periods of **rapid** sea level rise ( $\sim 1\text{m}/\text{century}$ )



Global sea level has been at least 6m higher in previous interglacials.

# Antarctic Ice Sheet



Current volume  $\sim 27 \times 10^6 \text{ km}^3$  ( $\sim 58\text{m}$  sea level equivalent)

Net mass loss currently  $\sim 100 \text{ Gt/yr}$  ( $\sim 0.3 \text{ mm/yr}$  sea level rise)

# Greenland Ice Sheet

Current volume  $\sim 2.7 \times 10^6 \text{ km}^3$  ( $\sim 7\text{m}$  sea level equivalent)

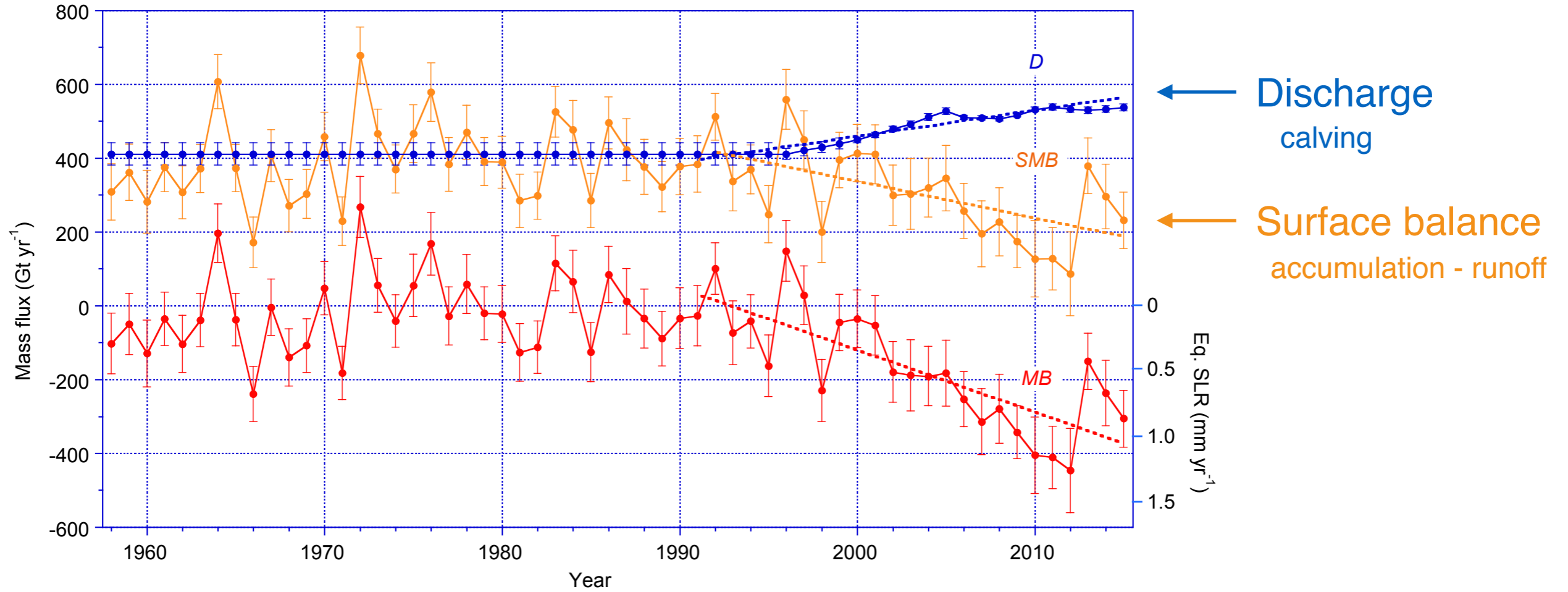
Net mass loss currently  $\sim 200 \text{ Gt/yr}$  (around  $0.6 \text{ mm/yr}$  sea level rise)

Timescale  $\sim 10,000$  years



Laura Stevens

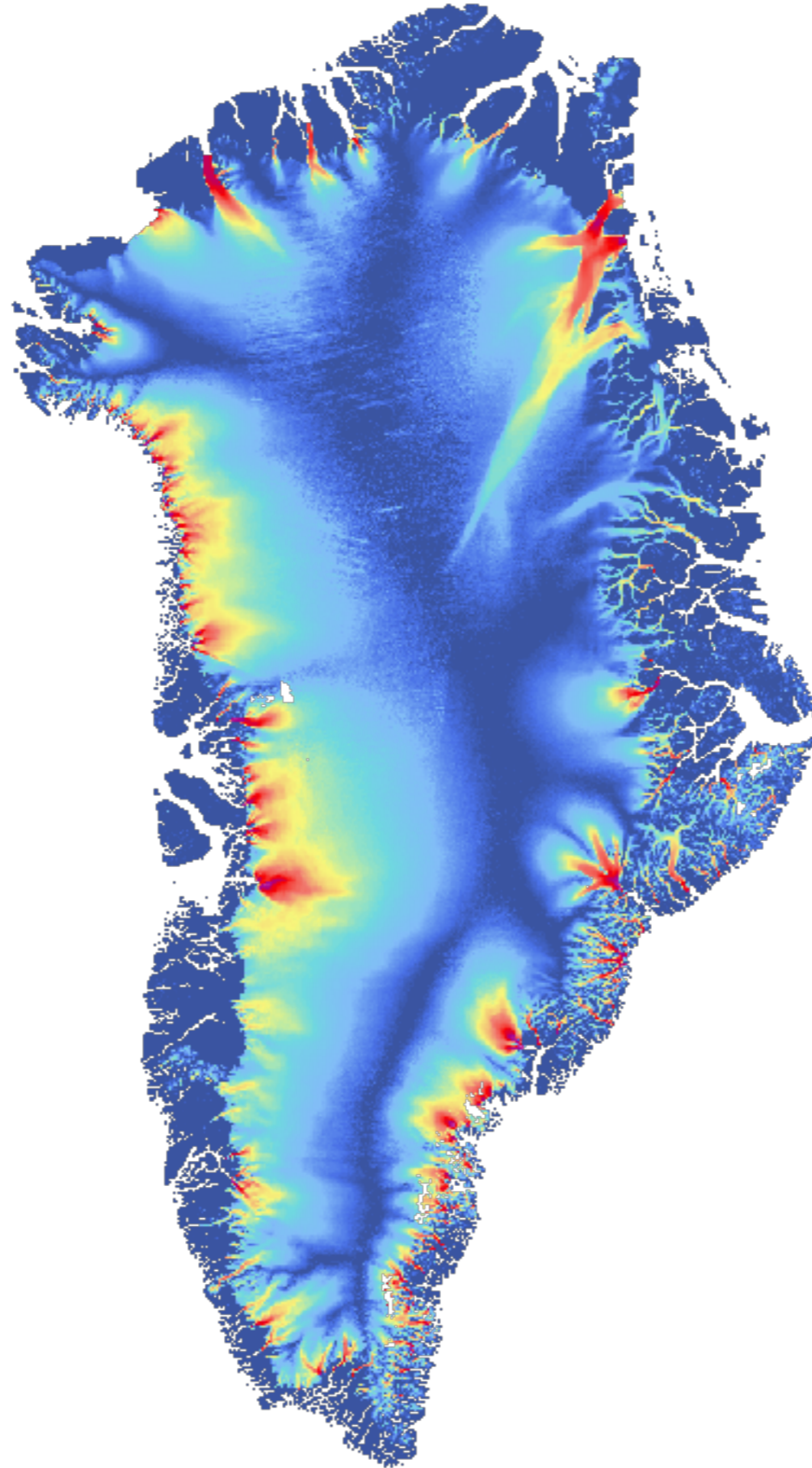
# Greenland ice sheet mass balance



van den Broeke et al 2016

Greenland is losing mass due to **decreased** SMB and **increased** discharge

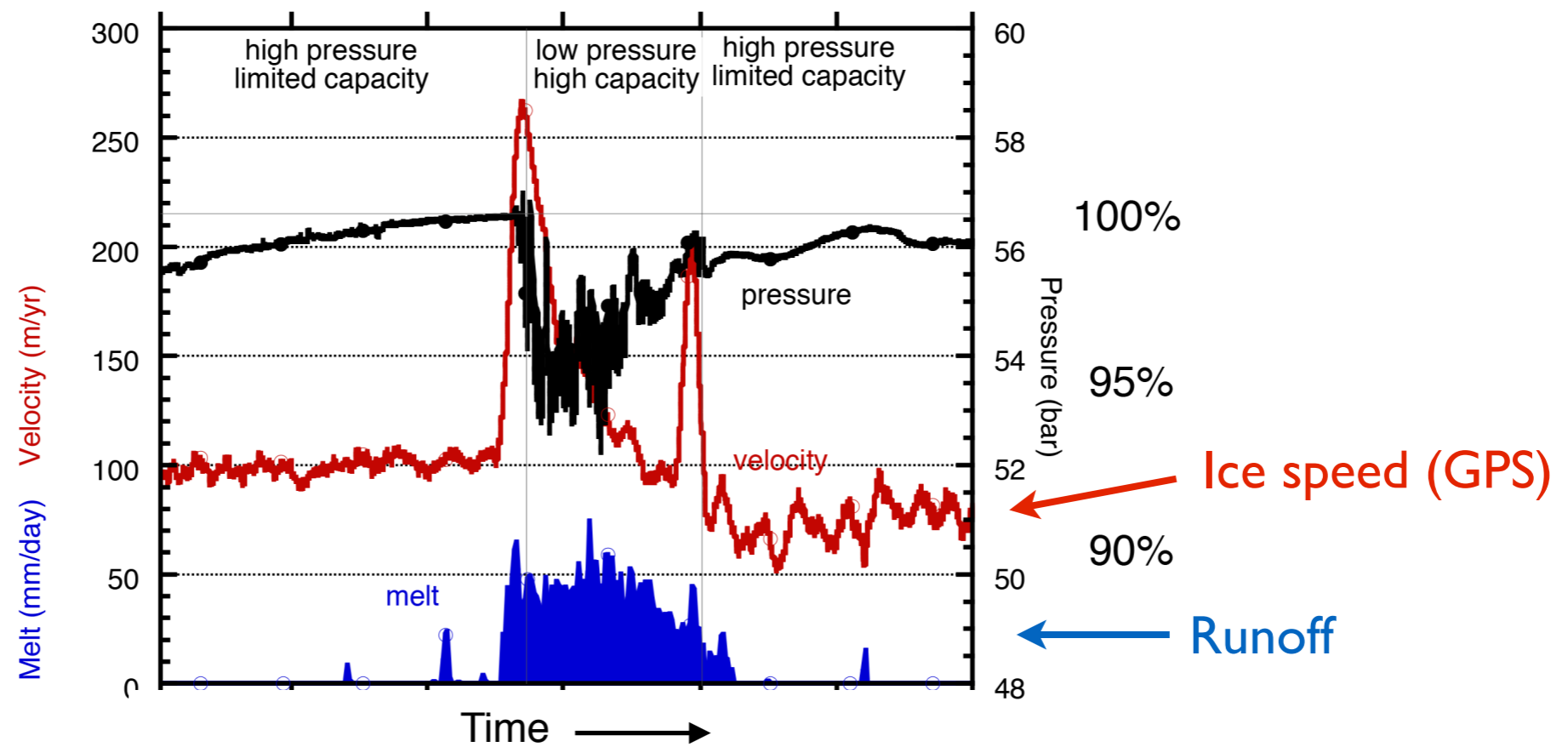
# Satellite-derived ice surface speeds



# Greenland ice sheet velocities

Summer drainage of surface meltwater to the bed causes large fluctuations in ice speed.

→ suggests potential for significant changes in ice velocity

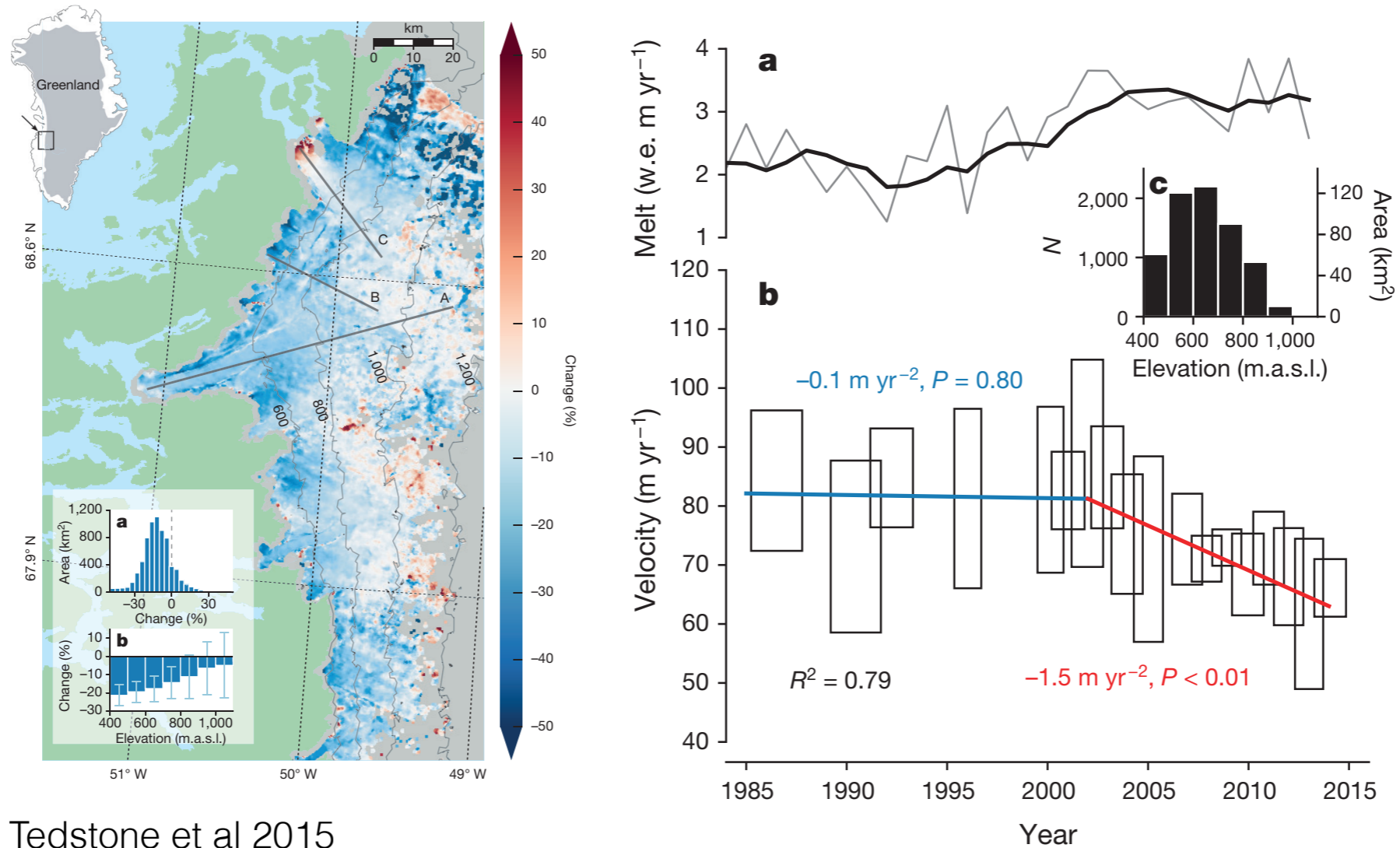


van de Wal et al 2015

# Greenland ice sheet velocities

Longer term measurements show a slight **decreasing** trend in **average** velocity, while runoff shows an increasing trend.

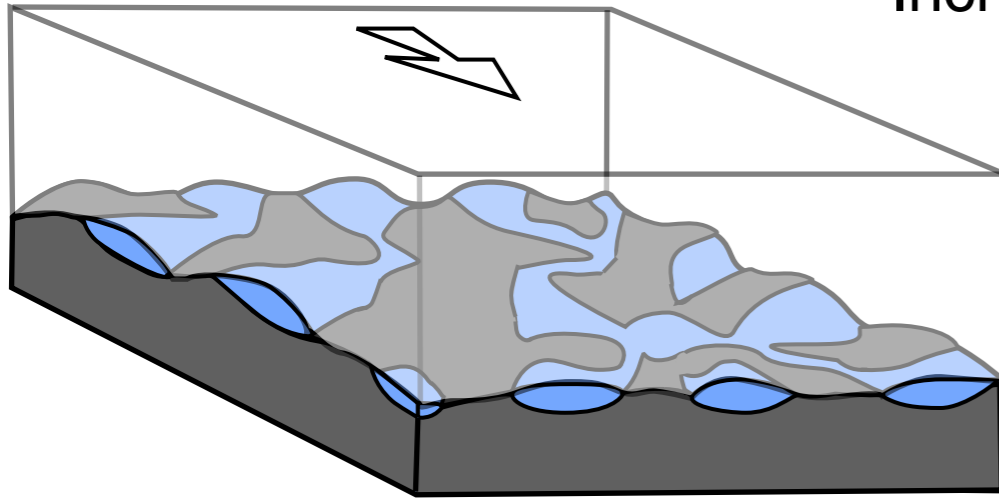
→ suggests possible negative relationship between runoff and average velocity?



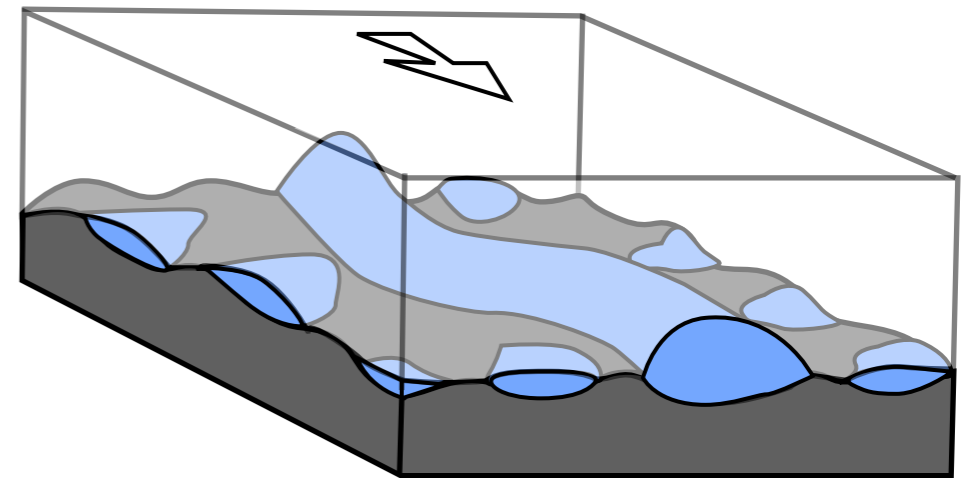
Tedstone et al 2015

# Evolution of the subglacial drainage system

Increased efficiency of drainage

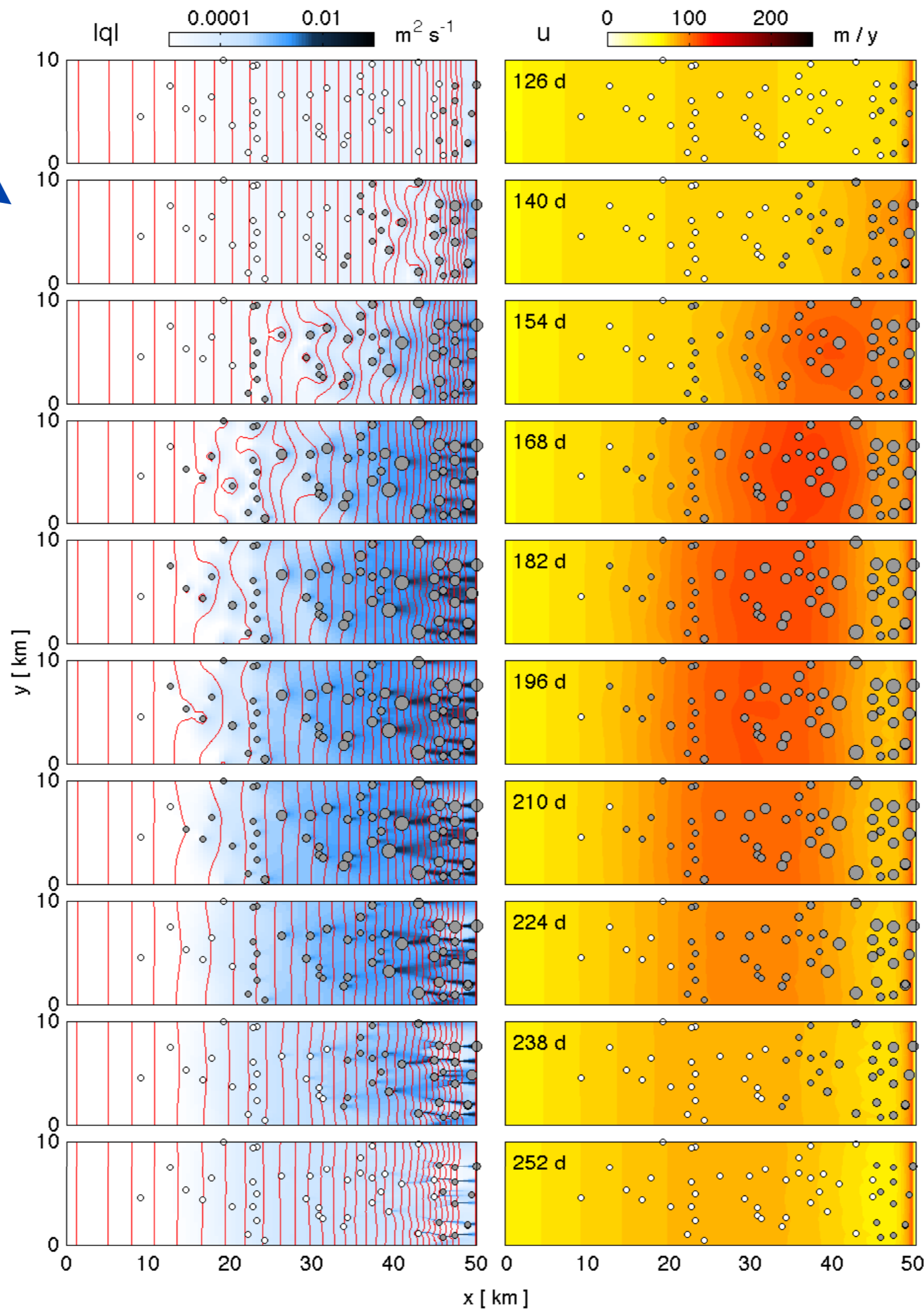


Isolated water pockets  
High average water pressure



Large melt-enlarged channels  
Lower average water pressure

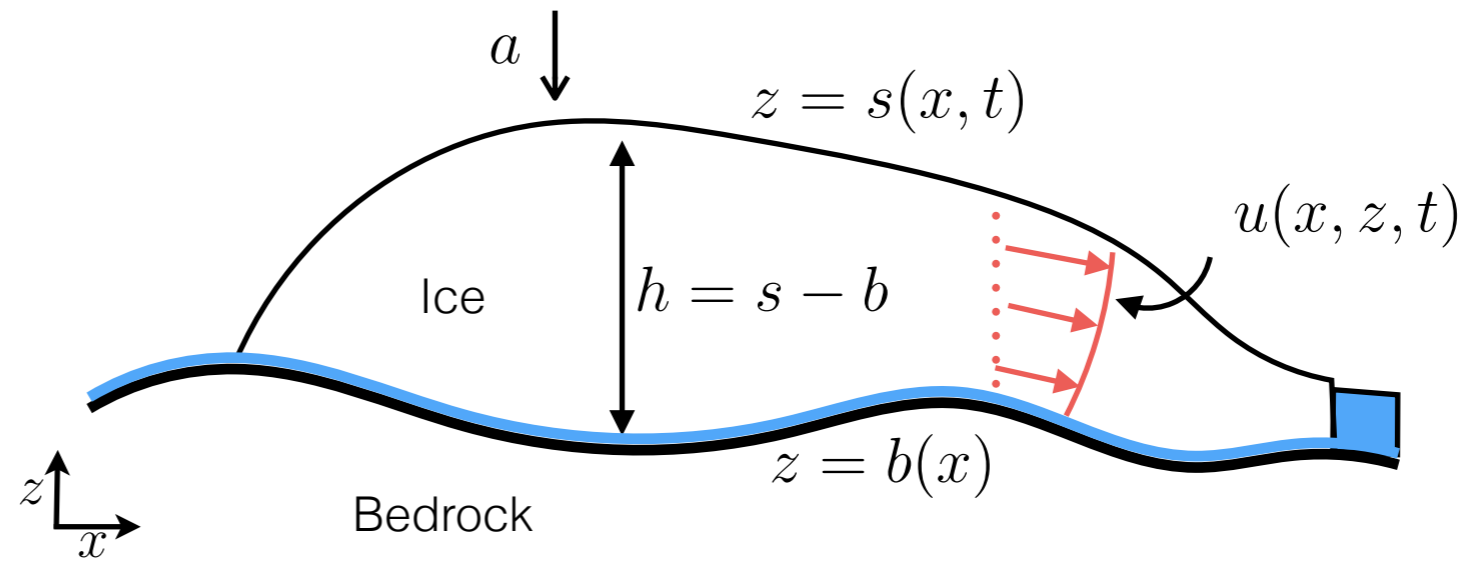
Subglacial discharge  
(areal  $m^2/s$ )



Ice speed

Time

# Mathematical model



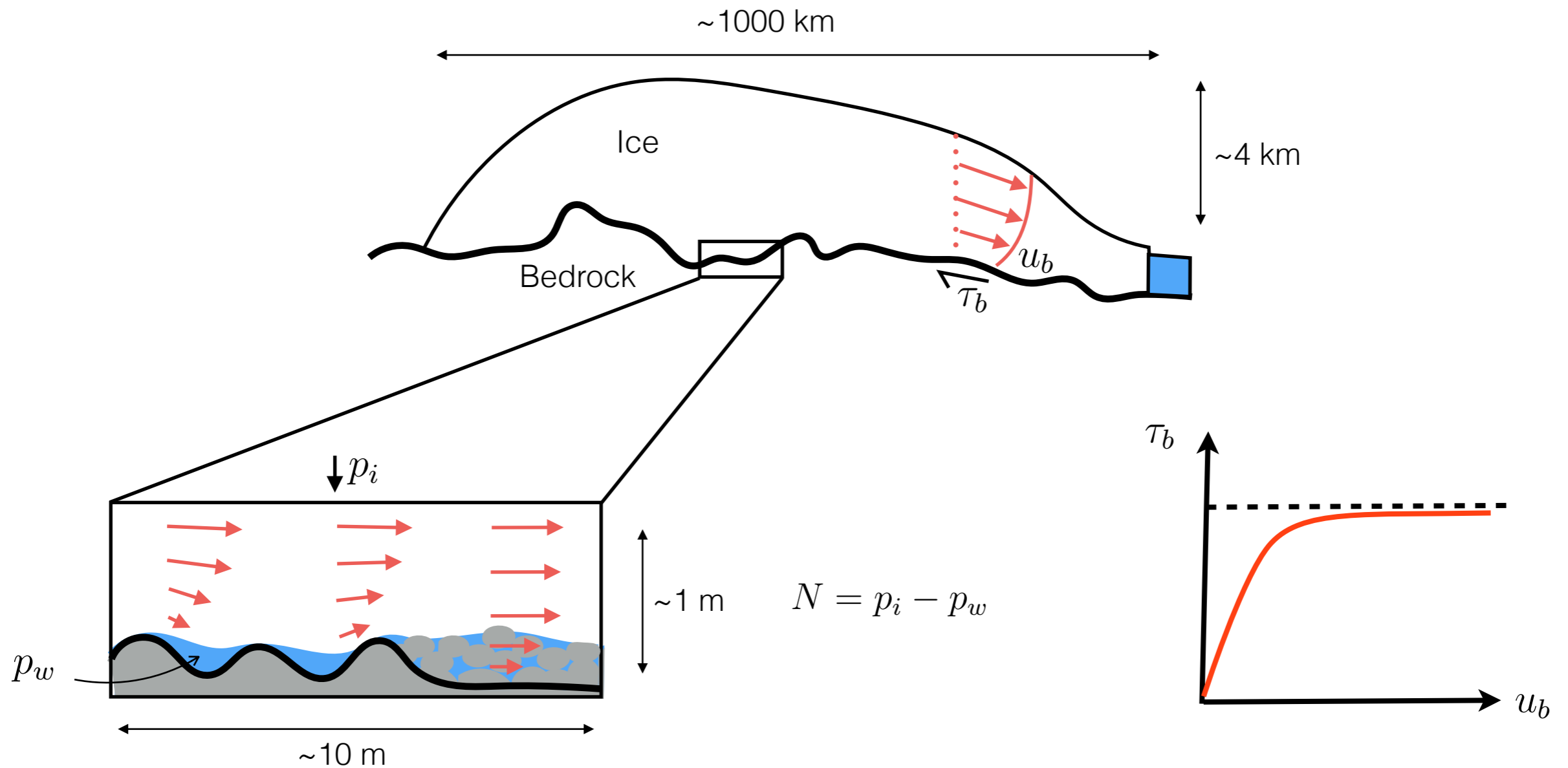
- Vertically-integrated mass conservation

$$\frac{\partial h}{\partial t} + \frac{\partial q}{\partial x} = a \quad q(x, t) = h\bar{u} = \int_b^s u \, dz \quad a = \text{net accumulation - melting}$$

- Force balance  $0 = \nabla \cdot \boldsymbol{\sigma} + \rho_i \mathbf{g}$

$$p_i = \rho_i g (s - z) \quad \tau_b = -\rho_i g h \frac{\partial s}{\partial x} + \frac{\partial}{\partial x} \left( 4h\eta_i \frac{\partial \bar{u}}{\partial x} \right) \quad \tau_b = f(\bar{u}, N) \quad N = p_i - p_w$$

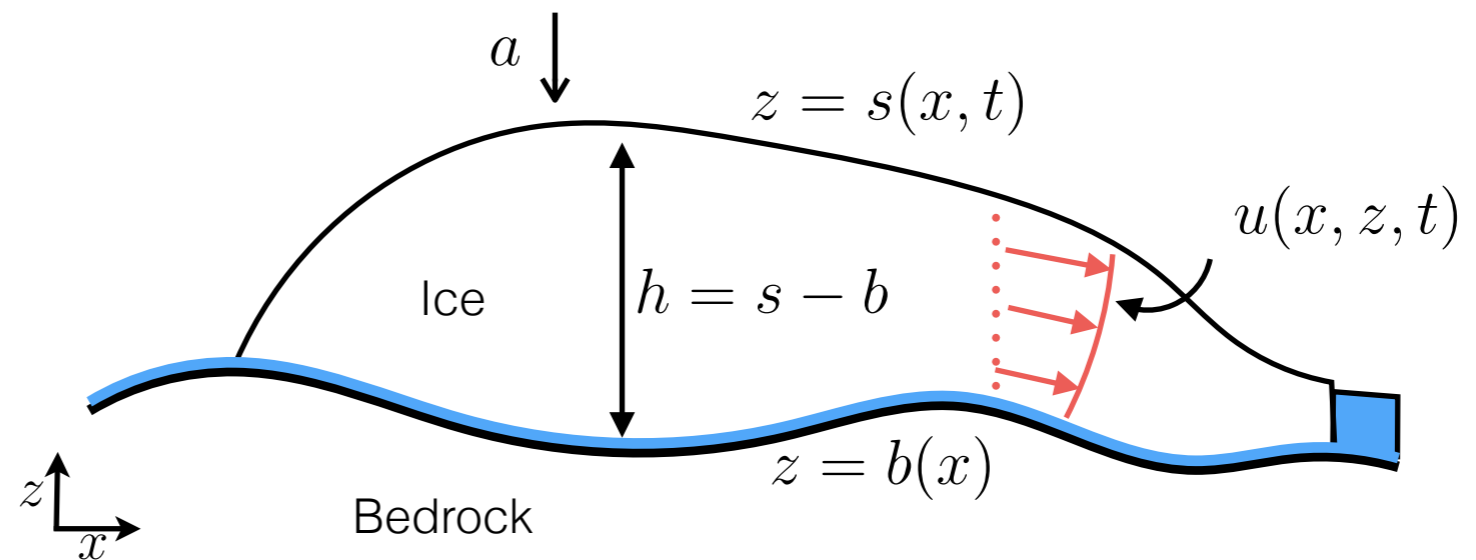
# Theoretical framework for modelling sliding



Theory and *some* measurements suggest a **friction / slip law** of the form

$$\tau_b = f(u_b, N) \quad \text{to be applied to the large-scale ice flow}$$

## Mathematical model



- Boundary conditions

land-terminating

$$h = 0, \quad q = 0 \quad \text{at} \quad x = x_m$$

marine-terminating

$$h\dot{x}_m = q - q_c, \quad 4h\eta_i \frac{\partial \bar{u}}{\partial x} = \frac{1}{2} (\rho_i g h^2 - \rho_o g b^2) \quad \text{at} \quad x = x_m$$

+ calving condition

$$h = f h_f \quad \text{at} \quad x = x_m$$

$$h_f = -\frac{\rho_o}{\rho_i} b \quad \text{flotation thickness}$$



# Glacier flow



© 2013 James Balog

Extreme Ice Survey - Time-lapse camera  
Khumbu glacier, Nepal

~10,000,000 x real time

# Glacier flow



© 2013 James Balog

Extreme Ice Survey - Time-lapse camera  
Khumbu glacier, Nepal

~10,000,000 x real time

# Marine-terminating / tidewater glaciers



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Extreme Ice Survey - Time-lapse camera  
Columbia Glacier, Alaska

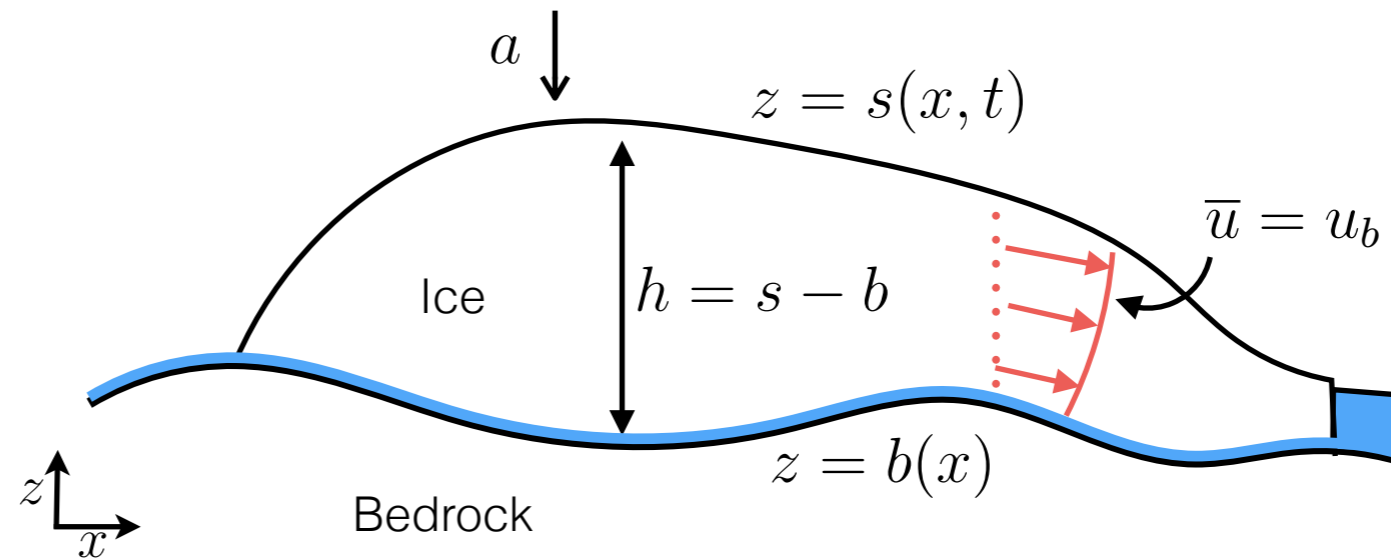
# Marine-terminating / tidewater glaciers



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Extreme Ice Survey - Time-lapse camera  
Columbia Glacier, Alaska

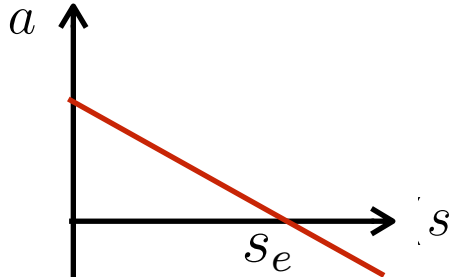
## Mathematical model



- If friction law is invertible,  $\bar{u} = F(\tau_b, N)$ , and bed topography relatively flat, the problem is seemingly **diffusive**

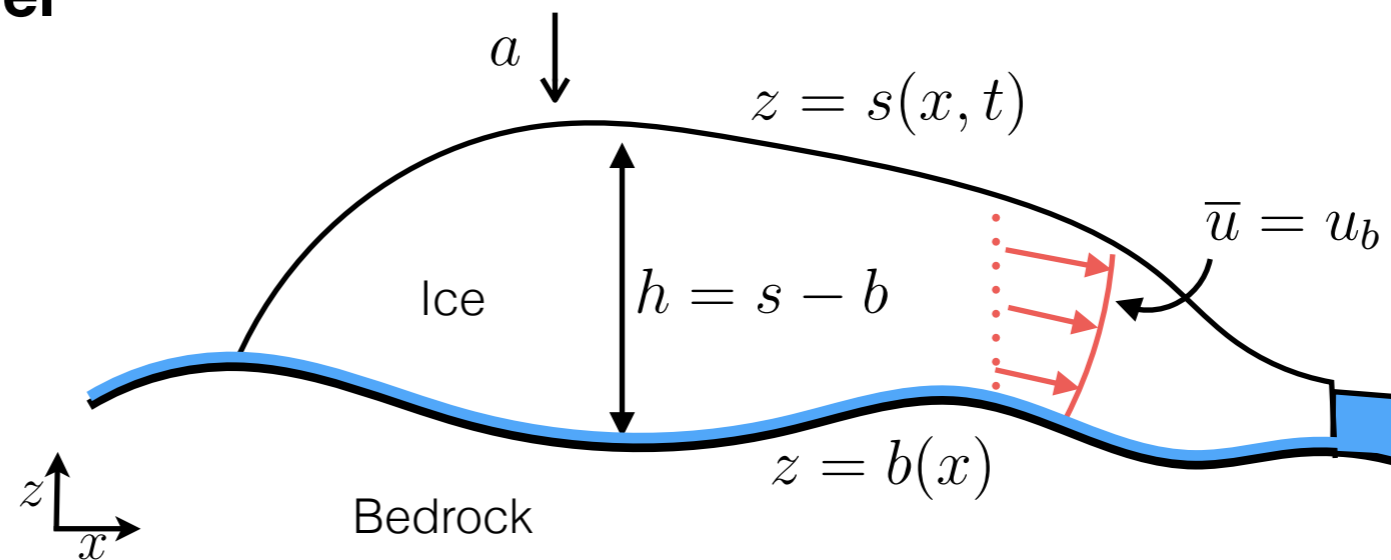
$$\frac{\partial h}{\partial t} = \frac{\partial}{\partial x} \left( K \frac{\partial h}{\partial x} \right) + a \quad K = \rho_i g h F_\tau$$

However, the diffusion coefficient may be highly non-linear, and accumulation rate varies with ice thickness.

- e.g.  $a = \lambda(s - s_e)$    $s_e$  equilibrium line altitude (ELA)

→ generic behaviour is 'blow-up' (cf. reaction-diffusion problems)

## A reduced 'plastic bed' model



- Vertically-integrated mass conservation

$$\frac{\partial h}{\partial t} + \frac{\partial q}{\partial x} = a \quad q = h\bar{u}$$

$$a = \lambda(s - s_e) \quad \text{net accumulation - melting}$$

$s_e$  equilibrium line altitude (ELA)

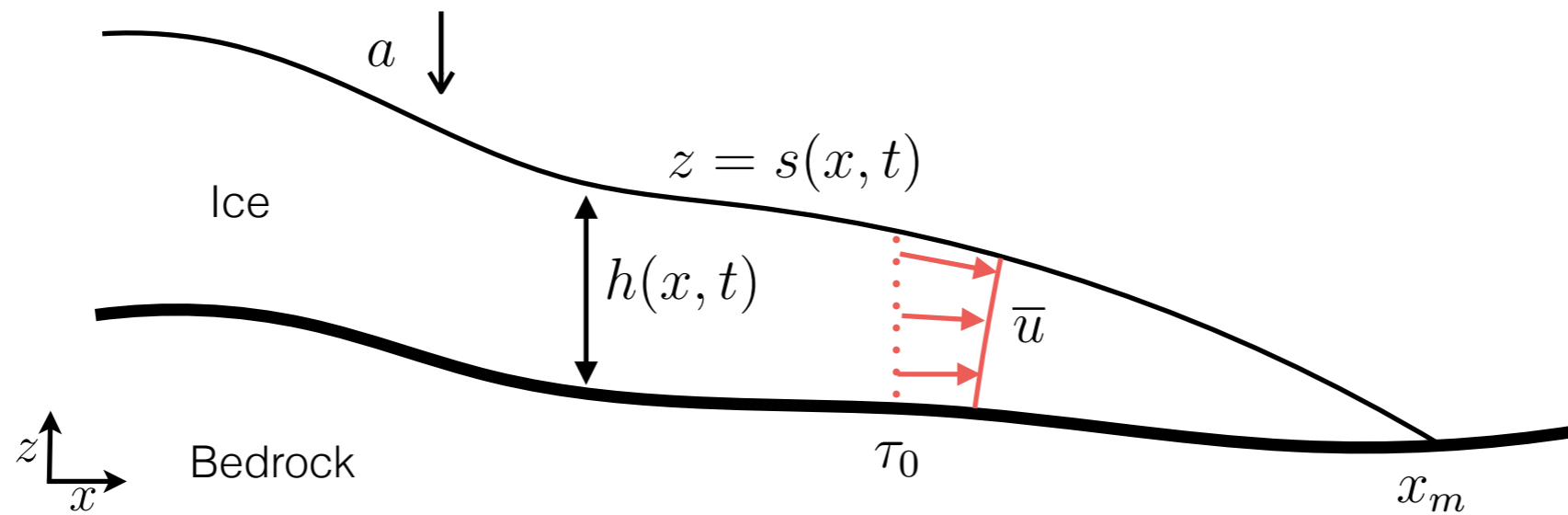
- Force balance + friction parameterisation

$$\tau_b = -\rho_i g h \frac{\partial s}{\partial x} + \frac{\partial}{\partial x} \left( 4h\eta_i \frac{\partial \bar{u}}{\partial x} \right)$$

$$\tau_b = \mu \langle N \rangle = \tau_0 \quad \text{bed 'strength'}$$

**Goal:** Consider effect of a long-term changes in  $\tau_0$  and  $s_e$

# Land-terminating glacier



- Boundary conditions  $h = 0, \quad q = 0$  at  $x = x_m(t)$

**Force balance**  $\tau_0 = -\rho_i g h \frac{\partial s}{\partial x} \rightarrow$  ice thickness & volume  $V = \int_0^{x_m} h dx$

e.g. for a flat bed profile  $h = \sqrt{\frac{2\tau_0}{\rho_i g}} (x_m - x)^{1/2}$

**Mass conservation**

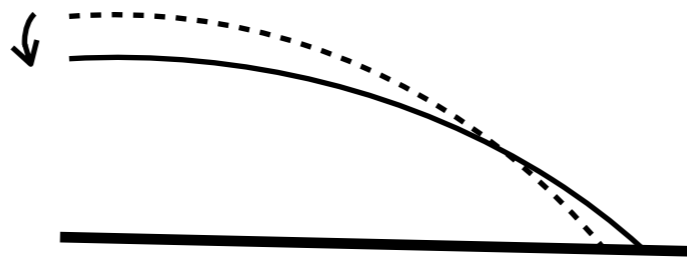
$$\rightarrow \frac{dV}{dt} = \int_0^{x_m} a dx$$

an ODE for the evolution of ice volume

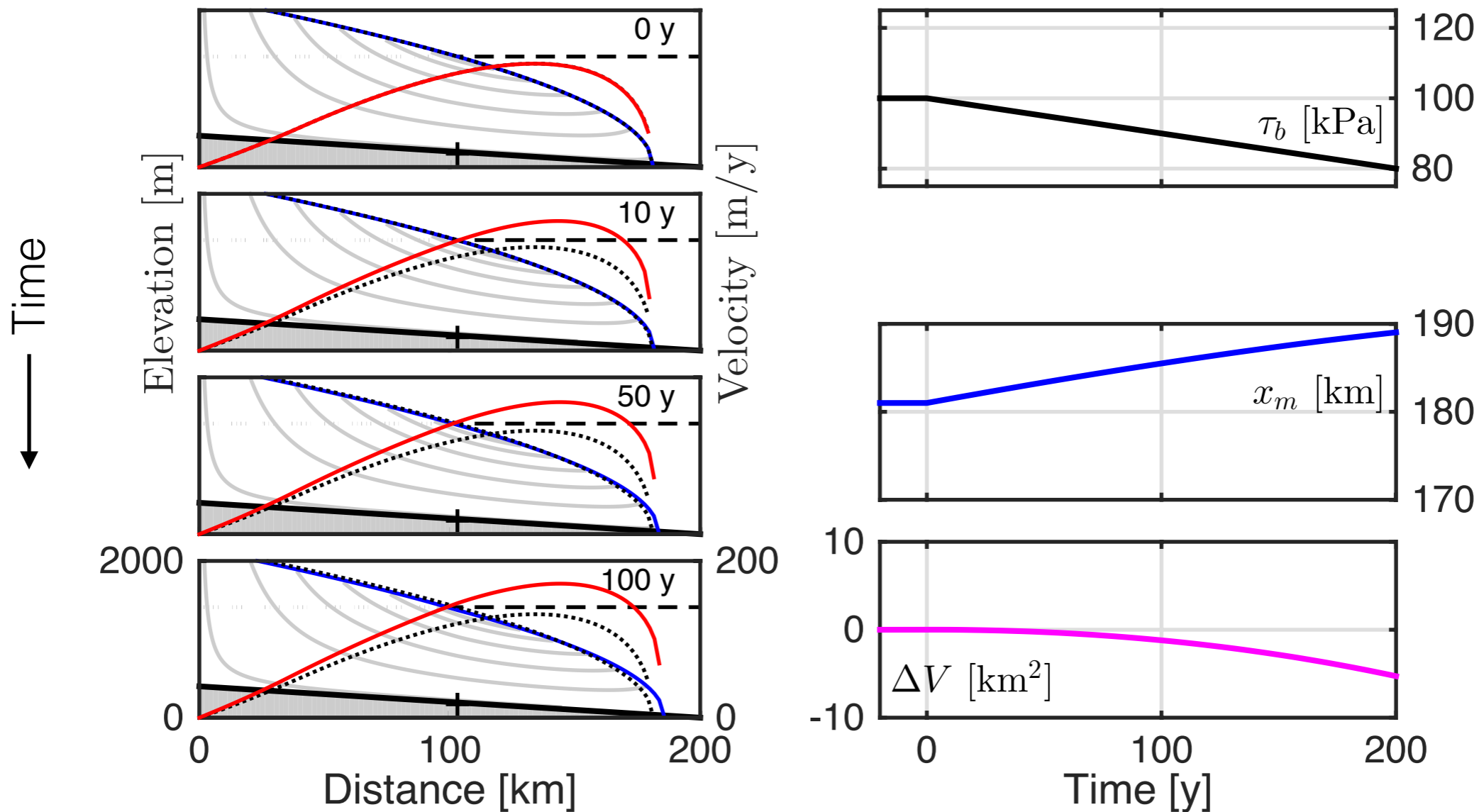
e.g.  $\frac{dV}{dt} = \lambda \left[ V - V^{2/3} \left( \frac{9\rho_i g}{8\tau_0} \right)^{1/3} s_e \right]$

# Land terminating glacier

A gradual **decrease** in bed strength results in **increased velocities** and **mass loss**

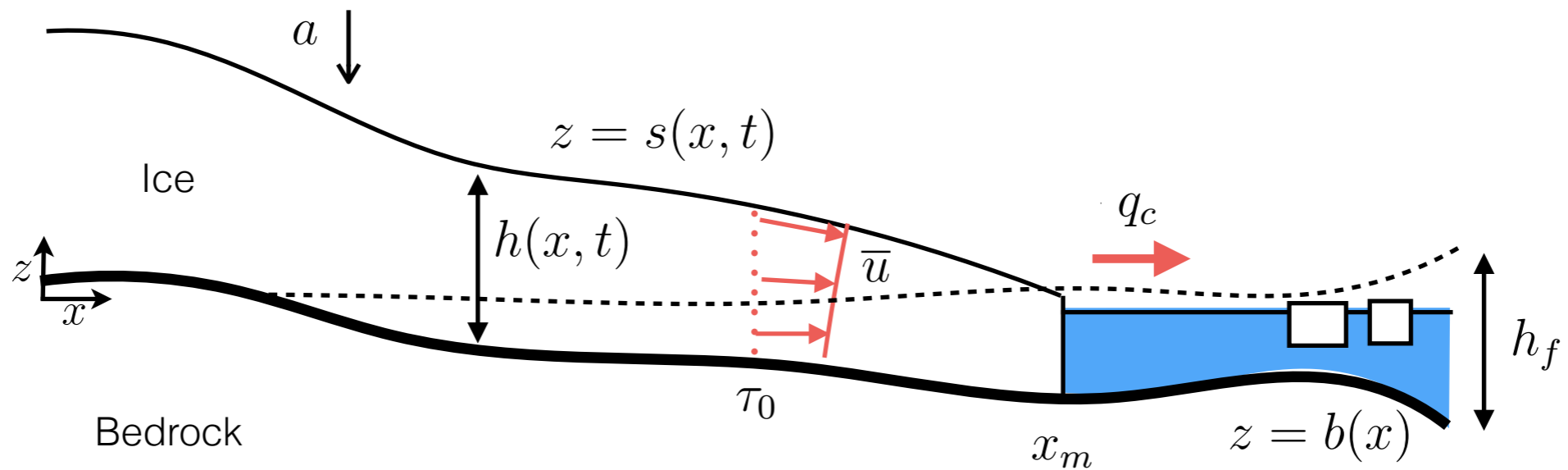


(an increase in bed strength causes the opposite)



Rate of ice loss controlled by SMB

# Marine-terminating glacier



- Boundary conditions  $h = fh_f$ ,  $h\dot{x}_m = q - q_c$ ,  $4h\eta_i \frac{\partial \bar{u}}{\partial x} = \frac{\rho_i g}{2} \left( h - \frac{\rho_i}{\rho_o} h_f^2 \right)$  at  $x = x_m(t)$

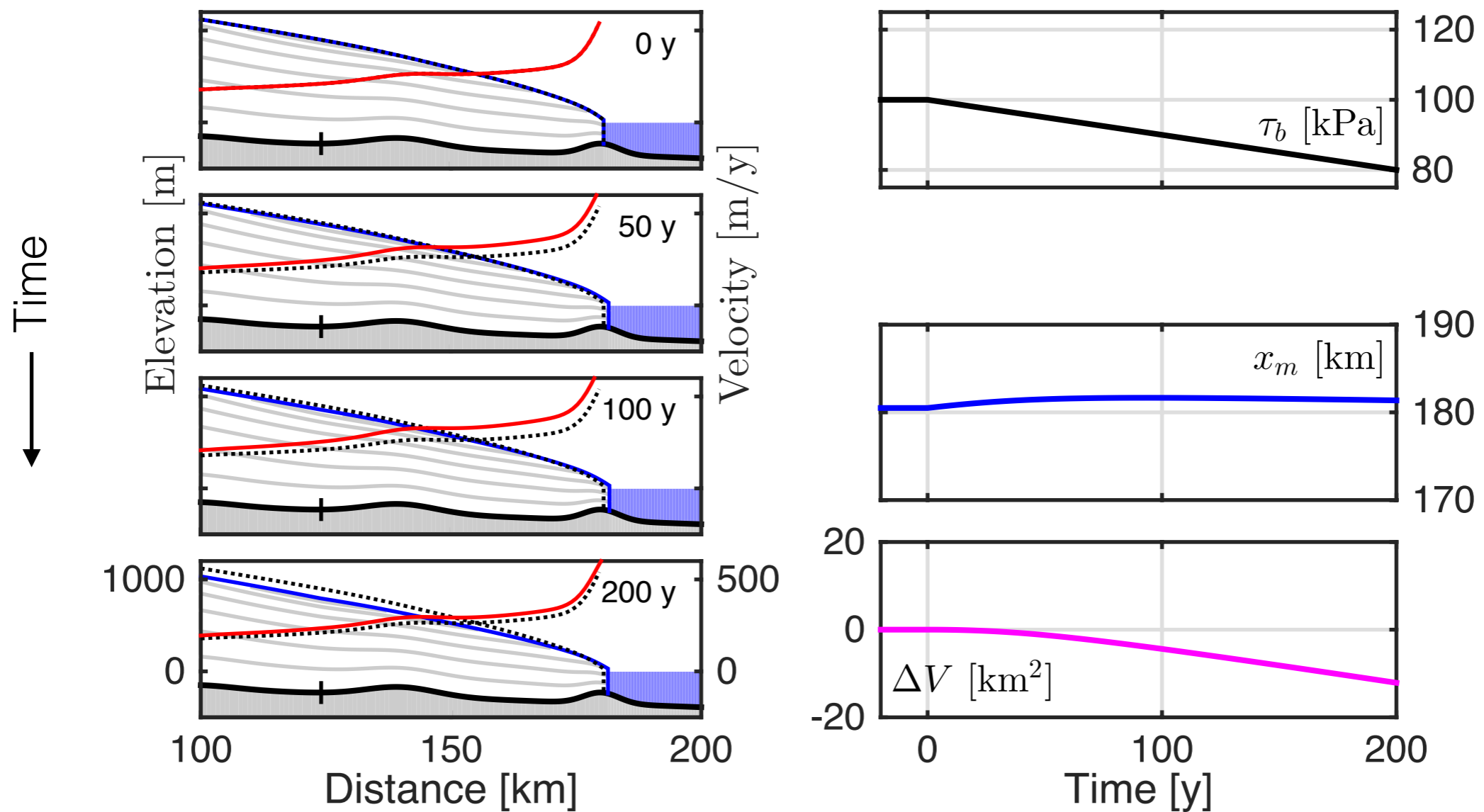
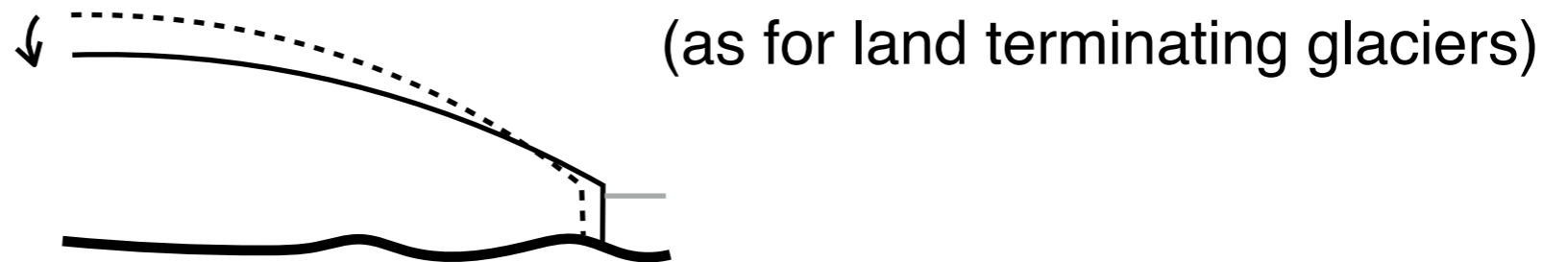
**Force balance**  $\tau_0 = -\rho_i g h \frac{\partial s}{\partial x} \rightarrow$  thickness + volume  $V = \int_0^{x_m} h dx$

**Mass conservation**  $\rightarrow \frac{dV}{dt} = \int_0^{x_m} a dx - q_c$  Calving flux  $q_c = \frac{\rho_i g}{\eta_i \mu} \hat{Q}(f) h_f^3$

an ODE for the evolution of ice volume

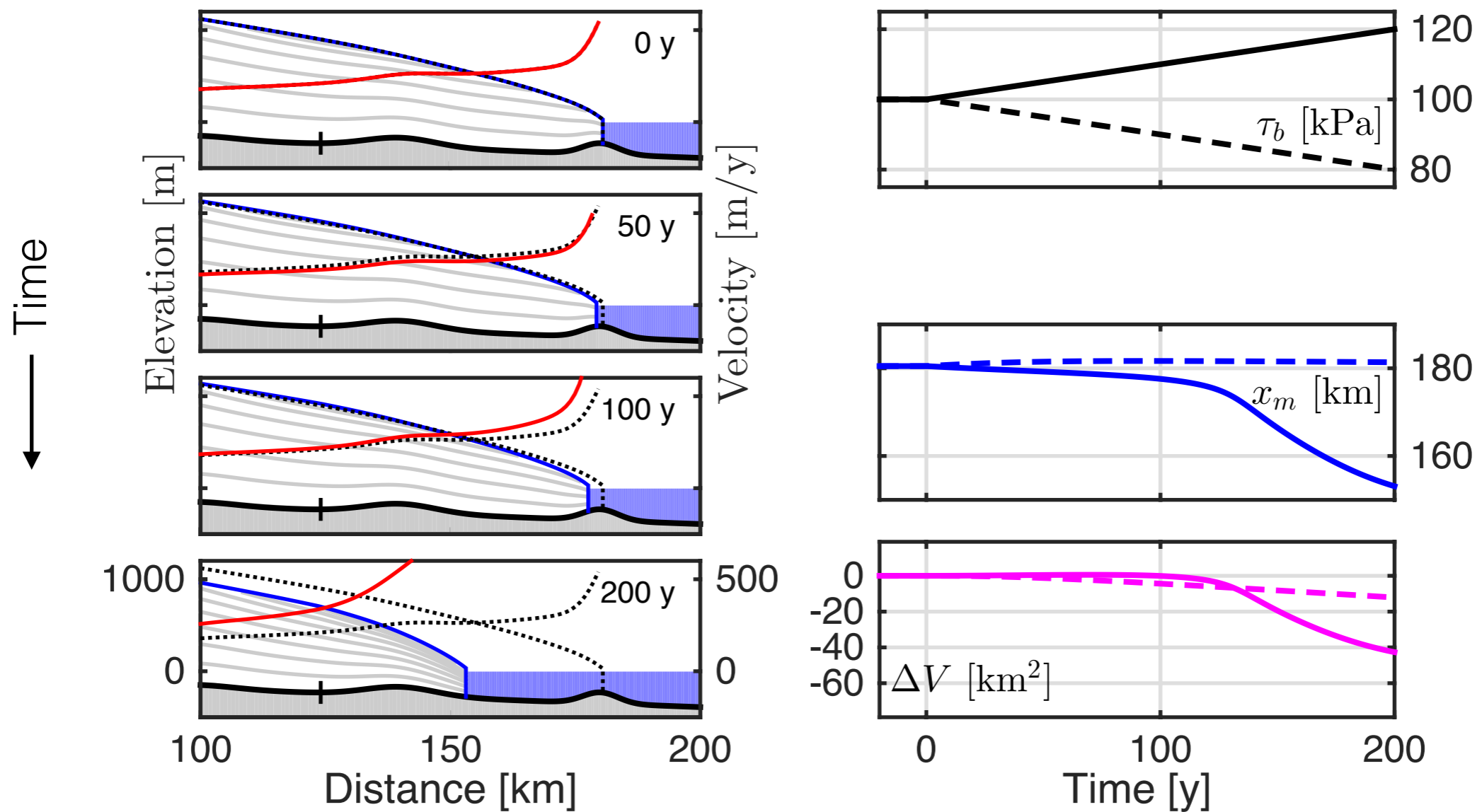
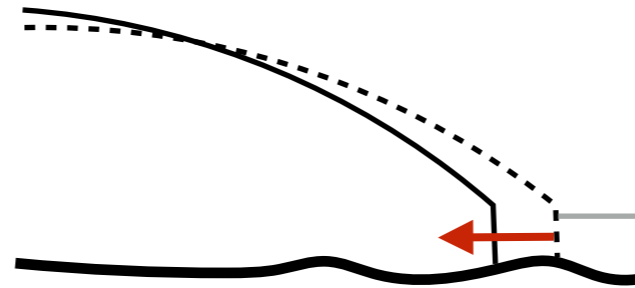
# Marine terminating glacier

A gradual **decrease** in bed strength results in **increased velocities** and **mass loss**



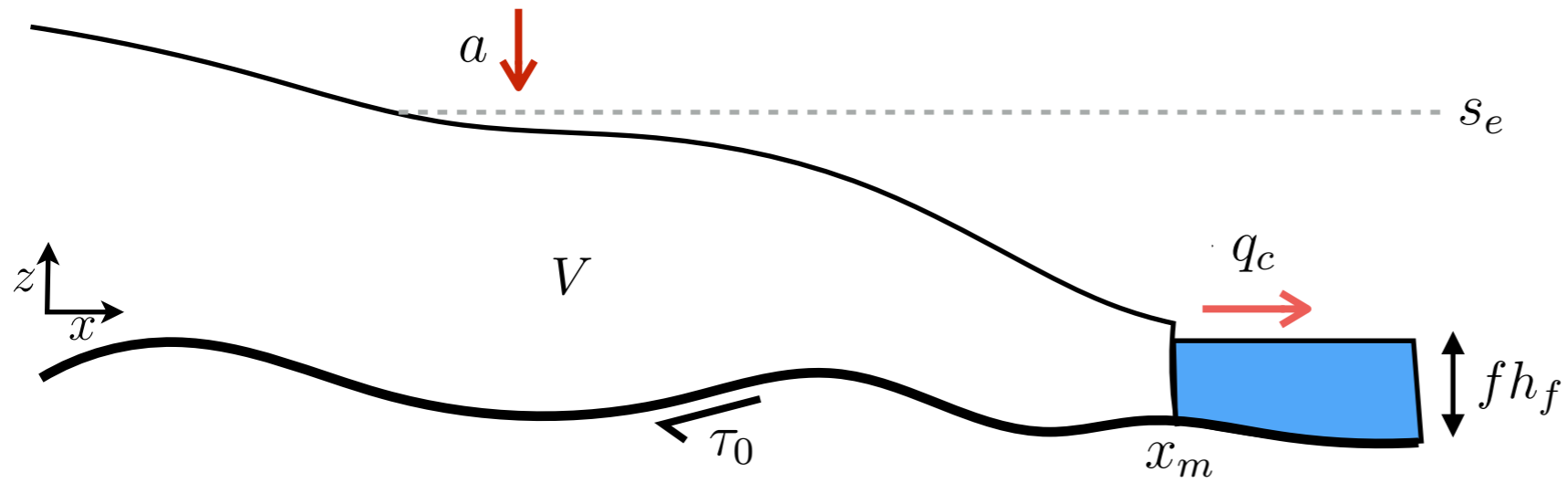
# Marine terminating glacier

An **increase** in bed strength results in **initially decreased velocities** ... but this initiates terminus retreat and acceleration.



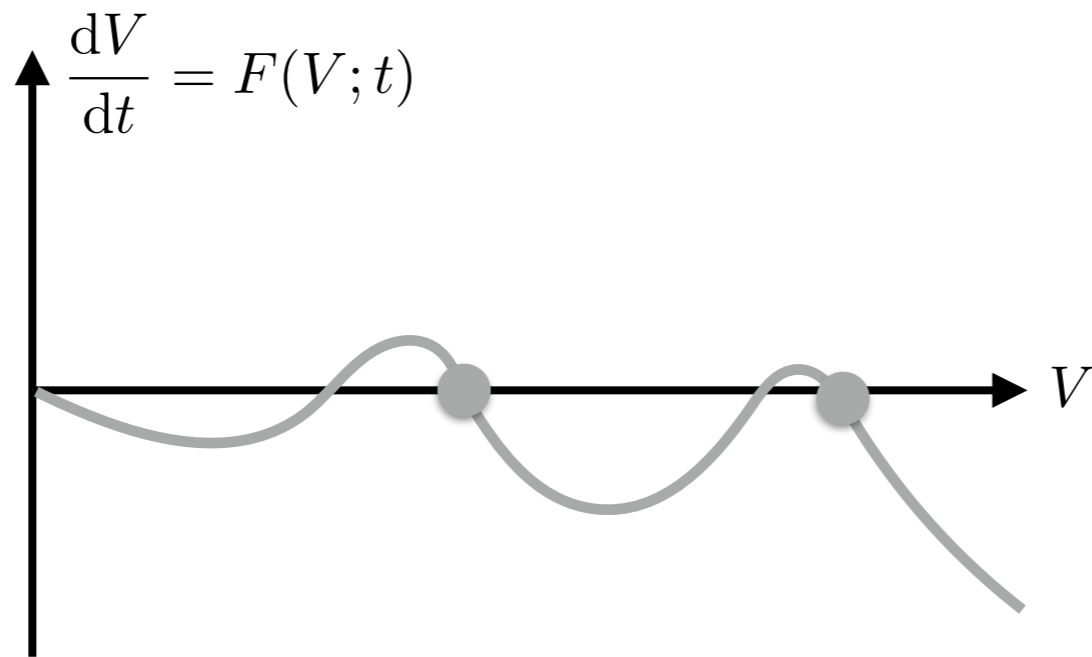
Rate of ice loss controlled by ice mechanics (& topography)

# Marine terminating glacier

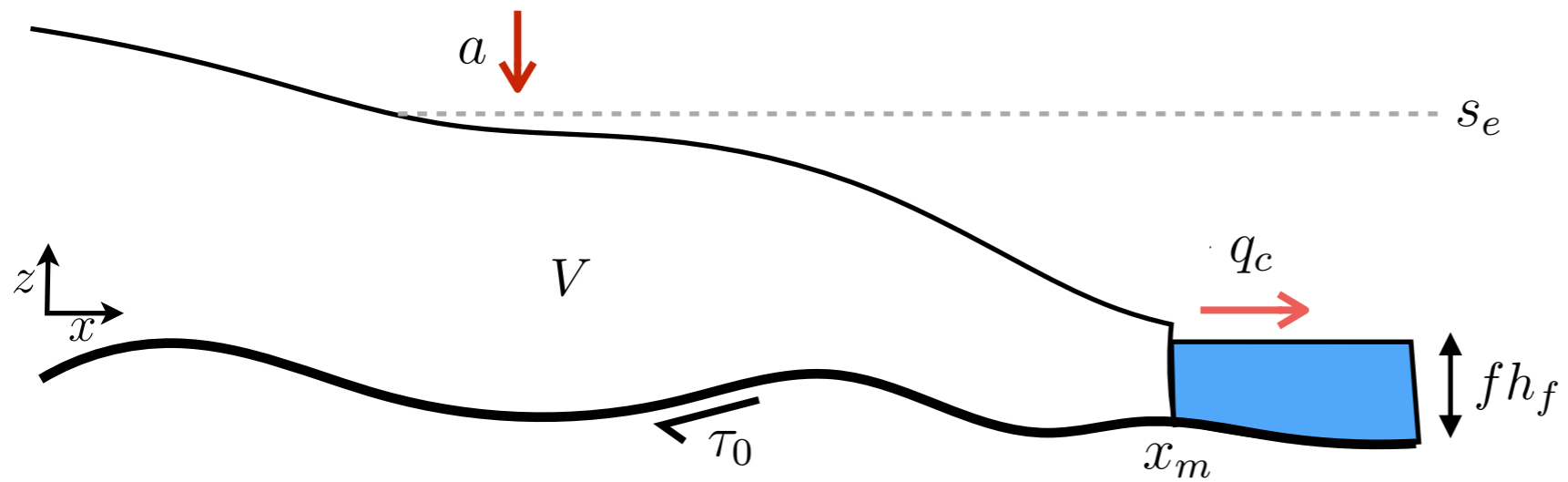


$$\frac{dV}{dt} = \int_0^{x_m} a \, dx - q_c$$

$$q_c = \frac{\rho_i g}{\eta_i \mu} \hat{Q}(f) h_f^3$$

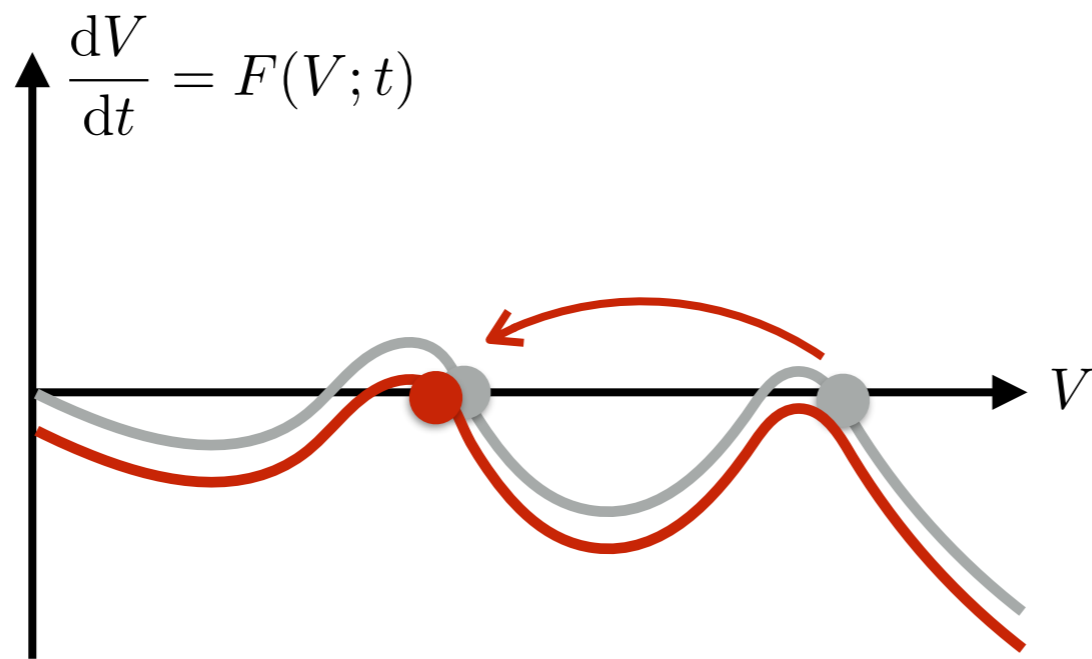


# Marine terminating glacier

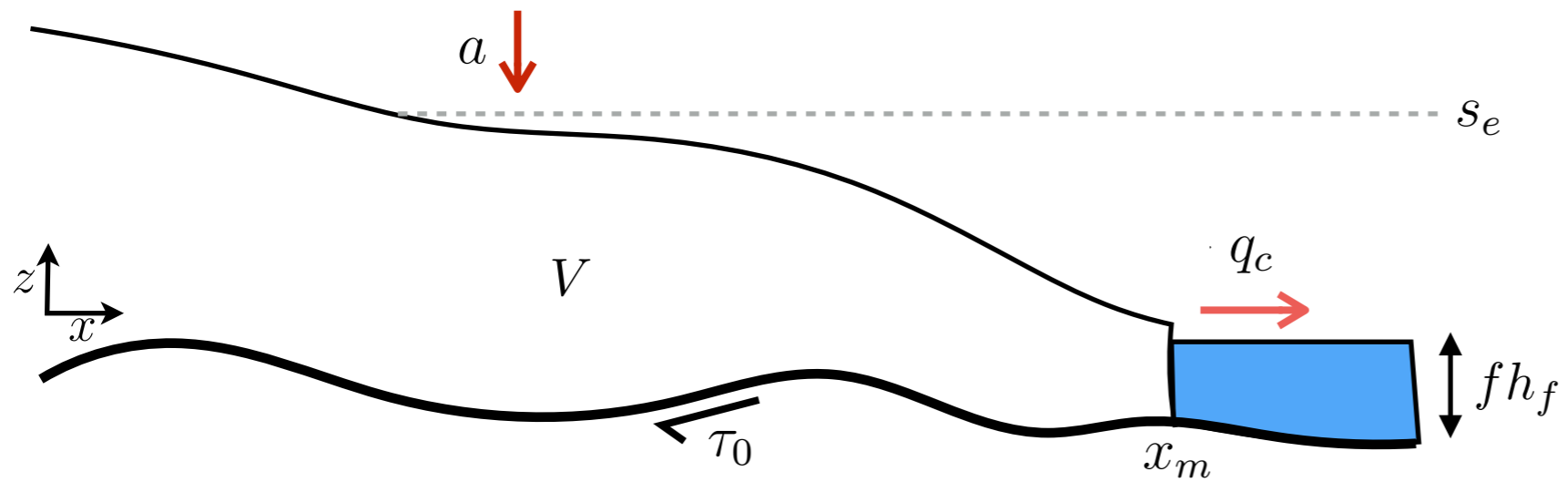


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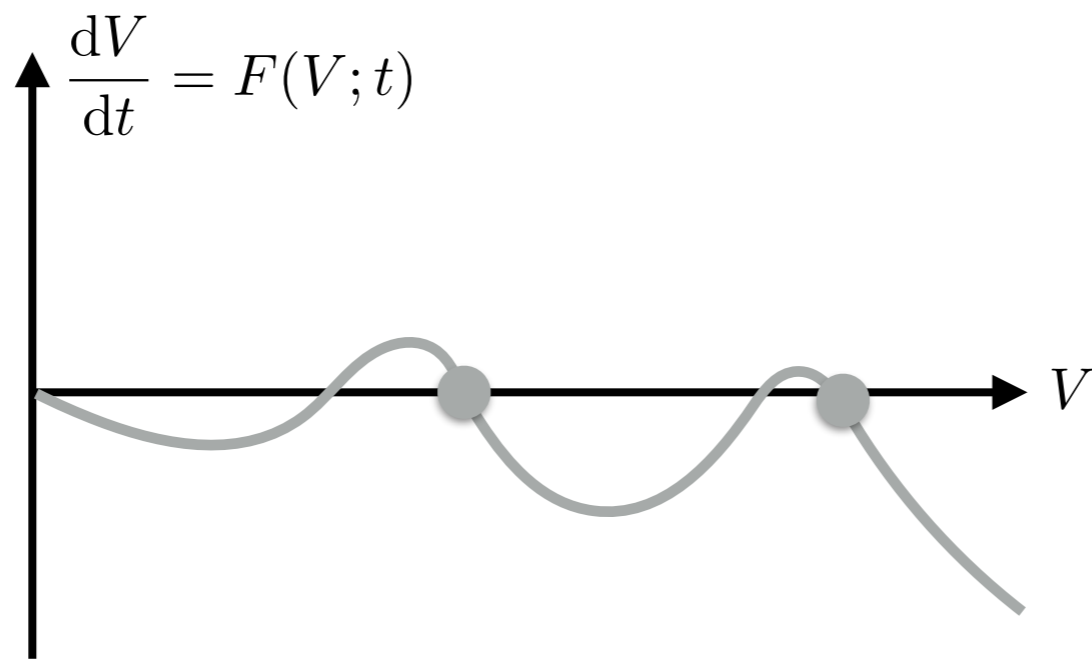


# Marine terminating glacier

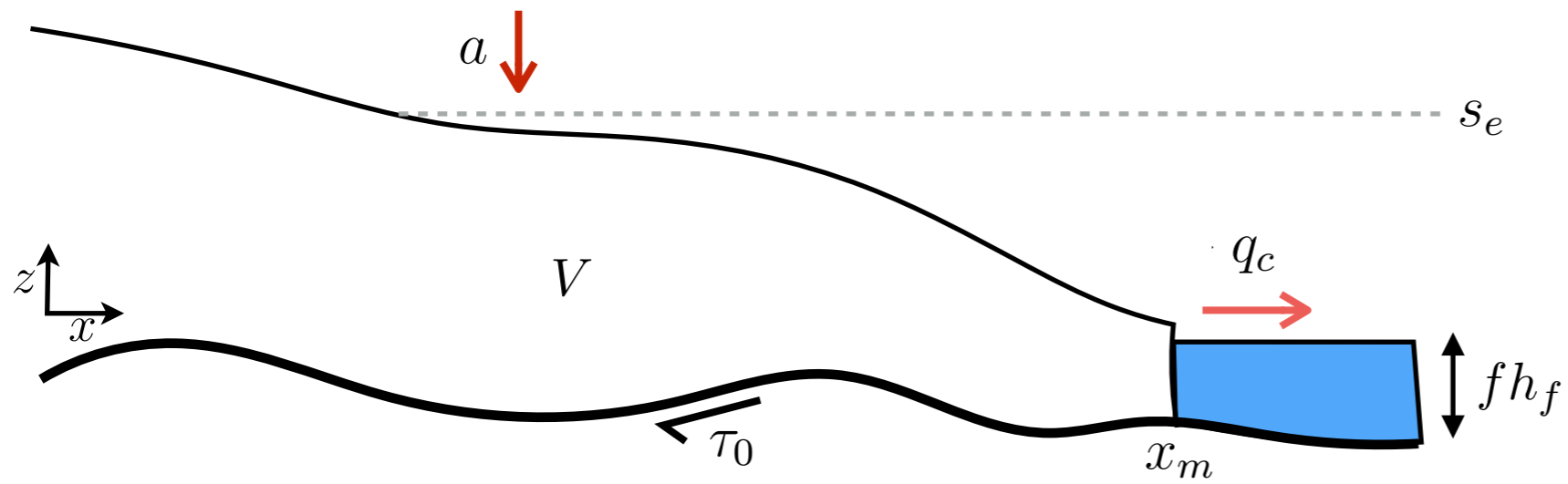


$$\frac{dV}{dt} = \int_0^{x_m} a \, dx - q_c$$

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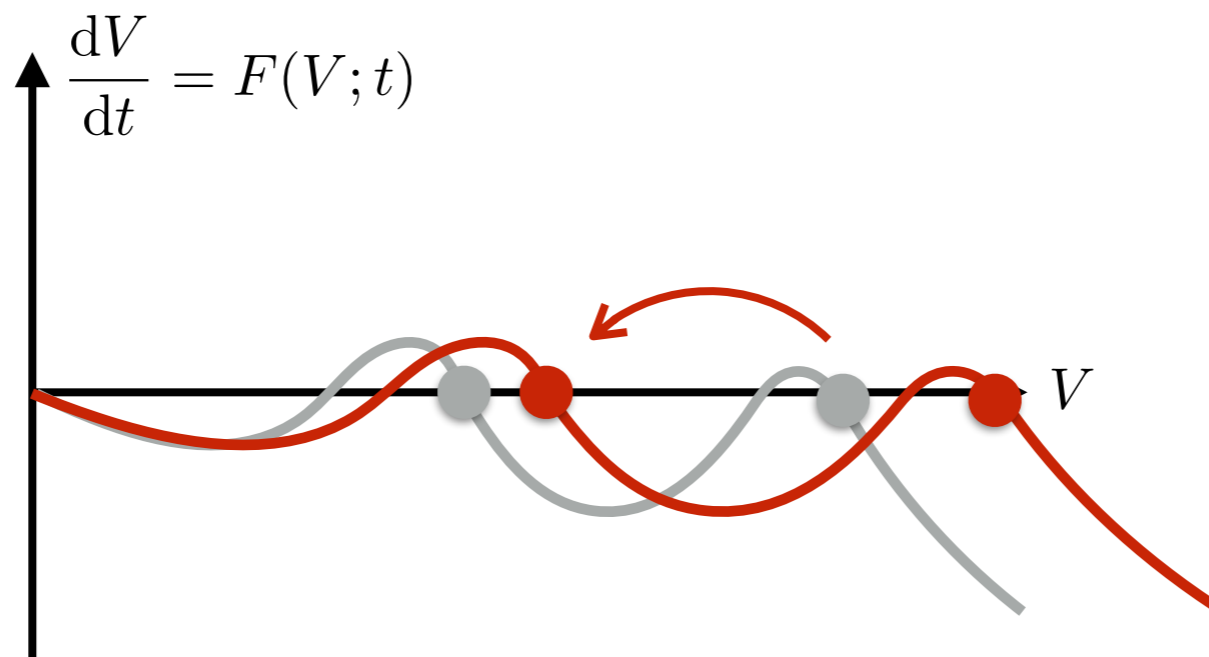


# Marine terminating glacier



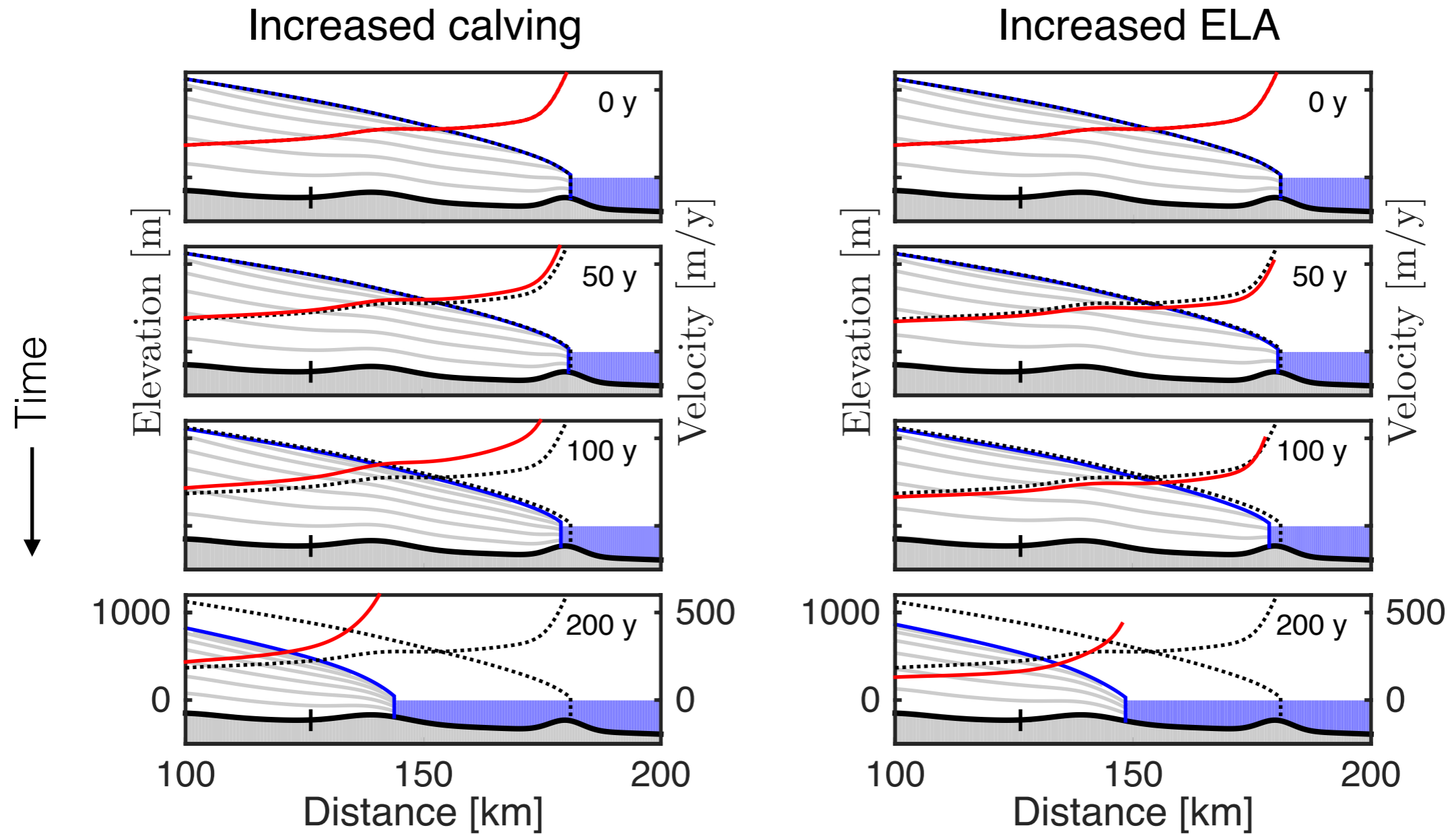
$$\frac{dV}{dt} = \int_0^{x_m} a \, dx - q_c$$

$$q_c = \frac{\rho_i g}{\eta_i \mu} \hat{Q}(f) h_f^3$$



# Marine terminating glacier

An essentially indistinguishable response occurs to an **increase in calving** or an **increase in ELA**



## Summary

Subglacial meltwater can both **increase** and **decrease** average ice speeds. Changes in **either direction** have potential to influence ice loss.

A simplified model suggests ice-sheet **slow down** can help induce **tidewater-glacier retreat**, and hence may facilitate rapid ice loss.

Recent retreat and speed-up of tidewater glaciers in Greenland may be as much a response to **inland** forcing as **ocean** forcing.