

In Form Continuum Methods for Industry ①

Exercises #2

1. Scale 2-D Navier Stokes eq's as suggested to get

$$\begin{cases} \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \\ R^* \left[\frac{\partial u}{\partial x} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right] = - \frac{\partial p}{\partial x} + \varepsilon^2 \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \\ \varepsilon^2 R^* \left[\frac{\partial v}{\partial x} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right] = - \frac{\partial p}{\partial y} + \varepsilon^4 \frac{\partial^2 v}{\partial x^2} + \varepsilon^2 \frac{\partial^2 v}{\partial y^2} \end{cases}$$

BCs $V = \frac{\partial u}{\partial y} = 0$ at $y=0$

& under the roller, the top surface is given by

$$h_0 h = h_m + R - \sqrt{R^2 - \frac{h_0^2 x^2}{\varepsilon^2}} \quad \text{in dimensionless variables}$$

Recall $\frac{R}{h_0} = 1/\varepsilon^2$, $\frac{h_m}{h_0} = \eta$

$$\therefore h(x) = \eta + \frac{1}{\varepsilon^2} \left[1 - \sqrt{1 - \varepsilon^2 x^2} \right]$$

$$\therefore \boxed{h(\eta) = \eta + \frac{1}{2} \varepsilon^2 + O(\varepsilon^4)}$$

BCs are $V = u h'(x)$, $u + \epsilon^2 h'(x) V = \gamma \sqrt{1 + \epsilon^2 h'(x)^2}$ (2)
at $y = h(x)$

now neglect terms of order ϵ^2 and ϵ^4 to get

$$\begin{aligned} \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} &= 0 \\ \frac{\partial u}{\partial y} &= \frac{\partial v}{\partial x} \\ \frac{\partial v}{\partial y} &= 0 \end{aligned}$$

&
$$\begin{aligned} V &= \frac{\partial u}{\partial y} = 0 \quad \text{at } y = 0 \\ u &= \gamma, \quad v = \gamma h'(x) \quad \text{at } y = h(x) \end{aligned} \quad (2.11)$$

(2.10):
$$\gamma \int_{-\alpha}^{\beta} \frac{dx}{\left(\eta + \frac{1}{2}x^2\right)^2} = \int_{-\alpha}^{\beta} \frac{dx}{\left(\eta + \frac{1}{2}x^2\right)^3}$$

Let $\underline{x = \sqrt{2\eta} X}$

$$\gamma \eta \int_{-\alpha/\sqrt{\eta}}^{\beta/\sqrt{\eta}} \frac{dX}{(1+X^2)^2} = \int_{-\alpha/\sqrt{\eta}}^{\beta/\sqrt{\eta}} \frac{dX}{(1+X^2)^3}$$

Recall $\underline{\alpha = \sqrt{2(1-\eta)}}$

(3)

$$\therefore \gamma = \frac{F_2 \left(\frac{\beta}{\sqrt{2\eta}} \right) + F_2 \left(\sqrt{\frac{1-\eta}{\eta}} \right)}{\eta \left[F_1 \left(\frac{\beta}{\sqrt{2\eta}} \right) + F_2 \left(\sqrt{\frac{1-\eta}{\eta}} \right) \right]}$$

with F_1 & F_2 as given.

(4)

```
In[1]:= Integrate[1 / (1 + x^2)^2, x]
```

```
Out[1]:= 1/2 (x/(1 + x^2) + ArcTan[x])
```

```
In[2]:= F1[x_] := 1/2 (x/(1 + x^2) + ArcTan[x])
```

```
In[3]:= Integrate[1 / (1 + x^2)^3, x]
```

```
Out[3]:= 1/8 (x(5 + 3 x^2)/(1 + x^2)^2 + 3 ArcTan[x])
```

```
In[4]:= F2[x_] := 1/8 (x(5 + 3 x^2)/(1 + x^2)^2 + 3 ArcTan[x])
```

```
In[5]:= ga[b_, y_] := (F2[b / Sqrt[2 y]] + F2[Sqrt[1 - y] / Sqrt[y]]) / y /  
(F1[b / Sqrt[2 y]] + F1[Sqrt[1 - y] / Sqrt[y]])
```

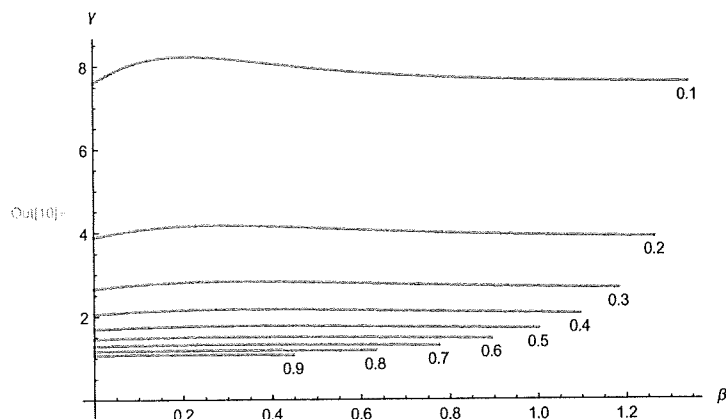
```
In[6]:= pl[y_] := Graphics[Text[y, {Sqrt[2 (1 - y)], ga[0, y] - 0.3}]]
```

```
In[7]:= pp[y_] := Plot[Evaluate[ga[b, y]], {b, 0, Sqrt[2 (1 - y)]}]
```

```
In[8]:= Table[pp[y], {y, 0.1, 0.9, 0.1}];
```

```
In[9]:= Table[pl[y], {y, 0.1, 0.9, 0.1}];
```

```
In[10]:= Show[%%, %, PlotRange -> All, AxesOrigin -> {0, 0}, AxesLabel -> {β, γ}]
```



```
In[11]:= (* γ versus β for different values of η. Note that the value of γ at β = α is the same as that at β = 0.
```

```
In each case there is a range of values of γ for which there are two positive roots for β. *)
```

```
In[12]:= (* Now let η = 1 - ε, β = ε1/2b and γ = 1 + εΓ/3, where ε << 1. *)
```

```
In[13]:= Series[Evaluate[ga[Sqrt[ε] b, 1 - ε]], {ε, 0, 1}] // FullSimplify
```

```
Out[13]:= 1 + 1/6 (4 + (√2 - b) b) ε + O[ε]3/2
```

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```
In[14]:= % - 1 - e Γ / 3
```

```
Out[14]=  $\left(\frac{1}{6} \left(4 + \left(\sqrt{2} - b\right) b\right) - \frac{\Gamma}{3}\right) e + O[e]^{3/2}$ 
```

```
In[15]:= Normal[% / e]
```

```
Out[15]=  $\frac{1}{6} \left(4 + \left(\sqrt{2} - b\right) b\right) - \frac{\Gamma}{3}$ 
```

```
In[16]:= Solve[% == 0, b]
```

```
Out[16]=  $\left\{\left\{b \rightarrow \frac{1}{\sqrt{2}} - \frac{\sqrt{9 - 4 \Gamma}}{\sqrt{2}}\right\}, \left\{b \rightarrow \frac{1}{\sqrt{2}} + \frac{\sqrt{9 - 4 \Gamma}}{\sqrt{2}}\right\}\right\}$ 
```

```
In[17]:= b /. %[[1]]
```

```
Out[17]=  $\frac{1}{\sqrt{2}} - \frac{\sqrt{9 - 4 \Gamma}}{\sqrt{2}}$ 
```

```
In[18]:= (*The smaller root is the physically relevant solution. We need Γ>
2 so that this root is positive, and Γ<9/4 so that it is real*)
```

2 Start from

(6)

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad \frac{\partial p}{\partial x} = \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right), \quad \frac{\partial p}{\partial y} = 0$$

where now

$$\mu = k \left[2 \left(\frac{\partial u}{\partial y} \right)^2 + 2 \left(\frac{\partial v}{\partial y} \right)^2 + \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right)^2 \right]^{\frac{n-1}{2}}$$

Non-dimensionalize $\underline{u} = U(\bar{u}, \varepsilon \bar{v})$
 $(x, y) = h_0(\varepsilon^{-1} \bar{x}, \bar{y})$

then drop tildes. Note that $\frac{\partial u}{\partial y}$ is much greater than any other velocity gradient component (by a factor of ε^{-2} at least, so to lowest order,

$$\mu \sim k \left(\frac{U}{h_0} \right)^{n-1} \left| \frac{\partial u}{\partial y} \right|^{n-1}$$

so define $\boxed{\mu_0 = k \left(\frac{U}{h_0} \right)^{n-1}}$

and non-dimensionalize pressure with

$$p = P_a + \frac{\mu_0 U}{\varepsilon h_0} \quad \tilde{p} = P_a + \frac{k}{\varepsilon} \left(\frac{U}{h_0} \right)^n \tilde{p}$$

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to get (with tildes dropped)

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \text{ (1)}, \quad \frac{\partial p}{\partial y} = 0, \quad \frac{dp}{dx} = \frac{\partial}{\partial y} \left(\left| \frac{\partial u}{\partial y} \right|^{n-1} \frac{\partial u}{\partial y} \right)$$

with BCs (as before)

$$\boxed{\begin{aligned} u = \frac{\partial u}{\partial y} = 0 & \quad \text{at } y=0 \\ u = \gamma, v = \gamma h'(x) & \quad \text{at } y=h(x) \end{aligned}}$$

where $h(x) = \eta + \frac{1}{2}x^2$ * lowest order.

By integrating (1) w.r.t. y , we get zero mass conservation as before:

$$\boxed{\int_0^h u \, dy \equiv 1}$$

Now solve for u : $\left| \frac{\partial u}{\partial y} \right|^{n-1} \frac{\partial u}{\partial y} = \frac{dp}{dx} y$

[satisfying BC at $y=0$.]

Take 1/2 of both sides:

$$\left| \frac{\partial u}{\partial y} \right|^n = \left| \frac{dp}{dx} \right| |y| \Rightarrow \left| \frac{\partial u}{\partial y} \right| = \left| \frac{dp}{dx} \right|^{1/n} |y|^{1/n}$$

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 $\therefore h, y \geq 0,$

$$\boxed{\frac{\partial u}{\partial y} = \frac{dp}{dn} \left| \frac{dp}{dn} \right|^{\frac{1}{n}-1} y^{1/n}} \quad \text{with } u = \gamma \text{ at } y = h$$

$$\Rightarrow \boxed{u = \gamma - \frac{dp}{dn} \left| \frac{dp}{dn} \right|^{\frac{1}{n}-1} \left(\frac{n}{n+1} \right) \left[h^{\frac{n+1}{n}} - y^{\frac{n+1}{n}} \right]}$$

$$\therefore \int_0^h u dy = \gamma h - \frac{dp}{dn} \left| \frac{dp}{dn} \right|^{\frac{1}{n}-1} \left(\frac{n}{n+1} \right) \left(\frac{n+1}{2n+1} \right) h^{\frac{2n+1}{n}}$$

$$\therefore \boxed{\gamma h - \frac{n}{2n+1} h^{\frac{2n+1}{n}} \frac{dp}{dn} \left| \frac{dp}{dn} \right|^{\frac{1}{n}-1} = 1}$$

with $\underline{p(-\alpha) = p(\beta) = 0}$

$$\therefore \underline{\frac{n}{2n+1} \frac{dp}{dn} \left| \frac{dp}{dn} \right|^{\frac{1}{n}-1} = \frac{\gamma}{h^{1+1/n}} - \frac{1}{h^{2+1/n}}}$$

$$\therefore \frac{n}{2n+1} \left| \frac{dp}{dn} \right|^{\frac{1}{n}} = |RHS|$$

$$\left(\frac{n}{2n+1} \right)^n \left| \frac{dp}{dn} \right| = |RHS|^n$$

$$\therefore \left(\frac{n}{2n+1} \right)^n \frac{dp}{dn} = RHS |RHS|^{n-1}$$

$$\therefore \int_{-\alpha}^{\beta} \left(\frac{\gamma}{h^{1+1/n}} - \frac{1}{h^{2+1/n}} \right) \left| \frac{\gamma}{h^{1+1/n}} - \frac{1}{h^{2+1/n}} \right|^{n-1} dh = 0$$

$$\cong \quad (A - (\lambda + \epsilon)I) \underline{x} = \underline{b} \quad (9)$$

where λ is an eigenvalue of A .

Try $\underline{x} \sim \underline{x}_0 + \epsilon \underline{x}_1 + \dots$

to get at lowest order

$$(A - \lambda I) \underline{x}_0 = \underline{b}$$

But $\det(A - \lambda I) = 0$ so this generically has no solution.

Instead, let $\underline{x} = \frac{1}{\epsilon} \underline{X}$ to get

$$\boxed{(A - (\lambda + \epsilon)I) \underline{X} = \epsilon \underline{b}}$$

$\underline{X} \sim \underline{X}_0 + \epsilon \underline{X}_1 + \dots$ gives

$$(A - \lambda I) \underline{X}_0 = \underline{0}$$

$$(A - \lambda I) \underline{X}_1 = \underline{b} + \underline{X}_0$$

etc...

Now since $\det(A - \lambda I) = 0$, there are non trivial solutions for $\underline{X}_0$ of the form...

$$\boxed{\underline{x}_0 = \alpha \underline{v}}$$

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where $\underline{v}$ is the eigenvector of A with eigenvalue λ (wlog normalized) so that $|\underline{v}|=1$, and the amplitude α is arbitrary.

At $O(\epsilon)$ we get

$$(A - \lambda I) \underline{x}_1 = \underline{b} + \alpha \underline{v}$$

But $\det(A - \lambda I) = 0 \Rightarrow$ this has no solutions for $\underline{x}_1$ unless the RHS is orthogonal to everything in the kernel of $A^a - \lambda I$.
Here A is real symmetric, so $A^a = A$ & we

need
$$(\underline{b} + \alpha \underline{v}) \cdot \underline{v} = 0$$

for solutions for $\underline{x}_1$ to exist. This solvability condition determines the amplitude of the leading-order solution:

$$\boxed{\alpha = -\underline{b} \cdot \underline{v}}$$

4 Given $\begin{cases} \epsilon^2 \text{St } r \frac{\partial u}{\partial t} = \frac{\partial}{\partial r} \left(r \frac{\partial u}{\partial r} \right) + \epsilon^2 r \frac{\partial^2 u}{\partial t^2} \\ \frac{\partial u}{\partial r} - \epsilon^2 R'(t) \frac{\partial u}{\partial t} = -\epsilon^2 u \sqrt{1 + \epsilon^2 R'(t)^2} \end{cases} \quad \text{at } r = R(t) \quad (11)$

Integrate w.r.t. r :

$$\epsilon^2 \text{St} \int_0^R r \frac{\partial u}{\partial t} dr = \left[r \frac{\partial u}{\partial r} \right]_0^R + \epsilon^2 \int_0^R r \frac{\partial^2 u}{\partial t^2} dr$$

$$\therefore \cancel{\epsilon^2} \text{St} \frac{\partial}{\partial t} \int_0^R u r dr = R \cancel{\epsilon^2} \left[R' \frac{\partial u}{\partial t} - u \sqrt{1 + \epsilon^2 R'^2} \right]_{r=R} + \cancel{\epsilon^2} \int_0^R r \frac{\partial^2 u}{\partial t^2} dr$$

$$\therefore \text{St} \frac{\partial}{\partial t} \int_0^R u r dr = \frac{\partial}{\partial t} \int_0^R \frac{\partial u}{\partial t} r dr - R u \sqrt{1 + \epsilon^2 R'^2} \Big|_{r=R}$$

using Leibnitz rule: $\frac{\partial}{\partial t} \int_0^{R(t)} f(r, t) dr$

$$= \int_0^R \frac{\partial f}{\partial t} dr + R'(t) f(r, t) \Big|_{r=R}$$

S11

(12)

```
In[1]:= R[z_] := R1 z + R2 z^2 + R3 z^3 + R4 z^4 + R5 z^5
```

```
In[2]:= u[z_, t_] := 1 + t f1[z/t] + t^2 f2[z/t] + t^3 f3[z/t] + t^4 f4[z/t] + t^5 f5[z/t]
```

```
In[3]:= s[t_] := s1 t + s2 t^2 + s3 t^3 + s4 t^4 + s5 t^5 (StD[u[z, t], t] -  
D[u[z, t], z, z] - 2 R'[z]/R[z] D[u[z, t], z] + 2 u[z, t]/R[z]) /. z -> t y
```

```
In[4]:= (StD[u[z, t], t] - D[u[z, t], z, z] - 2 R'[z]/R[z] D[u[z, t], z] + 2 u[z, t]/R[z]) /.  
z -> t y;
```

```
In[5]:= Series[%, {t, 0, 3}] // Simplify
```

```
Out[5]= 
$$\frac{2 - 2 R1 f1'[y] - R1 y f1''[y]}{R1 y t} - \frac{1}{R1^2 y}$$


$$\left( 2 R2 y - R1 (2 + R1 St y) f1[y] + R1 y (2 R2 + R1 St y) f1'[y] + 2 R1^2 f2'[y] + R1^2 y f2''[y] \right) +$$


$$\left( \frac{2 R2^2 y}{R1^3} - \frac{2 R3 y}{R1^2} - \frac{2 R2 f1[y]}{R1^2} + 2 \left( St + \frac{1}{R1 y} \right) f2[y] + \frac{2 R2^2 y f1'[y]}{R1^2} - \right.$$


$$\left. \frac{4 R3 y f1'[y]}{R1} - \frac{2 R2 f2'[y]}{R1} - St y f2'[y] - \frac{2 f3'[y]}{y} - f3''[y] \right) t + \left( 3 St f3[y] - \frac{1}{R1^4} \right.$$


$$\left. 2 \left( (R2^3 - 2 R1 R2 R3 + R1^2 R4) y^2 + R1 (-R2^2 + R1 R3) y f1[y] + R1^2 R2 f2[y] - \frac{R1^3 f3[y]}{y} \right) - \right.$$


$$\left. St y f3'[y] - \frac{1}{R1^3 y} 2 \left( (R2^3 - 3 R1 R2 R3 + 3 R1^2 R4) y^3 f1'[y] + \right. \right.$$


$$\left. R1 (- (R2^2 - 2 R1 R3) y^2 f2'[y] + R1 (R2 y f3'[y] + R1 f4'[y])) - f4''[y] \right) t^2 +$$


$$\left( 4 St f4[y] + \frac{1}{R1^5} 2 \left( (R2^4 - 3 R1 R2^2 R3 + 2 R1^2 R2 R4 + R1^2 (R3^2 - R1 R5)) y^3 - R1 (R2^3 - 2 R1 R2 R3 + \right. \right.$$


$$\left. R1^2 R4) y^2 f1[y] - R1^2 (-R2^2 + R1 R3) y f2[y] - R1^3 R2 f3[y] + \frac{R1^4 f4[y]}{y} \right) -$$


$$St y f4'[y] - \frac{1}{R1^4 y} 2 \left( - (R2^4 - 4 R1 R2^2 R3 + 4 R1^2 R2 R4 + 2 R1^2 (R3^2 - 2 R1 R5)) y^4 f1'[y] + \right.$$


$$R1 \left( (R2^3 - 3 R1 R2 R3 + 3 R1^2 R4) y^3 f2'[y] + \right.$$


$$\left. R1 (- (R2^2 - 2 R1 R3) y^2 f3'[y] + R1 (R2 y f4'[y] + R1 f5'[y])) \right) - f5''[y] \Big) t^3 + O[t]^4$$

```

```
In[6]:= eq = Simplify[% t];
```

```
In[7]:= u[s[t], t] - 1 // Simplify;
```

(13)

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```
In[8]:= Series[%, {t, 0, 5}] // Simplify
```

$$\begin{aligned} \text{Out[8]} = & f1[s1] t + (f2[s1] + s2 f1'[s1]) t^2 + \left(f3[s1] + s3 f1'[s1] + s2 f2'[s1] + \frac{1}{2} s2^2 f1''[s1] \right) t^3 + \\ & \left(f4[s1] + s3 f2'[s1] + s2 f3'[s1] + s2 s3 f1''[s1] + \right. \\ & \quad \left. f1'[s1] \left(s4 + \frac{2 s5}{R1 y} - \frac{2 s5 f1'[y]}{y} - s5 f1''[y] \right) + \frac{1}{2} s2^2 f2''[s1] + \frac{1}{6} s2^3 f1^{(3)}[s1] \right) t^4 + \\ & \left(f5[s1] + s3 f3'[s1] + s2 f4'[s1] + f2'[s1] \left(s4 + \frac{2 s5}{R1 y} - \frac{2 s5 f1'[y]}{y} - s5 f1''[y] \right) + \right. \\ & \quad \left. \frac{1}{2} f1''[s1] \left(s3^2 + 2 s2 \left(s4 + \frac{2 s5}{R1 y} - \frac{2 s5 f1'[y]}{y} - s5 f1''[y] \right) \right) + \right. \\ & \quad \left. s2 s3 f2''[s1] - \frac{1}{R1^2 y} s5 f1'[s1] (2 R2 y - R1 (2 + R1 St y) f1[y] + \right. \\ & \quad \left. R1 y (2 R2 + R1 St y) f1'[y] + 2 R1^2 f2'[y] + R1^2 y f2''[y]) + \frac{1}{2} s2^2 f3''[s1] + \right. \\ & \quad \left. \frac{1}{2} s2^2 s3 f1^{(3)}[s1] + \frac{1}{6} s2^3 f2^{(3)}[s1] + \frac{1}{24} s2^4 f1^{(4)}[s1] \right) t^5 + O[t]^6 \end{aligned}$$

```
In[9]:= bc1 = Simplify[% / t];
```

```
In[10]:= (D[u[z, t], z] - 1 - s'[t] /. z -> s[t]) // Simplify;
```

```
In[11]:= Series[%, {t, 0, 4}] // Simplify
```

```
Out[11]= (-1 - s1 + f1'[s1]) + (f2'[s1] + s2 (-2 + f1''[s1])) t +
  (-3 s3 + f3'[s1] + s3 f1''[s1] + s2 f2''[s1] +  $\frac{1}{2}$  s22 f1(3)[s1]) t2 +
  (-4 s4 -  $\frac{8 s5}{R1 y}$  + f4'[s1] -  $\frac{2 s5 f1'[y] (-4 + f1''[s1])}{y}$  + s4 f1''[s1] +
   $\frac{2 s5 f1''[s1]}{R1 y}$  + 4 s5 f1''[y] - s5 f1''[s1] f1''[y] + s3 f2''[s1] +
  s2 f3''[s1] + s2 s3 f1(3)[s1] +  $\frac{1}{2}$  s22 f2(3)[s1] +  $\frac{1}{6}$  s23 f1(4)[s1]) t3 +
  ( $\frac{10 R2 s5}{R1^2}$  +  $\frac{10 s5 f2'[y]}{y}$  + f5'[s1] +  $\frac{1}{R1 y}$  s5 (2 + R1 St y) f1[y] (-5 + f1''[s1]) -
   $\frac{2 R2 s5 f1''[s1]}{R1^2}$  -  $\frac{2 s5 f2'[y] f1''[s1]}{y}$  + s4 f2''[s1] +  $\frac{2 s5 f2''[s1]}{R1 y}$  - s5 f1''[y] f2''[s1] +
  5 s5 f2''[y] - s5 f1''[s1] f2''[y] + s3 f3''[s1] + s2 f4''[s1] +  $\frac{1}{2}$  s32 f1(3)[s1] +
  s2 s4 f1(3)[s1] +  $\frac{2 s2 s5 f1^{(3)}[s1]}{R1 y}$  - s2 s5 f1''[y] f1(3)[s1] -  $\frac{1}{R1 y}$  s5 f1'[y]
  (-10 R2 y - 5 R1 St y2 + y (2 R2 + R1 St y) f1''[s1] + 2 R1 f2''[s1] + 2 R1 s2 f1(3)[s1]) +
  s2 s3 f2(3)[s1] +  $\frac{1}{2}$  s22 f3(3)[s1] +  $\frac{1}{2}$  s22 s3 f1(4)[s1] +
   $\frac{1}{6}$  s23 f2(4)[s1] +  $\frac{1}{24}$  s24 f1(5)[s1]) t4 + O[t]5
```

```
In[12]:= bc2 = Simplify[%];
```

```
In[13]:= Simplify[eq /. t -> 0]
```

```
Out[13]=  $\frac{2 - 2 R1 f1'[y] - R1 y f1''[y]}{R1 y}$ 
```

```
In[14]:= DSolve[% == 0, f1, y]
```

```
Out[14]= {{f1 -> Function[{y},  $\frac{y}{R1} - \frac{C[1]}{y} + C[2]$ ]}}
```

```
In[15]:= f1[y] /. %[[1]] /. C[1] -> 0
```

```
Out[15]=  $\frac{y}{R1} + C[2]$ 
```

```
In[16]:= f1[y_] = %
```

```
Out[16]=  $\frac{y}{R1} + C[2]$ 
```

```
In[17]:= (* Must be bounded as y->0 *)
```

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```

In[18]:= {bc1, bc2} /. t -> 0
Out[18]=  $\left\{ \frac{s1}{R1} + C[2], -1 + \frac{1}{R1} - s1 \right\}$ 

In[19]:= Solve[% == 0, {C[2], s1}] // Simplify
Out[19]=  $\left\{ \left\{ C[2] \rightarrow \frac{-1 + R1}{R1^2}, s1 \rightarrow -1 + \frac{1}{R1} \right\} \right\}$ 

In[20]:= s1 = (s1 /. %[[1]]); f1[y_] = (Simplify[f1[y] /. %[[1]]]);
In[21]:= {eq, bc1, bc2} = Simplify[{eq, bc1, bc2} / t];
In[22]:= (* Now repeat procedure *)
In[23]:= Simplify[eq /. t -> 0];
In[24]:= DSolve[% == 0, f2, y]
Out[24]=  $\left\{ \left\{ f2 \rightarrow \text{Function}\left[\{y\}, \frac{1}{3 R1^3} \left( 3 (-1 + R1) y + \frac{1}{2} R1 (2 - 4 R2 - St + R1 St) y^2 - \frac{3 R1^3 C[1]}{y} \right) + C[2] \right] \right\} \right\}$ 

In[25]:= f2[y] /. %[[1]] /. C[1] -> 0 // Simplify
Out[25]=  $\frac{1}{6 R1^3} y (-6 + R1^2 St y + R1 (-2 + 4 R2 + St) y) + C[2]$ 

In[26]:= f2[y_] = %;
In[27]:= {bc1, bc2} /. t -> 0 // Simplify;
In[28]:= Solve[% == 0, {C[2], s2}] // Simplify;
In[29]:= s2 = (s2 /. %[[1]]); f2[y_] = (Simplify[f2[y] /. %[[1]]]);
In[30]:= {eq, bc1, bc2} = Simplify[{eq, bc1, bc2} / t];
In[31]:= Simplify[eq /. t -> 0];
In[32]:= DSolve[% == 0, f3, y] // Simplify
Out[32]=  $\left\{ \left\{ f3 \rightarrow \text{Function}\left[\{y\}, -\frac{1}{6 R1^5} (-1 + R1) (5 - 4 R1 + 8 R2 - 4 R1 R2 + 2 St - 3 R1 St + R1^2 St) y - \frac{1}{18 R1^4} (-1 + R1) (-6 + 12 R2 + 5 St - 4 R1 St + 8 R2 St - 4 R1 R2 St + 2 St^2 - 3 R1 St^2 + R1^2 St^2) y^2 - \frac{1}{18 R1^3} (-1 + 7 R2 - 10 R2^2 + 9 R1 R3 + 2 St - 2 R1 St - R2 St + R1 R2 St) y^3 - \frac{C[1]}{y} + C[2] \right] \right\} \right\}$ 

In[33]:= f3[y] /. %[[1]] /. C[1] -> 0 // Simplify;
In[34]:= f3[y_] = %;
In[35]:= {bc1, bc2} /. t -> 0 // Simplify;
In[36]:= Solve[% == 0, {C[2], s3}] // Simplify;

```

```

In[37]:= s3 = (s3 /. %[[1]]); f3[y_] = (Simplify[f3[y] /. %[[1]]]);
In[38]:= ss = Series[s[t], {t, 0, 3}] // Simplify
Out[38]= 
$$\left(-1 + \frac{1}{R1}\right) t - \frac{1}{6 R1^3} ((-1 + R1) (-1 - 4 R2 + (-1 + R1) St)) t^2 + \frac{1}{54 R1^5} (-1 + R1) (-4 - 46 R2^2 - 9 St - 5 St^2 + 2 R1^3 St^2 - R2 (23 + 27 St) - R1^2 (27 R3 + St (5 + 11 R2 + 9 St)) + R1 (3 + 30 R2^2 + 27 R3 + 14 St + 12 St^2 + R2 (15 + 38 St))) t^3 + O[t]^4$$

In[39]:= (* We need s'(t) ≥ 0 and hence R1 < 1 *)
In[40]:= (* As R1 → 1 we then need s''(t) ≥ 0 *)
In[41]:= Simplify[D[ss, t, t] /. t → 0]
Out[41]= 
$$-\frac{1}{3 R1^3} (-1 + R1) (-1 - 4 R2 + (-1 + R1) St)$$

In[42]:= Simplify[% / (1 - R1)] /. R1 → 1
Out[42]= 
$$\frac{1}{3} (-1 - 4 R2)$$

In[43]:= (* Therefore as R1 → 1 we need R2 ≥ -1/4 *)
In[44]:= (* As R1 → 1 and R2 → -1/4, we need s'''[0] ≥ 0 *)
In[45]:= Simplify[D[ss, t, t, t] /. t → 0]
Out[45]= 
$$\frac{1}{9 R1^5} (-1 + R1) (-4 - 46 R2^2 - 9 St - 5 St^2 + 2 R1^3 St^2 - R2 (23 + 27 St) - R1^2 (27 R3 + St (5 + 11 R2 + 9 St)) + R1 (3 + 30 R2^2 + 27 R3 + 14 St + 12 St^2 + R2 (15 + 38 St)))$$

In[46]:= (* R3 must exceed the zero of this coefficient *)
In[47]:= (* Let R2 → -1/4 and R1 → 1 carefully! *)
In[48]:= Solve[Evaluate[D[ss, t, t, t] /. t → 0] == 0, R3] // Simplify
Out[48]= 
$$\left\{ \left\{ R3 \rightarrow -\frac{1}{27 (-1 + R1) R1} (4 + 46 R2^2 + 9 St + 5 St^2 - 2 R1^3 St^2 + R1^2 St (5 + 11 R2 + 9 St) + R2 (23 + 27 St) - R1 (3 + 30 R2^2 + 14 St + 12 St^2 + R2 (15 + 38 St))) \right\} \right\}$$

In[49]:= FullSimplify[(R3 /. %[[1]]) /. R2 → -1/4]
Out[49]= 
$$\frac{9 + 2 (-1 + R1) St (-9 + 4 (-5 + 2 R1) St)}{216 R1}$$

In[50]:= FullSimplify[% /. R1 → 1]
Out[50]= 
$$\frac{1}{24}$$

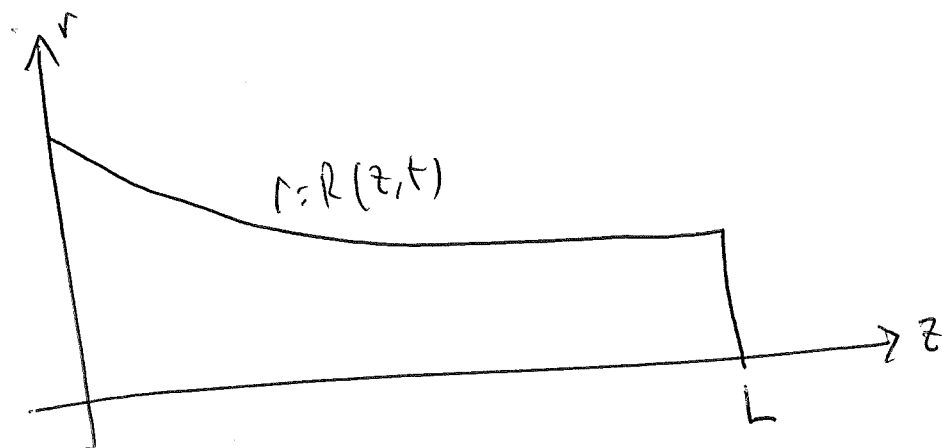
In[51]:= (* This gives us the required maximal radius profile

$$R_m(z) \sim z - \frac{z^2}{4} + \frac{z^3}{24} + \dots$$
 *)

```


6/11

(17)



BCs on free surface $r = R(z, t)$ are

① Kinematic BC $\boxed{u = \frac{\partial R}{\partial t} + w \frac{\partial R}{\partial z}}$ at $r = R(z, t)$

② Dynamic BC $\underline{\underline{\sigma}} \cdot \underline{\underline{n}} = \underline{\underline{0}}$ where $\underline{\underline{n}} = \frac{\underline{\underline{e}}_r - \frac{\partial R}{\partial z} \underline{\underline{e}}_z}{\sqrt{1 + \left(\frac{\partial R}{\partial z}\right)^2}}$

$\therefore \boxed{\begin{aligned} \sigma_{rr} &= \frac{\partial R}{\partial z} \sigma_{rt} \\ \sigma_{rt} &= \frac{\partial R}{\partial z} \sigma_{tz} \end{aligned}}$ at $r = R(z, t)$

Given $\frac{\partial}{\partial r}(ru) + \frac{\partial}{\partial z}(rw) = 0$

$\therefore \left[ru \right]_{r=0}^{r=R} + \int_0^R \frac{\partial}{\partial z}(rw) dr = 0$ (NB $ru \rightarrow 0$ as $r \rightarrow 0$)

$\therefore R \left(\frac{\partial R}{\partial t} + w \frac{\partial R}{\partial z} \right)_{r=R} + \int_0^R \frac{\partial}{\partial z}(rw) dr = 0$

$\therefore$ (by Leibniz)

$R \frac{\partial R}{\partial t} + \frac{\partial}{\partial z} \int_0^R wr dr = 0$

$$\therefore \boxed{\frac{\partial A}{\partial t} + \frac{\partial}{\partial t} (A \bar{\omega}) = 0} \quad \text{as required.}$$

(18)

Similarly from $\frac{\partial}{\partial r} (r \sigma_{rz}) + \frac{\partial}{\partial t} (r \sigma_{zt}) = 0$

$$\left[r \sigma_{rt} \right]_0^R + \int_0^R \frac{\partial}{\partial t} (r \sigma_{zt}) dr = 0$$

$$\therefore R \frac{\partial R}{\partial t} \sigma_{zt} \Big|_{r=R} + \int_0^R \frac{\partial}{\partial t} (r \sigma_{zt}) dr = 0$$

$$\therefore \frac{\partial}{\partial t} \int_0^R \sigma_{zt} r dr = 0 \quad (\text{Leibnitz})$$

$$\therefore \boxed{\frac{\partial}{\partial t} (A \bar{\sigma}_{zt}) = 0}$$

Now non-dimensionalize given equations & BCs:

$$u = \varepsilon U \tilde{u}, \quad \omega = U \tilde{\omega}, \quad z = L \tilde{z}, \quad r = \varepsilon L \tilde{r}, \quad t = \frac{L}{U} \tilde{t}$$

$$P = \frac{\mu U}{L} \tilde{p}, \quad \sigma_{zz} = \frac{\mu U}{L} \tilde{\sigma}_{zz}, \quad \sigma_{rz} = \frac{\varepsilon \mu U}{L} \tilde{\sigma}_{rz}$$

$$\sigma_{rr} = \frac{\varepsilon^2 \mu U}{L} \tilde{\sigma}_{rr}, \quad \sigma_{\theta\theta} = \frac{\varepsilon^2 \mu U}{L} \tilde{\sigma}_{\theta\theta}$$

Dimensional problem

(19)

$$\frac{1}{r} \frac{\partial}{\partial r} (ru) + \frac{\partial w}{\partial z} = 0$$

$$\frac{1}{r} \frac{\partial}{\partial r} (r \sigma_{rr}) - \frac{\sigma_{\theta\theta}}{r} + \frac{\partial \sigma_{rz}}{\partial z} = 0$$

$$\frac{1}{r} \frac{\partial}{\partial r} (r \sigma_{rz}) + \frac{\partial \sigma_{zz}}{\partial z} = 0$$

constitutive relations

$$\epsilon^2 \sigma_{rr} = -p + 2 \frac{\partial u}{\partial r}, \quad \sigma_{zz} = -p + 2 \frac{\partial w}{\partial z}$$

$$\epsilon^2 \sigma_{\theta\theta} = -p + \frac{2u}{r}, \quad \epsilon^2 \sigma_{rz} = \epsilon^2 \frac{\partial u}{\partial z} + \frac{\partial w}{\partial r}$$

B/c Boundedness as $r \rightarrow 0$ and

$$u = \frac{\partial R}{\partial t} + w \frac{\partial R}{\partial z}$$

$$\sigma_{rr} = \frac{\partial R}{\partial z} \sigma_{rz}, \quad \sigma_{rz} = \frac{\partial R}{\partial t} \sigma_{zz}$$

} at $r = R(z, t)$.

Now let $\epsilon \rightarrow 0$: leading order equations

$$\frac{\partial w}{\partial r} = 0 \Rightarrow \boxed{w = w(z, t)}$$

extensional flow.

Then

$$ru = -\frac{r^2}{2} \frac{\partial w}{\partial z}$$

(integration constant = 0)

$$\boxed{u = -\frac{r}{2} \frac{\partial w}{\partial z}}$$

then

$$p = 2 \frac{\partial u}{\partial r} = \frac{2u}{r} = - \frac{\partial w}{\partial z}$$

& hence

$$\sigma_{zz} = 3 \frac{\partial w}{\partial z}$$

so integrated eqs. obtained above lead to Frenkel model:

$$\begin{aligned} \frac{\partial A}{\partial t} + \frac{\partial}{\partial t} (A\omega) &= 0 \\ \frac{\partial}{\partial t} \left(3 A \frac{\partial w}{\partial z} \right) &= 0 \end{aligned}$$

□

7, (2-74) $\frac{d}{dz} (\tilde{A} + \tilde{w}) + \sigma e^{-mz} \tilde{A} = 0$

$$\frac{d\tilde{w}}{dz} + m(\tilde{A} + \tilde{w}) = m\tilde{F}$$

$$\tilde{A}(0) = \tilde{w}(0) = \tilde{w}(1) = 0$$

$$\therefore \tilde{A} = \tilde{F} - \tilde{w} - \frac{1}{m} \frac{d\tilde{w}}{dz}$$

Then $-\frac{1}{m} \frac{d^2\tilde{w}}{dz^2} + \sigma e^{-mz} \left[\tilde{F} - \tilde{w} - \frac{1}{m} \frac{d\tilde{w}}{dz} \right] = 0$

i.e. $\frac{d^2\tilde{w}}{dz^2} + \sigma e^{-mz} \left[\frac{d\tilde{w}}{dz} + m(\tilde{w} - \tilde{F}) \right] = 0$

subject to $\begin{cases} \tilde{w}(0) = \tilde{w}(1) = 0 \\ \frac{d\tilde{w}}{dz}(0) = m\tilde{F} \end{cases}$

Now make suggested substitution:

$$\tilde{w} = \tilde{F}(1 - \phi), \quad x = \frac{\sigma}{m} e^{-mz}$$

$$\Rightarrow \frac{d}{dz} = -mx \frac{d}{dx} ; \quad \frac{d\tilde{w}}{dz} = mx \tilde{F} \frac{d\phi}{dx}$$

$$m^2 x \frac{d}{du} \left(x \frac{d\phi}{du} \right) + mx \left[-mx \frac{d\phi}{du} + m\phi \right] = 0 \quad (22)$$

$$\therefore x^2 \frac{d^2\phi}{du^2} + x \frac{d\phi}{du} - x^2 \frac{d\phi}{du} + x\phi = 0$$

$$\therefore \boxed{\frac{d^2\phi}{du^2} + \left(\frac{1}{x} - 1 \right) \frac{d\phi}{du} + \frac{1}{x} \phi = 0}$$

bls :

$$\phi=1, \quad x \frac{d\phi}{du} = 1, \quad x = \lambda = \sigma/m$$

$$\phi=1, \quad x = \frac{\sigma}{m} e^{-u} = \lambda/D$$

Try $\phi = x-1$; $\phi' = 1$, $\phi'' = 0$

so LHS = $0 + \left(\frac{1}{x} - 1 \right) + \frac{1}{x} (x-1) = 0 \quad \checkmark$

So now try $\boxed{\phi(u) = (x-1)V(x)}$

Then $(x-1)V'' + 2V' + \left(\frac{1}{x} - 1 \right) \left((x-1)V' + V \right) + \left(\frac{x-1}{x} \right) V = 0$

$$\Rightarrow x(x-1)V'' + [2x - x^2 + 2x - 1]V' = 0$$

$$\Rightarrow \frac{V''}{V'} = \frac{x^2 - 4x + 1}{x(x-1)} = 1 - \frac{1}{x} - \frac{2}{x-1}$$

$$\therefore V' = \frac{C \cdot e^x}{x(x-1)^2}$$

BCs for V :

$$V(\lambda) = \frac{1}{\lambda-1}, \quad V'(\lambda) = -\frac{1}{\lambda(\lambda-1)^2},$$

$$V(\lambda/D) = \frac{1}{\lambda/D-1}$$

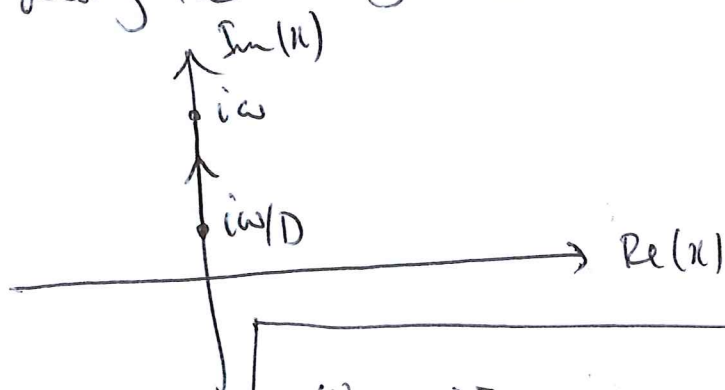
Hence evaluate C

$$V'(x) = -\frac{e^{x-\lambda}}{x(x-1)^2}$$

$$\& \int_{\lambda/D}^{\lambda} V'(x) dx = V(\lambda) - V(\lambda/D)$$

$$\dots \Rightarrow \int_{\lambda/D}^{\lambda} \frac{e^{x-\lambda}}{x(x-1)^2} dx = \frac{1}{\lambda/D-1} - \frac{1}{\lambda-1}$$

If $\lambda = i\omega$ where $\omega \in \mathbb{R}^+$ ($\omega \log$), then we integrate along the imaginary axis: let $x = i\xi$



$$\int_{\omega/D}^{\omega} \frac{e^{i\xi}}{\xi(1-i\xi)^2} d\xi + e^{i\omega} \left[\frac{1}{1-i\omega/D} - \frac{1}{1-i\omega} \right] = 0$$

```
In[1]:= Integrate[Exp[I x] / x / (1 - I x)^2, x]
```

```
Out[1]:= 
$$\frac{i \left( e^{i x} + (1 - i x) \text{ExpIntegralEi}[i x] \right)}{i + x}$$

```

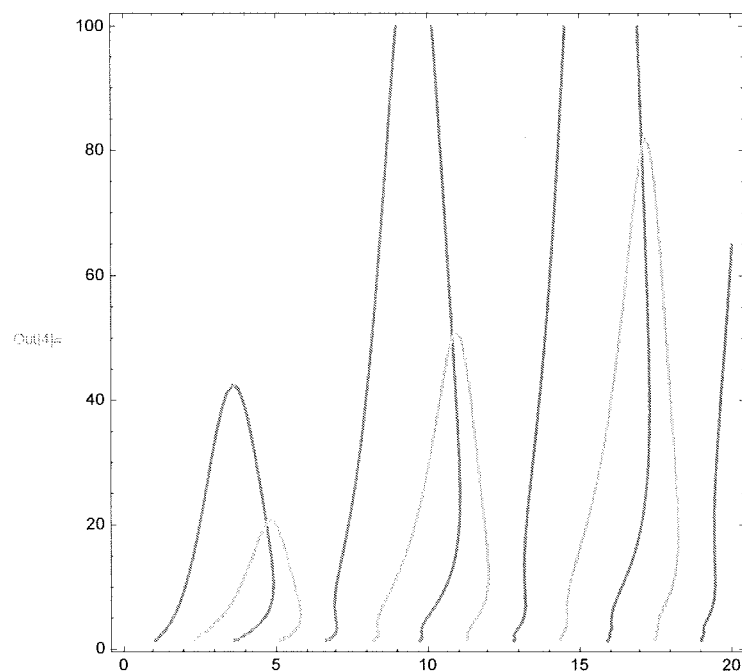
```
In[2]:= (% /. x -> w) - (% /. x -> w / Dr) + Exp[I w] (1 / (1 - I w / Dr) - 1 / (1 - I w)) // FullSimplify
```

```
Out[2]:= 
$$\frac{\text{Dr} \left( e^{i w} - e^{\frac{i w}{\text{Dr}}} \right) + (\text{Dr} - i w) \left( \text{ExpIntegralEi}[i w] - \text{ExpIntegralEi}\left[\frac{i w}{\text{Dr}}\right] \right)}{\text{Dr} - i w}$$

```

```
In[3]:= F[w_, Dr_] = %;
```

```
In[4]:= ContourPlot[Evaluate[{Re[F[w, Dr]] == 0, Im[F[w, Dr]] == 0}],  
  {w, 0, 20}, {Dr, 1.5, 100}, Contours -> {0}]
```



```
In[5]:= FindRoot[Evaluate[{Re[F[w, Dr]], Im[F[w, Dr]]} == 0], {{w, 4.5}, {Dr, 20}}]
```

```
Out[5]:= {w -> 4.66015, Dr -> 20.218}
```


$$8_{11} \quad A_0(x) = 1 + \alpha \left(x - \frac{1}{2}\right)^2, \quad \alpha > 0$$

$$\text{So } \begin{cases} A = 1 - f(\tau) + \alpha \left(\tau - \frac{1}{2}\right)^2 \\ Z = \tau + f(\tau) \int_0^\tau \frac{d\tau'}{1 - f(\tau') + \alpha \left(\tau' - \frac{1}{2}\right)^2} \end{cases}$$

Note critical value $f = 1$ where $A \rightarrow 0$ at $\tau = \frac{1}{2}$
and $Z \rightarrow \infty$

Calculate Integral:

$$Z = \tau + \frac{f}{\sqrt{\alpha(1-f)}} \left[\tan^{-1} \left(\frac{1}{2} \sqrt{\frac{\alpha}{1-f}} \right) + \tan^{-1} \left(\left(\tau - \frac{1}{2}\right) \sqrt{\frac{\alpha}{1-f}} \right) \right]$$

$$\therefore S = 1 + \frac{2f}{\sqrt{\alpha(1-f)}} \tan^{-1} \left(\frac{1}{2} \sqrt{\frac{\alpha}{1-f}} \right)$$

$$\text{As } f \rightarrow 1, \quad S \sim 1 + \frac{2}{\sqrt{\alpha(1-f)}} \cdot \frac{\pi}{2}$$

$$\Rightarrow f \sim 1 - \frac{\pi^2}{\alpha S^2} \quad \text{as } S \rightarrow \infty$$

$$\text{In}[1] = 1 / (1 - f + a (x - 1/2)^2)$$

$$\text{Out}[1] = \frac{1}{1 - f + a \left(-\frac{1}{2} + x\right)^2}$$

$$\text{In}[2] = \text{Integrate}[\%, x]$$

$$\text{Out}[2] = \frac{\text{ArcTan}\left[\frac{\sqrt{a} \left(-\frac{1}{2} + x\right)}{\sqrt{1-f}}\right]}{\sqrt{a} \sqrt{1-f}}$$

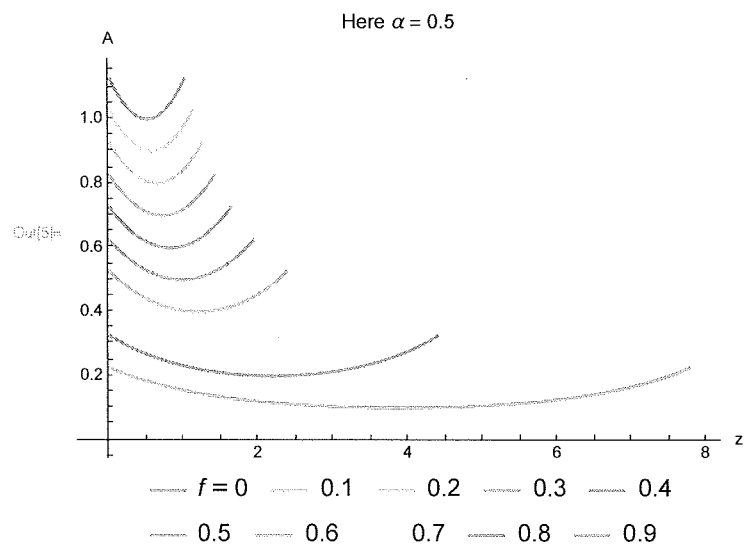
$$\text{In}[3] = (\% /. x \rightarrow ze) - (\% /. x \rightarrow 0) // \text{FullSimplify}$$

$$\text{Out}[3] = \frac{\text{ArcCot}\left[\frac{2\sqrt{1-f}}{\sqrt{a}}\right] + \text{ArcTan}\left[\frac{\sqrt{a} \left(-\frac{1}{2} + ze\right)}{\sqrt{1-f}}\right]}{\sqrt{a} \sqrt{1-f}}$$

$$\text{In}[4] = z = ze + f \%$$

$$\text{Out}[4] = ze + \frac{f \left(\text{ArcCot}\left[\frac{2\sqrt{1-f}}{\sqrt{a}}\right] + \text{ArcTan}\left[\frac{\sqrt{a} \left(-\frac{1}{2} + ze\right)}{\sqrt{1-f}}\right] \right)}{\sqrt{a} \sqrt{1-f}}$$

$$\text{In}[5] = \text{ParametricPlot}\left[\begin{array}{l} \text{Evaluate}[\text{Table}[\{z, 1 - f + a (ze - 1/2)^2\} /. a \rightarrow 0.5, \{f, 0, 0.9, 0.1\}], \\ \{ze, 0, 1\}, \text{PlotRange} \rightarrow \text{All}, \text{AspectRatio} \rightarrow 1/\text{GoldenRatio}, \\ \text{AxesOrigin} \rightarrow \{0, 0\}, \text{AxesLabel} \rightarrow \{z, "A"\}, \text{PlotLabel} \rightarrow \text{"Here } \alpha = 0.5", \\ \text{PlotLegends} \rightarrow \text{Join}[\{f == 0\}, \text{Table}[f, \{f, 0.1, 0.9, 0.1\}]]] \end{array}\right]$$



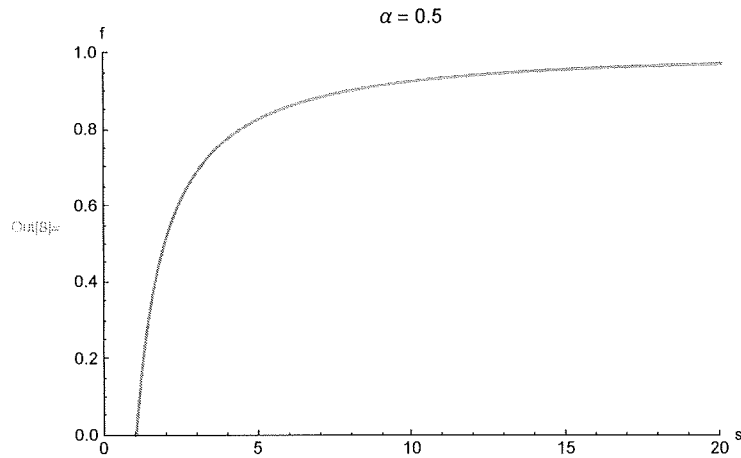
$$\text{In}[6] = z /. ze \rightarrow 1 // \text{FullSimplify}$$

$$\text{Out}[6] = 1 + \frac{2 f \text{ArcCot}\left[\frac{2\sqrt{1-f}}{\sqrt{a}}\right]}{\sqrt{a} \sqrt{1-f}}$$

$$\text{In}[7] = s = \%$$

$$\text{Out}[7] = 1 + \frac{2 f \text{ArcCot}\left[\frac{2\sqrt{1-f}}{\sqrt{a}}\right]}{\sqrt{a} \sqrt{1-f}}$$

```
In[5]:= p1 = ParametricPlot[Evaluate[{s, f}] /. a -> 0.5, {f, 0, 0.99},
  AspectRatio -> 1/GoldenRatio, PlotRange -> {{0, 20}, {0, 1}},
  AxesLabel -> {"s", "f"}, PlotLabel -> "α = 0.5"]
```



```
In[9]:= s - ss /. f -> 1 - Pi^2/a/ss^2 - f3/ss^3 - f4/ss^4 // Simplify // PowerExpand
```

$$\text{Out[9]} = 1 - ss - \frac{2 \left(\pi^2 ss^2 + a (f4 + f3 ss - ss^4) \right) \text{ArcCot} \left[\frac{2 \sqrt{a f4 + a f3 ss + \pi^2 ss^2}}{a ss^2} \right]}{a ss^2 \sqrt{a f4 + a f3 ss + \pi^2 ss^2}}$$

```
In[10]:= Assuming[a > 0 && f2 > 0 && ss > 0, Series[%, {ss, Infinity, 1}] // Simplify]
```

$$\text{Out[10]} = \left(1 - \frac{4}{a} - \frac{a f3}{2 \pi^2} \right) + \frac{\frac{3 a^2 f3^2}{8 \pi^4} - \frac{a f4}{2 \pi^2} - \frac{\pi^2}{a}}{ss} + O\left[\frac{1}{ss}\right]^2$$

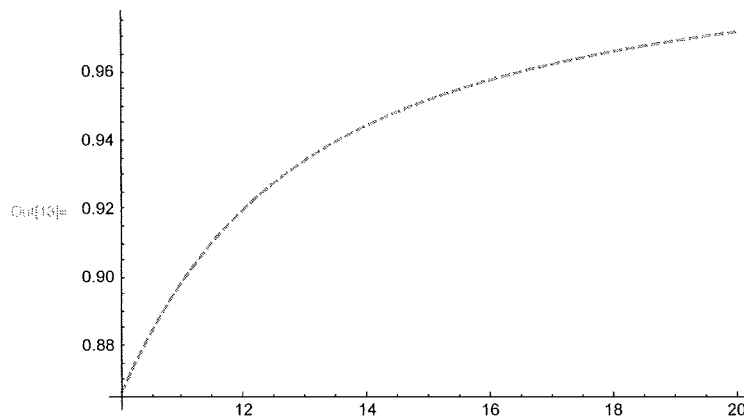
```
In[11]:= Solve[% == 0, {f3, f4}] // Simplify
```

$$\text{Out[11]} = \left\{ \left\{ f3 \rightarrow \frac{2(-4+a)\pi^2}{a^2}, f4 \rightarrow \frac{\pi^2(48+3a^2-2a(12+\pi^2))}{a^3} \right\} \right\}$$

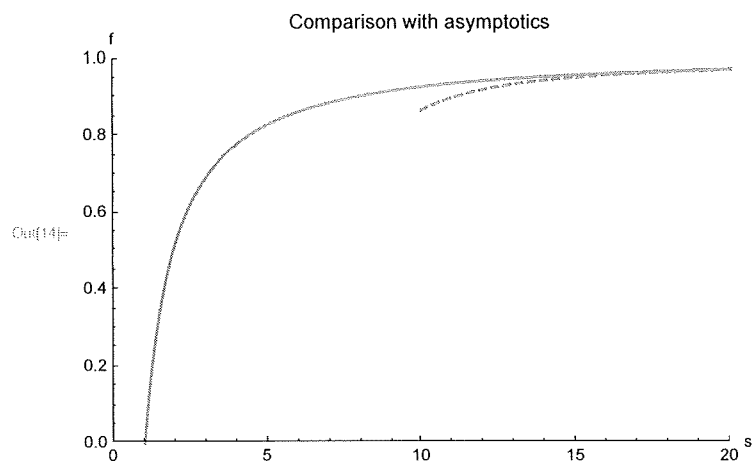
```
In[12]:= 1 - Pi^2/a/ss^2 - f3/ss^3 - f4/ss^4 /. %[[1]]
```

$$\text{Out[12]} = 1 - \frac{\pi^2(48+3a^2-2a(12+\pi^2))}{a^3 ss^4} - \frac{2(-4+a)\pi^2}{a^2 ss^3} - \frac{\pi^2}{a ss^2}$$

```
In[13]:= Plot[Evaluate[% /. a -> 0.5], {ss, 10, 20}, PlotStyle -> Dashing[{0.01, 0.01}]]
```



`In[14] := Show[p1, %, PlotLabel -> "Comparison with asymptotics"]`



Q// Write Trouton model with inertia in the form: (29)

$$\frac{DA}{Dt} = -A \frac{\partial w}{\partial z}, \quad R A \frac{DW}{Dt} = \frac{\partial}{\partial z} \left(3A \frac{\partial w}{\partial z} \right)$$

where $D/Dt = \partial/\partial t + w \partial/\partial z$

Now $\frac{\partial}{\partial z} \left(A \frac{\partial w}{\partial z} \right) = - \frac{\partial}{\partial z} \left[\frac{\partial A}{\partial t} + w \frac{\partial A}{\partial z} \right]$

$$= - \frac{D}{Dt} \left(\frac{\partial A}{\partial z} \right) - \frac{\partial w}{\partial z} \frac{\partial A}{\partial z} = \frac{R}{3} A \frac{DW}{Dt}$$

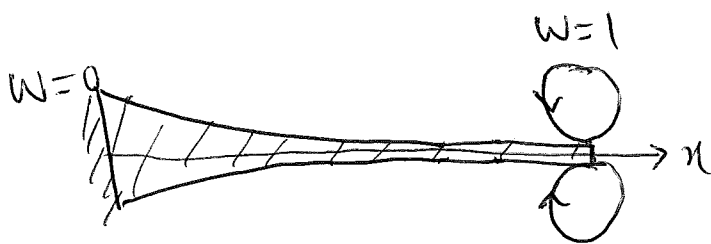
$$\therefore \frac{R}{3} A \frac{DW}{Dt} + \frac{D}{Dt} \left(\frac{\partial A}{\partial z} \right) - \frac{\partial A}{\partial z} \cdot \frac{1}{A} \frac{DA}{Dt} = 0$$

Divide by A & use product rule:

$$\boxed{\frac{D}{Dt} \left(\frac{R}{3} W + \frac{1}{A} \frac{\partial A}{\partial z} \right) = 0}$$

Initially at rest & uniform: $w=0, A=1$ at $t=0$

For $t=0$, hold $z=0$ fixed & draw off at speed 1
at $z=1$:



From above eqⁿ we get

$$\frac{R}{3}W + \frac{1}{A} \frac{\partial A}{\partial t} = \text{const. along characteristics}$$

$$dz/dt = W$$

From IIs, this constant is zero. So

$$W \equiv - \frac{3}{RA} \frac{\partial A}{\partial t}$$

& conservation of mass $\Rightarrow$ $\frac{R}{3} \frac{\partial A}{\partial t} = \frac{\partial^2 A}{\partial z^2}$ heat equation!

BCs $W(0)=0$, $W(1)=1$

$$\Rightarrow \frac{\partial A}{\partial t} = 0 \quad \text{at } z=0$$

$$\frac{\partial A}{\partial z} + \frac{R}{3}A = 0 \quad \text{at } z=1$$

& $A=1 \quad \text{at } t=0$

Just separate variables: $A(z,t) = \cos(kz) e^{-3k^2 t/R}$

works $\forall k$ satisfying $-k \sin k + \frac{R}{3} \cos k = 0$

i.e. $k \tan k = R/3$

Then by linear superposition:

$$A(z,t) = \sum_{n=1}^{\infty} C_n \cos(k_n z) e^{-3k_n^2 t/R}$$

(31)

where $C_n \int_0^1 \cos^2(k_n z) dz = \int_0^1 \cos(k_n z) dz$

by using initial conditions & orthogonality:

$$\left(\int_0^1 \cos(k_n z) \cos(k_m z) dz = 0 \text{ for } n \neq m \right)$$

Evaluate integrals: $C_n = \frac{2 \sinh(k_n)}{k_n + \sinh(k_n) \cosh(k_n)}$

```

In[1]:= kn[n_, R_] :=
  (k /. FindRoot[k Sin[k] == R Cos[k] / 3, {k, (n - 1) Pi, (n - 1 / 2) Pi}][[1]])

In[2]:= Integrate[Cos[k z], {z, 0, 1}] / Integrate[Cos[k z]^2, {z, 0, 1}] // FullSimplify

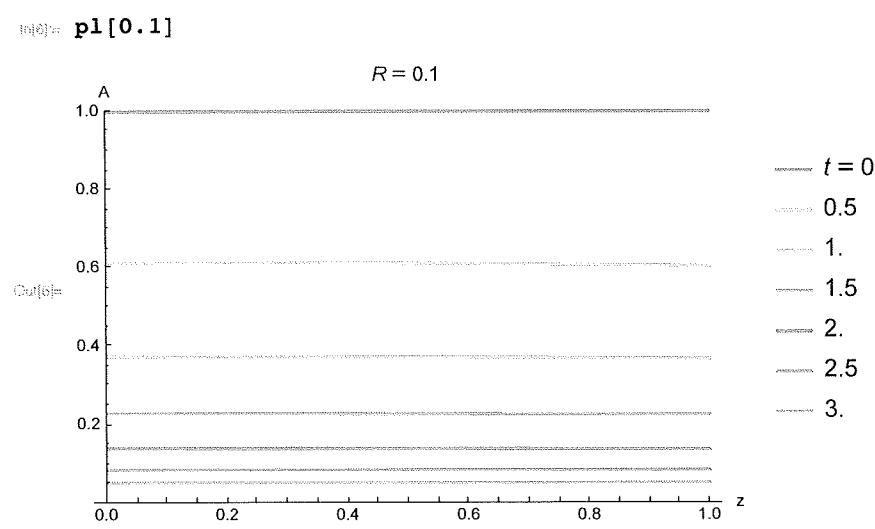
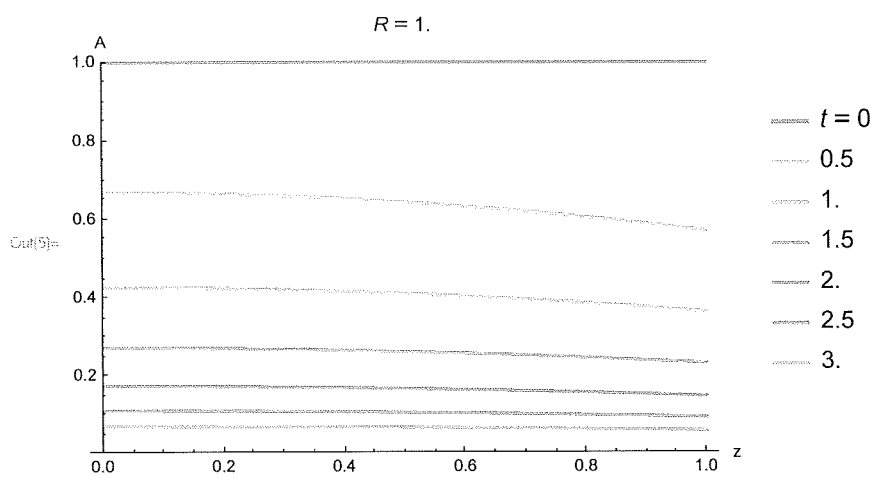
Out[2]:=
  2 Sin[k]
  -----
  k + Cos[k] Sin[k]

In[3]:= A[R_, m_, z_, t_] := Sum[
  -----
  2 Sin[kn[n, R]]
  kn[n, R] + Cos[kn[n, R]] Sin[kn[n, R]]
  Cos[kn[n, R] z] Exp[-3 kn[n, R]^2 t / R], {n, 1, m}]

In[4]:= pl[Re_] :=
  Plot[Evaluate[Join[{1}, Table[Evaluate[A[Re, 20, z, t]], {t, 0.5, 3, 0.5}]]],
  {z, 0, 1}, PlotLabel -> R == Re, AxesLabel -> {"z", "A"},
  PlotLegends -> Join[{t == 0}, Table[t, {t, 0.5, 3, 0.5}]], PlotRange -> {0, 1}]

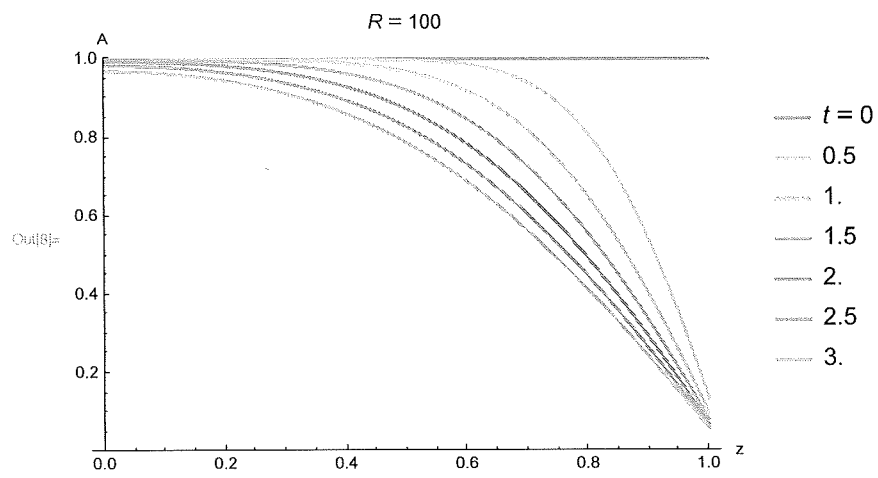
In[5]:= pl[1.0]

```



In[7]:= (* When R is small, A is approximately spatially uniform.*)


```
In[6] := pl[100]
```



```
In[7] := (* When R is large, there is a boundary layer near z=1. *)
```