

Nieether's Theorem & Factorisation algebras

Intuitively, Nieether's thm says "symmetries acting on a classical field theory yield conservation laws".

We want to see what this means in our framework.

Symmetries: These will be local L_∞ algebras.

Def: a L_∞ algebra is a $\mathbb{Z}$ -graded v.s. $\mathfrak{g}$ with $\forall n \geq 1$ a multilinear map $l_n: \mathfrak{g} \otimes \dots \otimes \mathfrak{g} \rightarrow \mathfrak{g}$ of degree $2-n$

s.t.

- (graded skew symmetry) $l_n(x_1, \dots, x_i, x_{i+1}, \dots, x_n) = (-1)^{x_i(x_{i+1}+1)} l_n(x_1, \dots, x_{i+1}, x_i, \dots, x_n)$
- (strong Jacobi identity) $\sum_{k=1}^n (-1)^k \sum_{\substack{\text{partition of } \{1, \dots, n\} \\ \text{into } \{i_1, \dots, i_k\} \\ \{j_1, \dots, j_{n-k}\}}} \text{sign}(\sigma) l_{n-k+1}(l_k(x_{i_1}, \dots, x_{i_k}), x_{j_1}, \dots, x_{j_{n-k}})$

think of this as a dgl, where l_1 is the differential, l_2 is $[\cdot, \cdot]$, but now Jacobi identity only holds up to homotopy.

Def: A local L_∞ algebra on mfd M consists of

- a graded v.s. $L \rightarrow M$
- a differential operator $d: L \rightarrow L$ of whom. deg 1, sq. zero.
- $\forall n \geq 2$ poly differential operators $l_n: L^{\otimes n} \rightarrow L$ which are alternating, deg $2-n$, & endow L with the structure of an L_∞ algebra.

Recall: differential operator $D: E_1 \rightarrow E_2$ map of sheaves of sections of E_1, E_2 which factors

as $E_1 \rightarrow \mathcal{J}(E_1) \rightarrow E_2$
 (sheaf of sections of jets)

Fact: Let $\mathcal{D}_M$ be the sheaf of diff operators on M with smooth coefficients
 differential operator $E \rightarrow F \iff$ map of $\mathcal{D}_M$ -modules $\mathcal{J}(E) \rightarrow \mathcal{J}(F)$

similarly: poly diff. operator $E \otimes \dots \otimes E_n \rightarrow F \iff$ map of $\mathcal{D}_M$ -modules $\mathcal{J}(E)^{\otimes_n} \rightarrow \mathcal{J}(F)$

Hence: local L_∞ alg structure on L is the same thing as L_∞ alg structure on $\mathcal{J}(L)$

in the category of $\mathcal{D}_M$ -modules.

e.g. L_∞ alg $\mathfrak{g}$ acts $\iff \mathfrak{g} \otimes \mathcal{D}_M^*$ acting

Classical field theories: Recall free classical field theory on a mfd M consists of

- $\mathbb{Z}$ -graded v.s. E with (-1) -shifted symplectic structure $\langle, \rangle$
- diff. operator $\mathcal{Q}: E \rightarrow E$ wh. degree +1, sq zero, skew-self-adj w.r.t $\langle, \rangle$
- i.e. $(E, \mathcal{Q})$ is elliptic.

This is a "free classical field theory" because E-L equation $(\Delta + \mathcal{Q}^2)\phi = 0$

Def: This only reverses linear diff. eqs. We require a more general concept.

Def: A classical field theory on n -d M consists of

local L_∞ alg L on M

invariant pairing of wh. deg. -3 $\langle, \rangle: L \otimes L \rightarrow \text{Der}_M[-3]$

sym., non-degen., $(\alpha_1 \mapsto \alpha_{n+1}) \mapsto \langle \ln(\alpha_1 \mapsto \alpha_n), \alpha_{n+1} \rangle$ ^{graded} antisymmetric for $\alpha_i \in L_c$ _{$n \geq 1$}

(L, d) is elliptic.

free field theory is a special case where the $\ln$ are trivial $\forall n \geq 2$

$\geq$ grading is shifted by $+1$

$\ln$ encode higher-order terms in the E-L eqs

g interacting scalar field theory

$$L(U): C^\infty(U)[-1] \xrightarrow{\Delta} C^\infty(U)[-2]$$

$$\text{pairing: } \langle \varphi_1, \varphi_2 \rangle = \int_U \varphi_1 \varphi_2$$

$$\ln \text{ trivial except for } n=3: \ln_3(\varphi_1, \varphi_2, \varphi_3) = \underbrace{\varphi_1 \varphi_2 \varphi_3}_{\text{all in deg 1}} = \underbrace{\varphi_1 \varphi_2}_{\text{in deg 2}} \varphi_3$$

$$E-L: d\varphi + \sum \frac{1}{n!} \ln(\varphi \mapsto \varphi) = 0 \quad \text{i.e.} \quad \Delta\varphi + \frac{1}{3}\varphi^3 = 0.$$

$$\text{observables: } \text{Obs}^U(U) = CE^*(L(U))$$

Recall: For L_∞ alg $\mathfrak{g}$, $CE^*(\mathfrak{g})$ is $\widehat{\text{Sym}}(\mathfrak{g}[\mathfrak{g}^\vee]) = \prod_{n=0}^{\infty} ((\mathfrak{g}[\mathfrak{g}^\vee])^{\otimes n})_{S_n}$ with differentials dual to the $\ln$

m (Costello-Gwilliam) $\exists$ factorization algebra $\widetilde{\text{Obs}}^U$ that's weakly equivalent to Obs^U .

in the case of free field theories, this is $\widetilde{\text{Obs}}^U(U) = \widehat{\text{Sym}}(E_c^!(U))$

Def

the level of Lie algebras, action of $\mathfrak{g}$ on $\mathfrak{h} \hookrightarrow \mathfrak{g} \rightarrow \text{Der}(\mathfrak{h}) \hookrightarrow$ semidirect product $\mathfrak{g} \ltimes \mathfrak{h}$.

analogously, we define

Def: Action of L_∞ alg $\mathfrak{g}$ on L_∞ alg $\mathfrak{h}$ is an L_∞ alg structure on $\mathfrak{g} \ltimes \mathfrak{h}$, denoted $\mathfrak{g} \ltimes \mathfrak{h}$, s.t. $0 \rightarrow \mathfrak{h} \rightarrow \mathfrak{g} \ltimes \mathfrak{h} \rightarrow \mathfrak{g} \rightarrow 0$ exact. (map of L_∞ algs)

Def: Action of local L_∞ alg L on L_∞ alg M is a local L_∞ alg structure on $L \ltimes M$, denoted $L \ltimes M$, s.t. $0 \rightarrow M \rightarrow L \ltimes M \rightarrow L \rightarrow 0$ (map of sheaves of L_∞ algs)

Def: local L_∞ alg acts on classical field theory M if it preserves the pairing:

$$\forall \alpha_1, \alpha_2 \in L_c, \beta_1, \beta_2 \in M_c, \langle \ln_3(\alpha_1 \mapsto \alpha_2, \beta_1 \mapsto \beta_2), \beta_3 \rangle \text{ graded totally antisym in } \beta_i$$

central extensions associated to actions: Given local L_∞ alg acting on classical field theory M , our version of Noether's theorem involves a central extension $\tilde{L}_c$ of L_c

Def: $CE_{red}^*(\mathcal{F}(L)) = \prod_{n \geq 0} \text{Hom}_{\mathcal{C}^0\text{-Mod}}(\mathcal{F}(L)[1]^{\otimes n}, \mathcal{C}^0_M)^{\otimes n}$

(with different) manifold of equivariant sections

Chevalley-Eilenberg co of $\mathcal{F}(L)$, in the category of $\mathcal{C}^0$ -modules.

$n \geq 0$: quotient out constants, hence "reduced" "sheaf of logarithms"

Def: $CE_{red,loc}^*(L) = \text{Dens}_M \otimes_{\mathcal{C}^0_M} C_{red}^*(\mathcal{F}(L))$

(sheaf of smooth densities on M) local Chevalley-Eilenberg complex

think of as acting on L_c via integration sheaf of dg algebras

"local action functionals on L_c "

Def: L local L_∞ alg on M . A k -shifted local central extension on M is an L_∞ structure on the presheaf $\tilde{L} = L_c \oplus \mathbb{C}[k]$ s.t.

- $0 \rightarrow \mathbb{C}[k] \rightarrow \tilde{L}_c \rightarrow L_c \rightarrow 0$ is a SES of presheaves
- the structure maps $\tilde{L}_c \rightarrow \mathbb{C}[k]$ determining the L_∞ structure on $\tilde{L}_c$ are local action functionals.

Lemma: $\exists$ bijection $\{k\text{-shifted local central extensions}\} / \cong \longleftrightarrow H^{k+2}(CE_{red,loc}^*(L)(M))$

(sanity check: elt in $\text{deg } k+2$ takes r elts in $\mathcal{F}(L)[1]$ which sum to $-2-k$ i.e. r inputs in $\mathcal{F}(L)$ with total $\text{deg } 2-k$ splits out a constant; this exactly defines L_r of $\text{deg } 2-r$; condition on L_∞ alg $\leftrightarrow$ cocycle condition)

We want to construct a central $+1$ -shifted central extension i.e. identify an elt in $H^1(CE_{red,loc}^*(L)(M))$

Sketch: Recall that M has a (-3) -shifted invariant pairing. This induces a β_0 structure on $CE_{red,loc}^*(U(M))[1]$

Idea: For a L_∞ alg h with (-3) -shifted invariant pairing, this induces $h[1]^\vee \cong h[2]$ so we get $h[1]^\vee \otimes h[1]^\vee \cong h[1]^\vee \otimes h[2] \rightarrow \mathbb{C}$ of $\text{deg } 1$

$\text{Sym}^2(h[1]^\vee) \quad \quad \quad \text{Sym}^2(h[1]^\vee)$

extend by Leibniz rule to get a β_0 structure on $CE^*(h)$. $\hookrightarrow \beta_0$ structure on $CE_{red}^*(h)$

shift β_0 get β_0 structure on $CE_{red}^*(h)[-1]$ (when M not compact, some analysis needed)

Analogous result holds for local L_∞ algs

Then: construct central ext of dglas

$$0 \rightarrow CE_{red,loc}^*(L)(M) \xrightarrow{\text{initial}} \text{InnerAct}(L, M) \xrightarrow{\text{dglas}} \text{Act}(L, M) \rightarrow 0$$

$\ker(CE_{red,loc}^*(L)(M) \xrightarrow{\text{initial}} \text{InnerAct}(L, M)) = CE_{red,loc}^*(U(M)(M))$

such that an action L on M corresponds to a Maurer-Cartan elt of $\text{Act}(L, M)$ (dglas) this must have $\text{deg } 1$

then we consider the obstruction to lifting S to a Maurer-Cartan elt in $\text{InnerAct}(L, M)$. We lift S ; this is $dS \in \{\tilde{S}, \tilde{S}\} \in \underbrace{CE_{\text{red}, \text{br}}^+(L)(M)[-1]}_{\substack{\text{in deg 2 here} \\ \Rightarrow \text{deg 1 in} \\ CE_{\text{red}, \text{br}}^+(L)(M)[-1]}}$

note $d\{\tilde{S}, \tilde{S}\} = 0$ as $\tilde{S}$ has deg 1 $\Rightarrow$ defines cocycle $[\alpha] \in H^1(CE_{\text{red}, \text{br}}^+(L)(M))$

really, we may state Noether's thm!

in: Let local obs alg L act on classical field theory M .

Then, $\exists$ map of presheaves of L_0 -algs

$$\hat{L}_c \rightsquigarrow \tilde{\text{Obs}}^c[-1]$$

sending the central elt $1 \in CE[-1] \subseteq \hat{L}_c$ to $1 \in \tilde{\text{Obs}}^c[-1]$.

ea of proof: $\forall U \in M \nabla \text{inj}^{\text{open}} \text{ obtain map}$

$$\text{InnerAct}(L, M) \rightarrow C_{\text{red}}^+(L_c(U)) \otimes \tilde{C}^+(M(U))$$

action $\rightsquigarrow S \in \text{InnerAct}(L, M)$ s.t. $dS \in \{\tilde{S}, \tilde{S}\} = \alpha$

$$\rightsquigarrow S(U) \in C_{\text{red}}^+(L_c(U)) \otimes \tilde{C}^+(M(U))$$

$$\text{cc. } dS(U) + \frac{1}{2} \{S(U), S(U)\} = \alpha(U)$$

$$\rightsquigarrow \hat{L}_c(U) \rightsquigarrow \tilde{\text{Obs}}^c(U)[-1]$$

functionals will cancel the derivative

By construction it's a map of presheaves

, not obvious how this corresponds to our usual notion of Noether's thm.

conservation laws

b: A conserved current is a map of presheaves $J: \bar{\Omega}_c^*[] \rightarrow \tilde{\text{Obs}}^c$
distributional forms

A $N \subseteq M$ codim 1 oriented submfld
then the delta distribution on N is an elt $[N] \in \bar{\Omega}_c^1(U)$ defined on any
neighborhood U of N . (compatible with inclusion $\rightsquigarrow$ genuine elt of $\tilde{\text{Obs}}^c(U)$)

now we have coordinates



Then $dJ[M] = J[N] - J[M'] \in \tilde{\text{Obs}}^c(M)$ $\Rightarrow$ when $dJ[N]$ is
codim 0 submfld. $\Rightarrow$ when $dJ[N]$ is
codimension invariant.

special case: spacetime $N \times I$ $\uparrow t \rightsquigarrow$ usual concept of conserved current



In the usual version of Noether's thm, we have an $\mathbb{R}$ -family of symmetries, so the local Lie alg is Ω_M^* .

In this case, can check that central ext is trivial, so we get

$$\Omega_{M,c}^*[\mathbb{I}] \rightarrow \widetilde{\text{Obs}}^{\mathcal{U}}$$

which for formal reasons gives $\bar{\Omega}_{M,c}^*[\mathbb{I}] \rightarrow \widetilde{\mathcal{O}}_S^{\mathcal{U}}$
(connected current).

Extension to factorisation algs

Define $\mathcal{U}_\alpha^{\text{Po}}(\mathcal{L})(\mathcal{U}) = \text{Sym}^*(\hat{\mathcal{L}}_\mathcal{L}(\mathcal{U})[\mathbb{I}]) \otimes_{\mathbb{C}[\mathcal{C}]} \mathbb{C}_{\mathcal{C}=1}$
 factorisation enveloping Po alg $\mathcal{U}_\alpha^{\text{Po}}(\mathcal{L})$ specialise central elt to 1

Then replace $\hat{\mathcal{L}}_\mathcal{L}$ by this in Noether's thm to get a homomorphism
 of Po FAs $\mathcal{U}_\alpha^{\text{Po}}(\mathcal{L}) \rightarrow \widetilde{\text{Obs}}^{\mathcal{U}}$.

Quantum version

classical field theory $\rightsquigarrow$ quantum field theory

definition is rather involved but it gives
 a BD FA of quantum observables as in previous lecture
 recall: BD-alg is a deformation of Po alg

Get map $\mathcal{U}_\alpha^{\text{BD}}(\mathcal{L}) \rightarrow \widetilde{\text{Obs}}^{\mathcal{U}}$ of BD FAs
 twisted enveloping BD factorisation alg

angular momentum for the scalar theory

first order formulation: $S(p, q) = \int (p, \dot{q}) dt + \frac{1}{2} \int (p, p) dt$

take $M = \mathbb{R}$ (time) $V \cong \mathbb{R}^N$ with inner product $(,)$

canonical field theory is given by

$$0 \rightarrow \begin{matrix} \overset{1}{Q^0=q} \\ \oplus \\ \underset{P^0=p}{C^\infty(\mathbb{R}) \otimes V} \end{matrix} \xrightarrow{d_{\text{od}}} \begin{matrix} \overset{2}{\Omega^1(\mathbb{R}) \otimes V} \\ \oplus \\ \underset{P^1}{\Omega^1(\mathbb{R}) \otimes V} \end{matrix} \xrightarrow{d_{\text{od}}} 0$$

differential: $\begin{matrix} Q \\ P \end{matrix} \longmapsto \begin{matrix} dQ - P dt \\ dP \end{matrix}$

all other L_n are 0

invariant pairing: $\langle Q, P \rangle = \int Q \wedge P$

we recover action $S(Q, P) = \int P \wedge dQ + \frac{1}{2} \int P \wedge P \wedge dt$
which gives eqⁿ above when restricted to fields in deg 0?

symmetries: $L = \Omega^1 \otimes \square(V)$

L action is given by the Maurer-Cartan elt

$$J^{\text{eq}}(p, q, \alpha) = \int (p, dQ) + \int (p, \alpha \cdot Q) + \frac{1}{2} \int (p, p) dt$$

$$\in \text{InnAct}(L, M) = \mathcal{C}_{\text{red, loc}}^+(L \otimes M) / \mathcal{C}_{\text{red, loc}}^*(M)$$

intuitively: $0 \rightarrow M \rightarrow L \ltimes M \rightarrow L \rightarrow 0$

$\forall \alpha \in L$, look at pre-image. This functional on Q, P defines L as structure on the pre-image.



Conserved current: $J: L_c(\mathbb{R}) \rightarrow \underbrace{\widetilde{\text{Obs}}^d[-1]}_{\text{recall this is } \approx \mathcal{C}^*(M(-1))}$

$$\alpha \longmapsto \int \langle p, \alpha \cdot Q \rangle$$

put in $\alpha = a \cdot \delta_t$: $\langle p(t), a \cdot q(t) \rangle$ constant which is conservation of angular momentum.