

Basal processes and geomorphology



Sediments and sliding

- Till rheology
- Deformation

Drainage in sediments

- Darcy flow
- Canals

Geomorphology

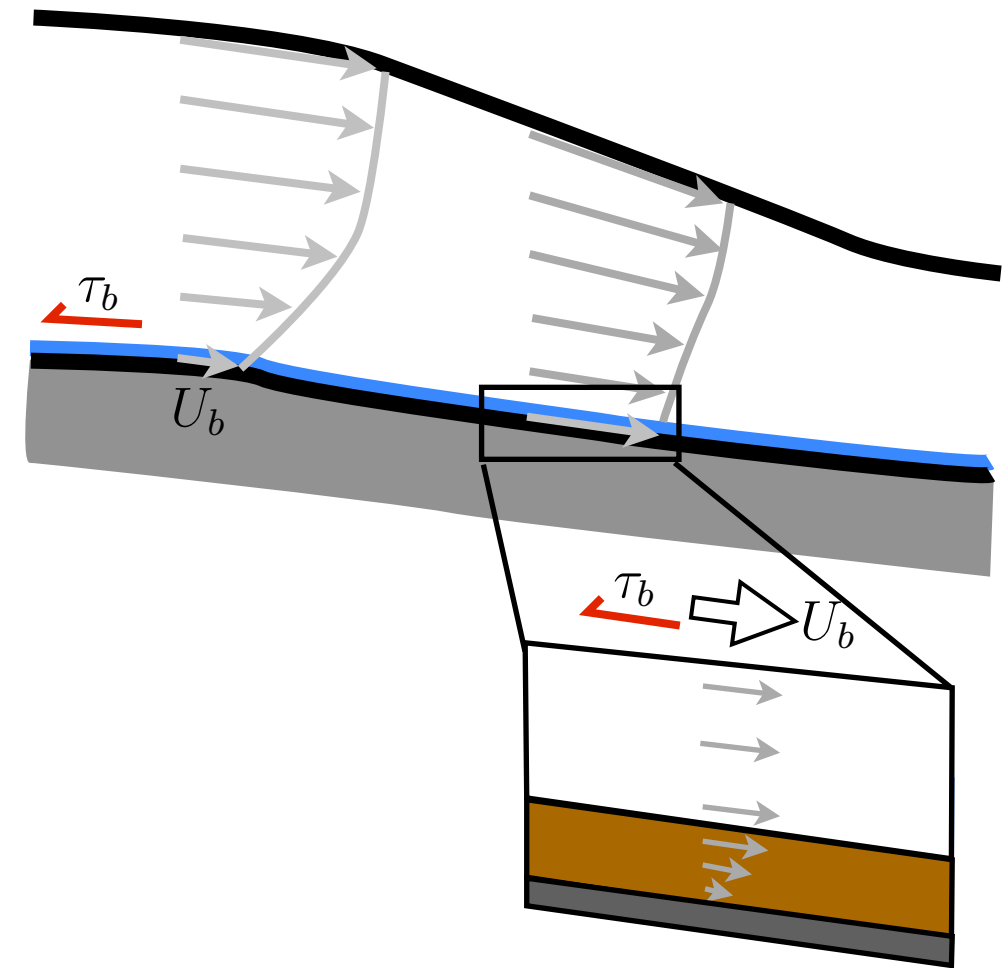
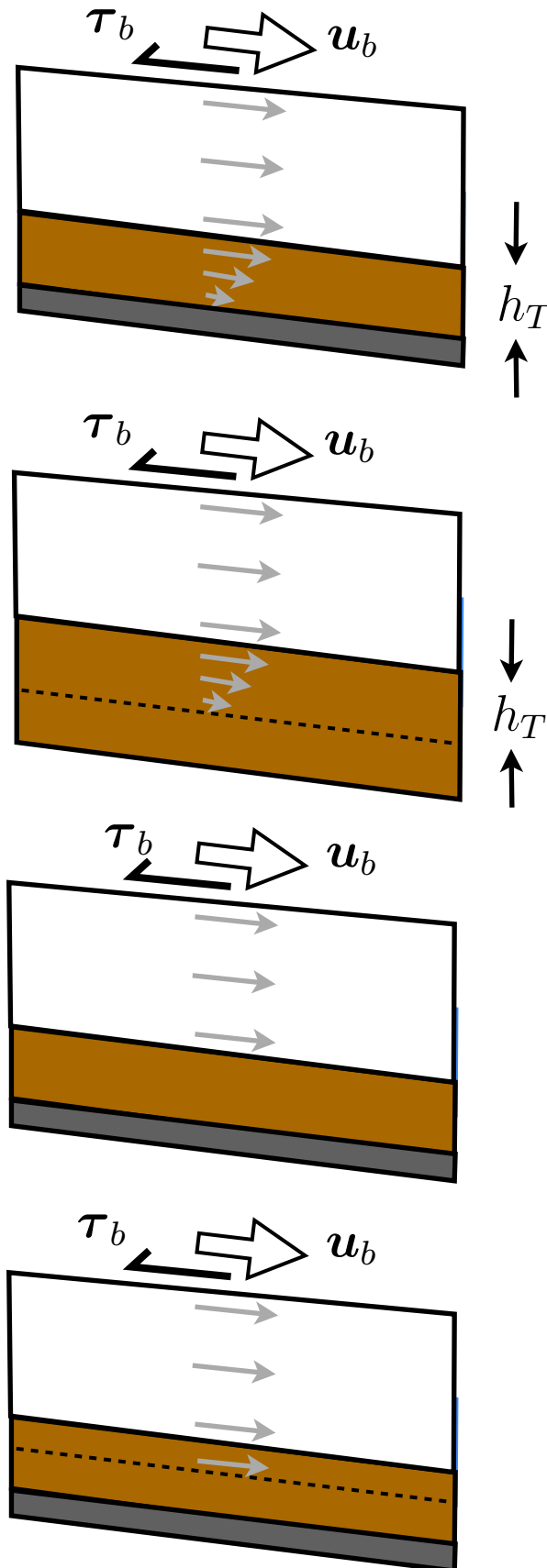
- Meltwater deposits
- Deformational deposits

Sediments and sliding

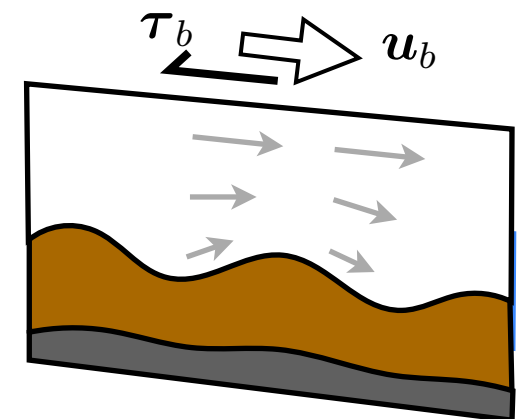
Sliding over sediments

‘Sliding’ could involve:

- Shear deformation of sediment layer
- Shear of a finite horizon of the sediment
- Slip at the ice-till interface
- Slip on slip-planes within the sediment layer



Macroscopic resistance may come from flow around sediment landforms



Till rheology

Hooke & Iverson 1998, Kamb 1991, Iverson 2011, Iverson & Zoet 2015

Laboratory experiments on samples show that till has a **yield stress**

$$\tau_f = c_0 + \mu\sigma_e$$

Yield stress depends on **effective stress**

$$\sigma_e = P - p_w \approx N$$



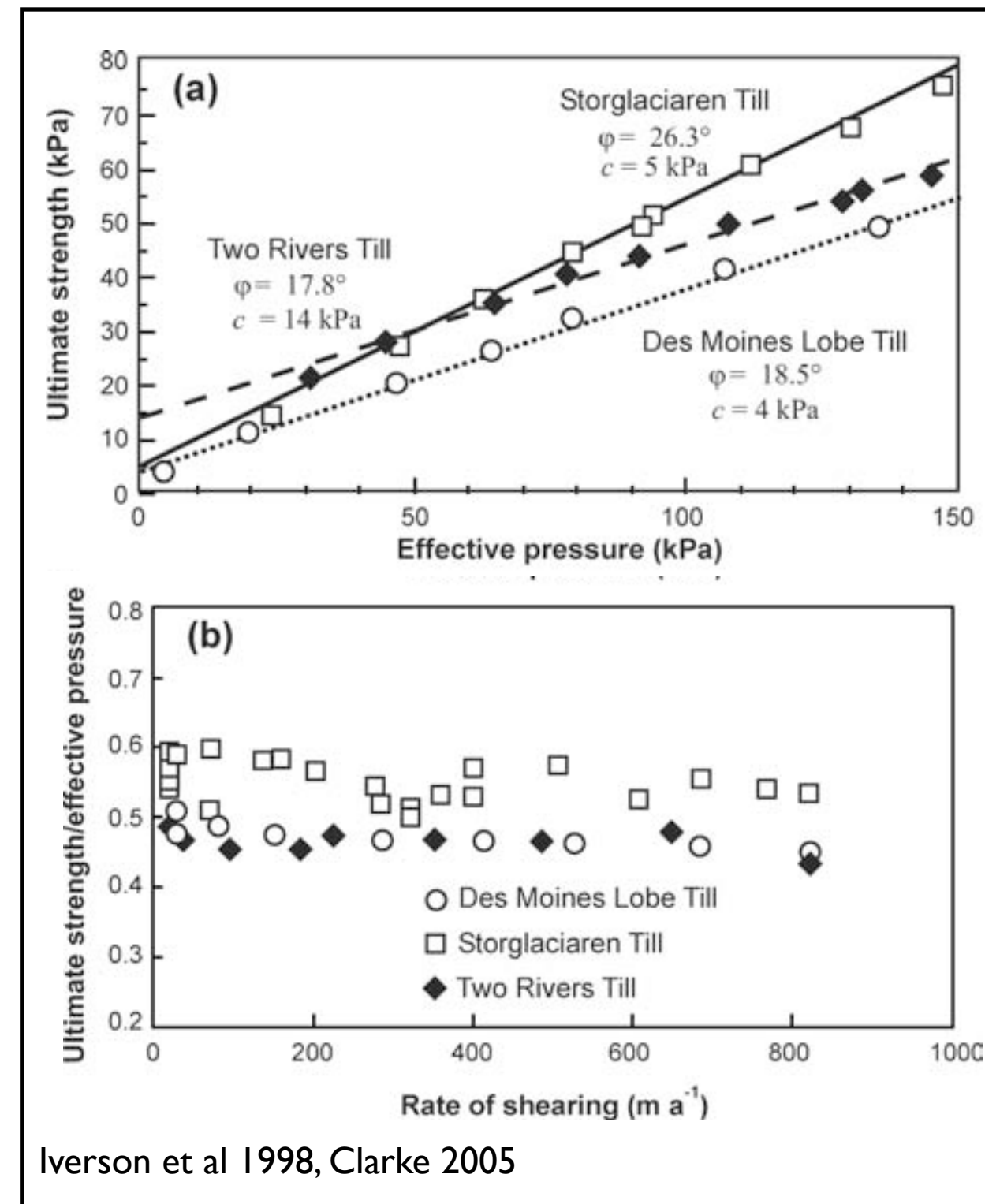
Effective pressure at ice-till interface

(effective stress increases with depth into the till - it is weakest at the top).

$$\mu = \tan \psi \approx 0.4 \quad \text{Coefficient of friction}$$

$$c_0 \approx 3 \text{ kPa} \quad \text{Cohesion}$$

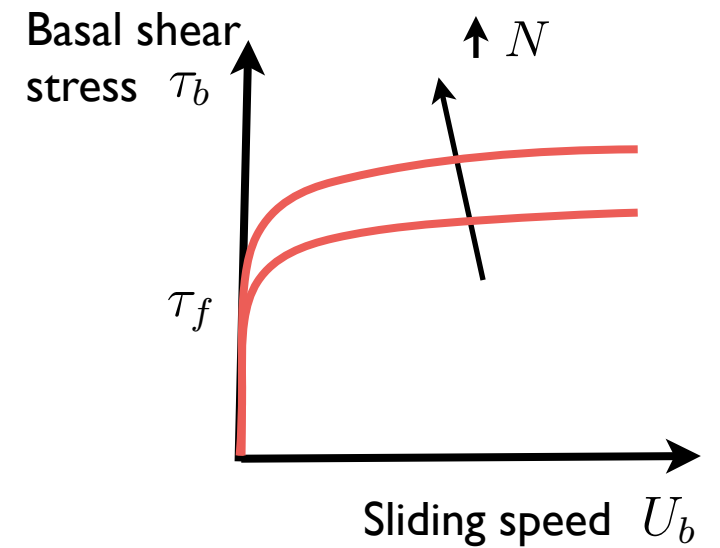
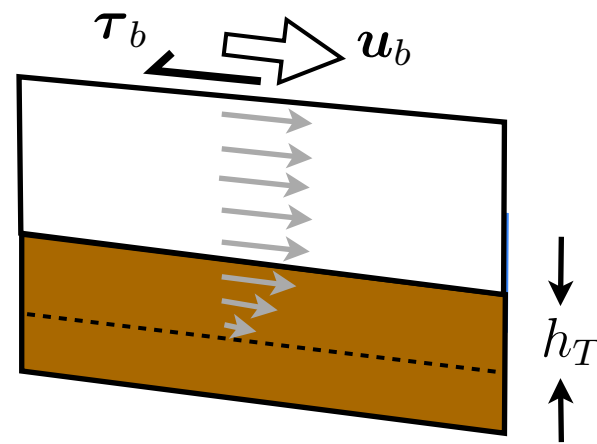
Experiments suggest **stress is almost independent of strain rate** (i.e. perfectly plastic).



Sliding over till

Viscous rheology

$$\dot{\epsilon} = A(\tau - \tau_f)^a \sigma_e^{-b} \quad \tau \geq \tau_f = \mu \sigma_e$$



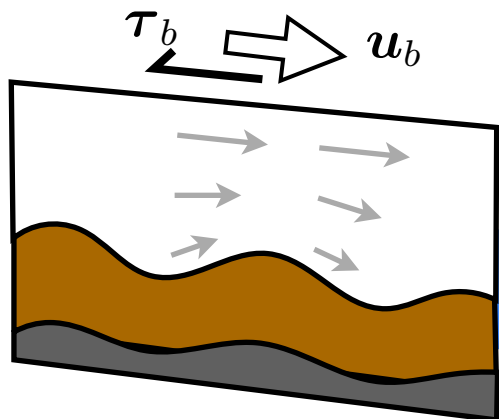
Pore water pressure roughly hydrostatic

⇒ Effective stress increases with depth through till $\sigma_e = N + \Delta\rho_{sw}g(Z_b - z)$

Deformation only if $\sigma_e \leq \tau_b/\mu$ ⇒ Deforming horizon $h_T = [\tau_b - \mu N]_+ / \mu \Delta\rho_{sw}g$

⇒ Sliding law

$$\tau_b = \mu N + C U_b^{1/a} N^{b/a}$$

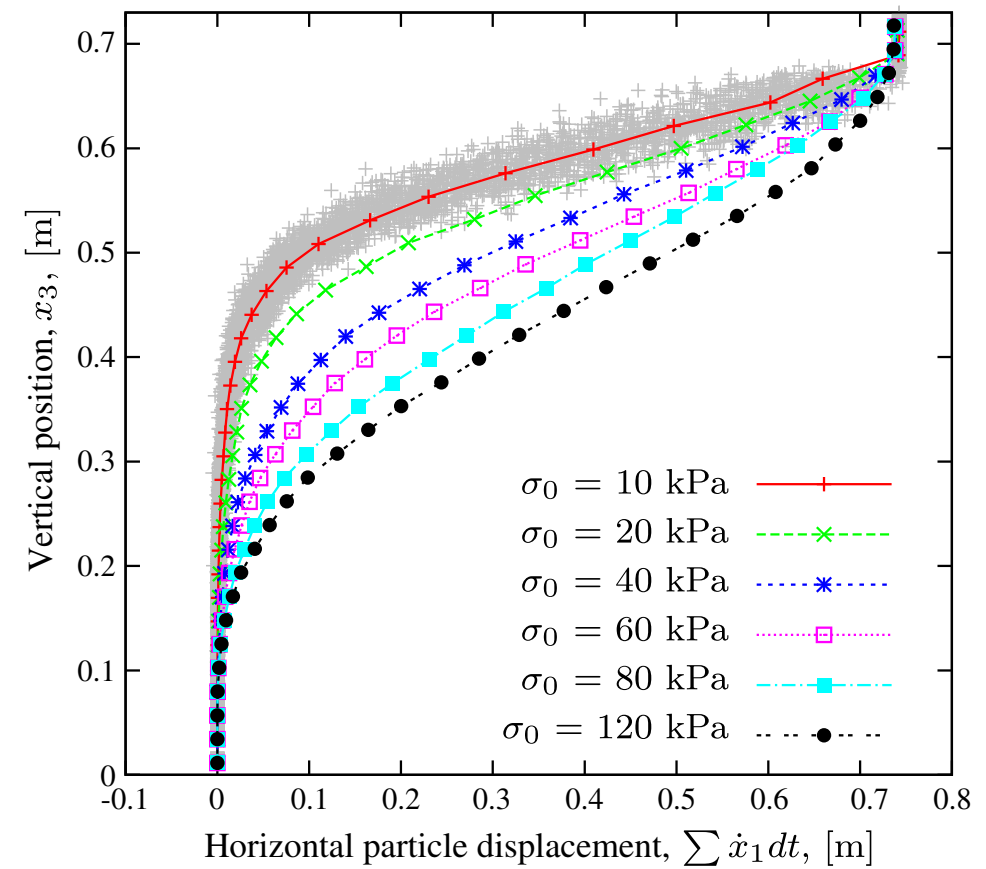
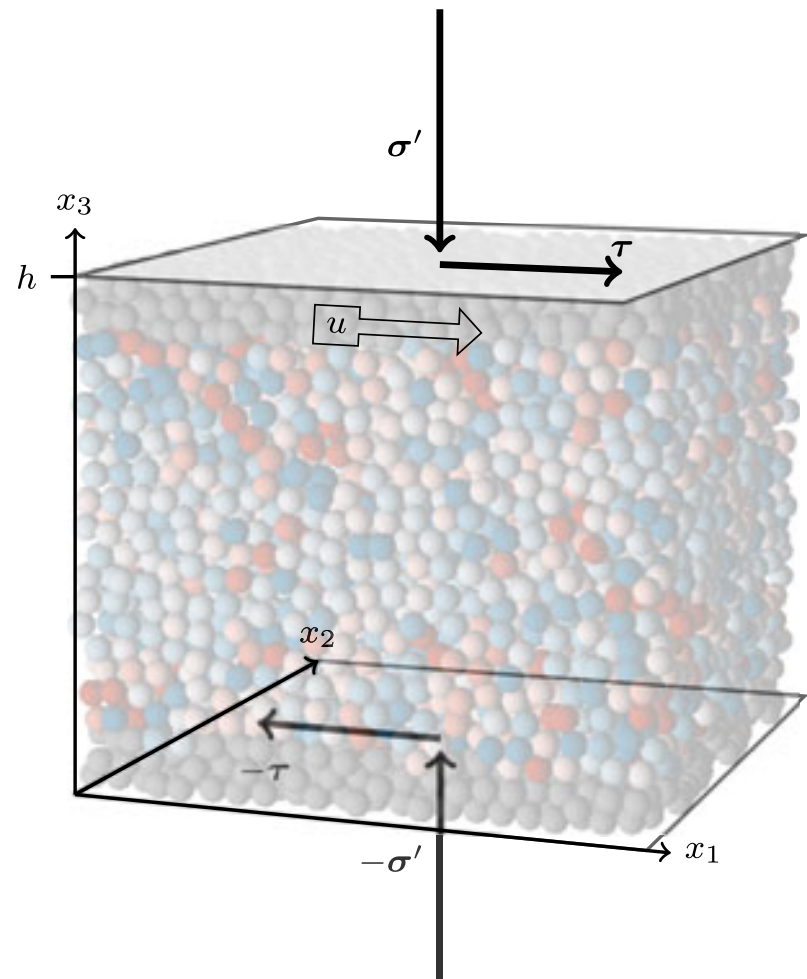


A similar law applies to describe ice flow over sediments with topography

$$\tau_b = \mu N + R U_b^{1/n}$$

Computational experiments

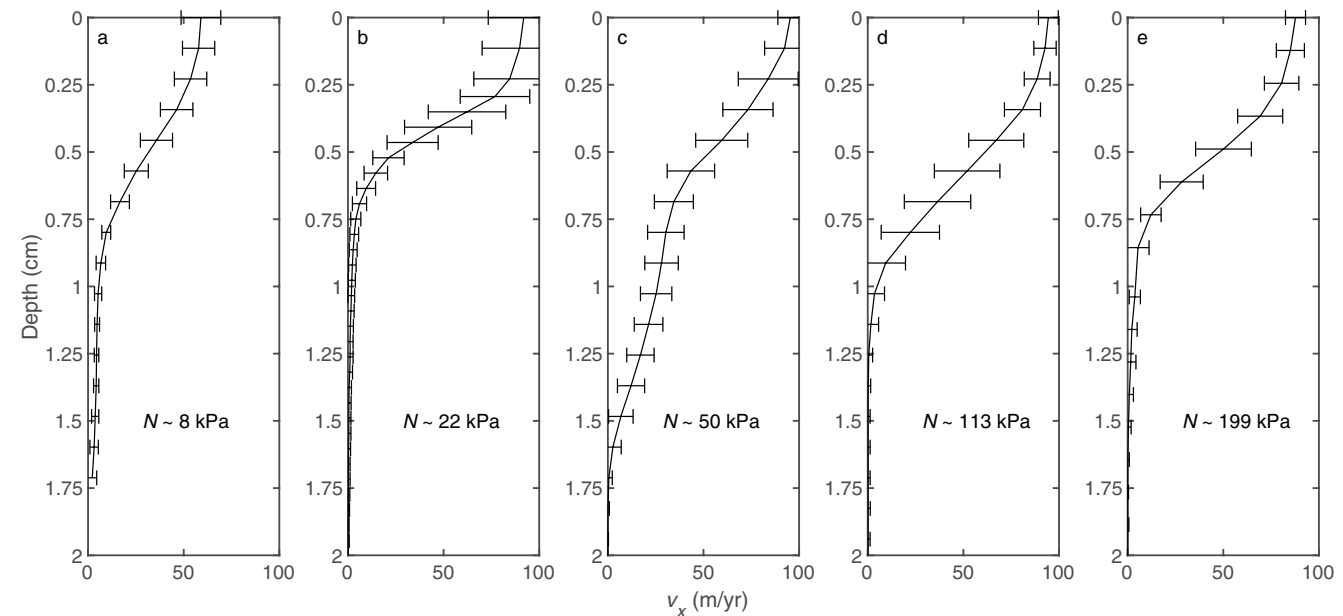
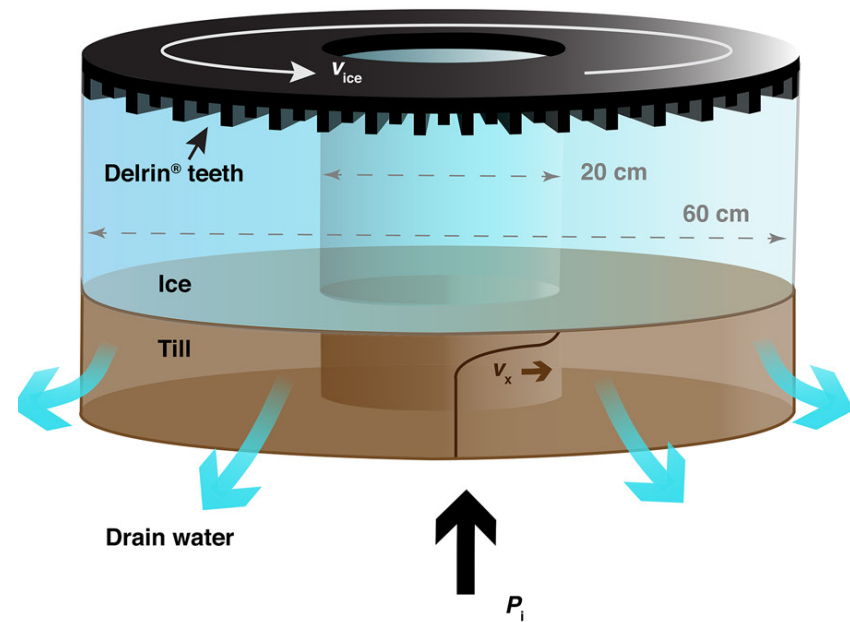
Discrete particle (DEM) experiments under imposed shearing velocity



Damsgaard et al 2013

Laboratory experiments

Laboratory ring shear experiment visualise till deformation, sediment flux, and ice-till slip

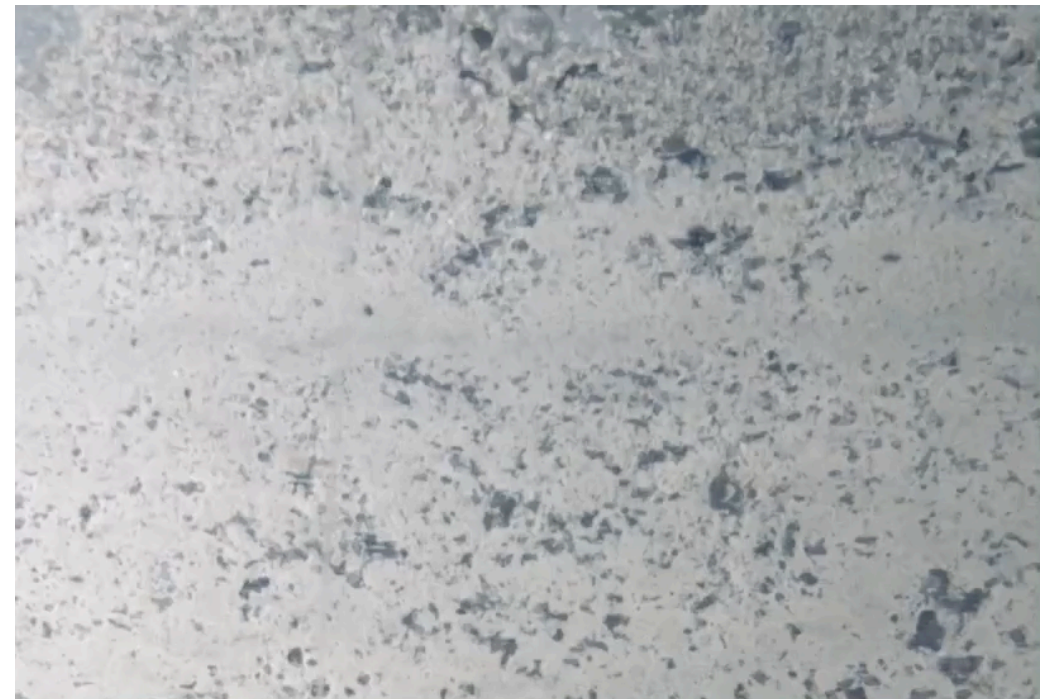


Hansen & Zoet 2022

Ice-till slip occurs at low effective pressure / low sliding speeds.

Depth of deformation increases then decreases with effective pressure.

Sediment flux scales approximately **linearly with sliding speed**, and **non-monotonically** with effective pressure.



Drainage in sediments

Drainage through till

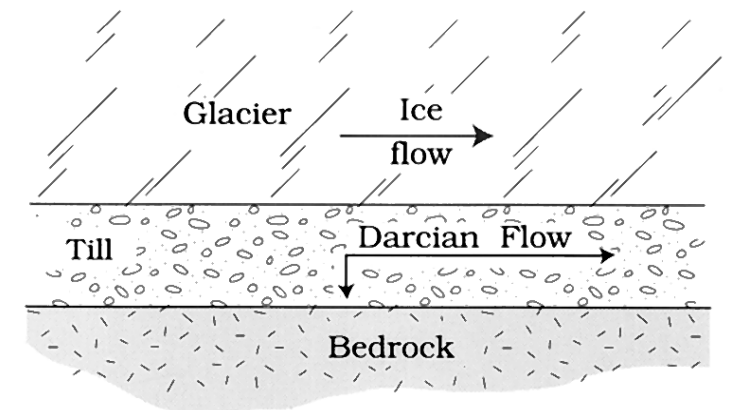
Estimates of **hydraulic conductivity** vary, but it is generally thought to be low.

Although water seeps **vertically** into the till, **horizontal** transport **through** the till is most likely **insufficient** to evacuate the water produced from melting.

eg. in Antarctica

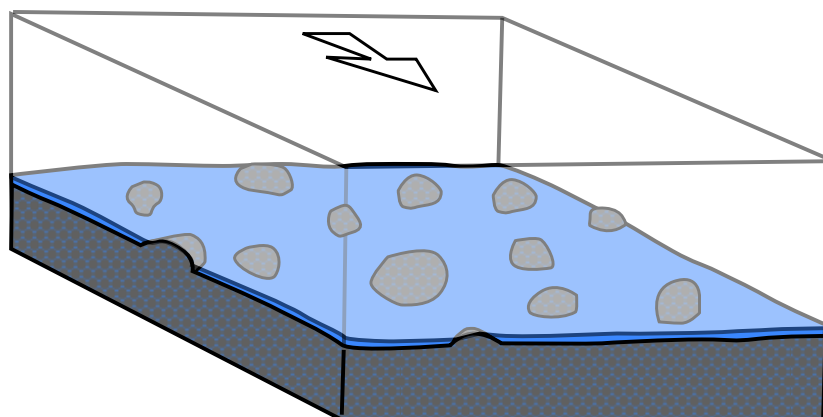
$$q = \frac{Kh_T}{\rho_w g} \nabla \phi \approx \frac{10^{-7} \cdot 10}{10^3 \cdot 10} \cdot 10 \frac{\text{m s}^{-1} \text{ m}}{\text{kg m}^{-3} \text{ m s}^{-2} \text{ m}} \text{ Pa} \approx 10^{-10} \text{ m}^2 \text{ s}^{-1}$$

$$\int m dx \approx 5 \cdot 10^3 \text{ mm y}^{-1} \text{ km} \approx 1.6 \times 10^{-4} \text{ m}^2 \text{ s}^{-1}$$



Fountain & Walder 1998

⇒ Water flows in a **patchy film** at the ice-till interface, or in some form of **channels or canals**.

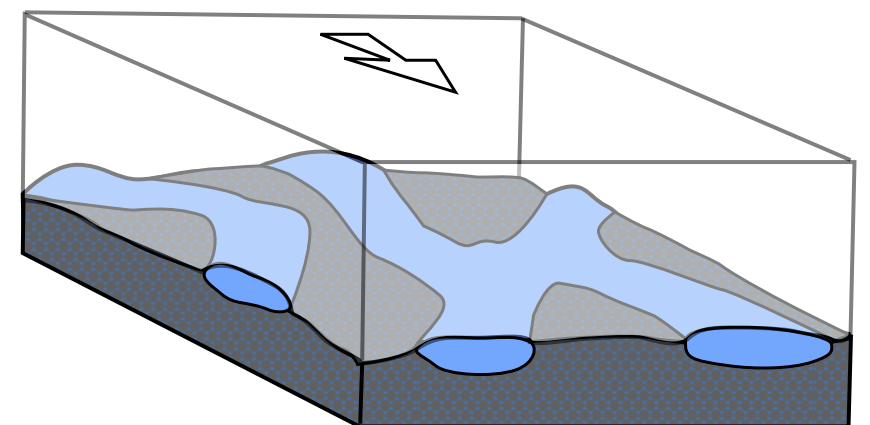


Patchy sheet

Alley 1989, Creyts & Schoof 2009

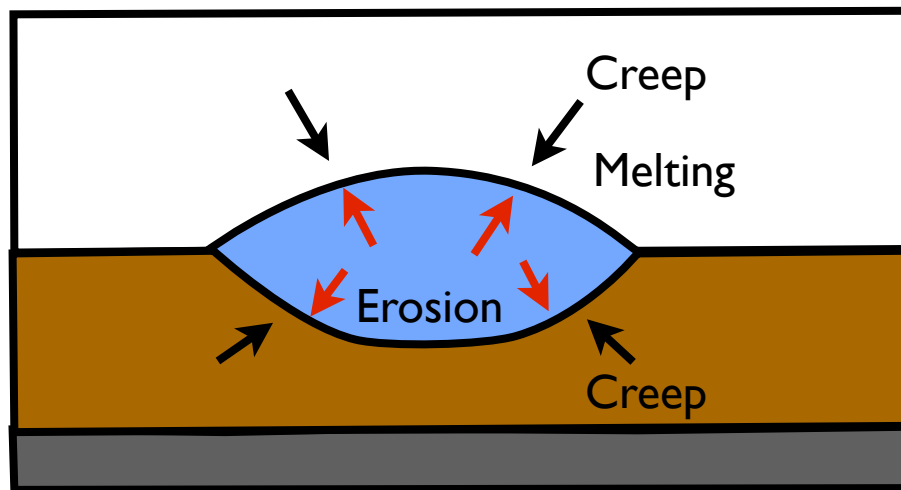
Canals

Walder & Fowler 1994



Canals

Walder & Fowler 1994, Ng 2000



Gravitational potential gradient

$$\Psi = \rho_i g \tan \alpha + (\rho_w - \rho_i) g \tan \theta$$

Walder & Fowler suggested two possibilities for steady states:

Channels - mostly melted into ice

$$N \propto \Psi^{7/15} Q^{1/15}$$

$$N > \tilde{N}$$

$$\tilde{N} \approx 0.8 \text{ MPa}$$

Canals - mostly eroded into sediment

$$N \propto \Psi^{-1/3} Q^{-1/3}$$

$$N < \tilde{N}$$

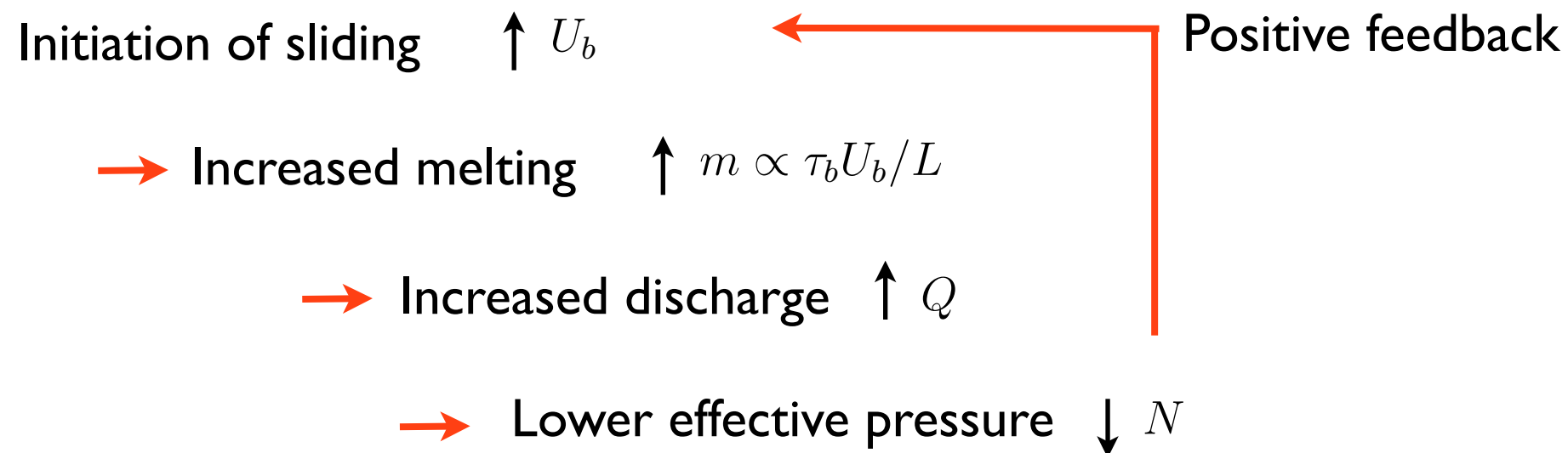
⇒ Effective pressure in canals **DECREASES** with increasing discharge Q
(so might expect distributed system)

The crucial difference seems to be that erosion tends to produce a **wide cross-section**.

Canals are favoured when the **potential gradient is small** (e.g. interior of ice sheets).

Interaction of sliding and drainage

A consequence of $\frac{\partial N}{\partial Q} < 0$ is the potential for a **positive feedback**

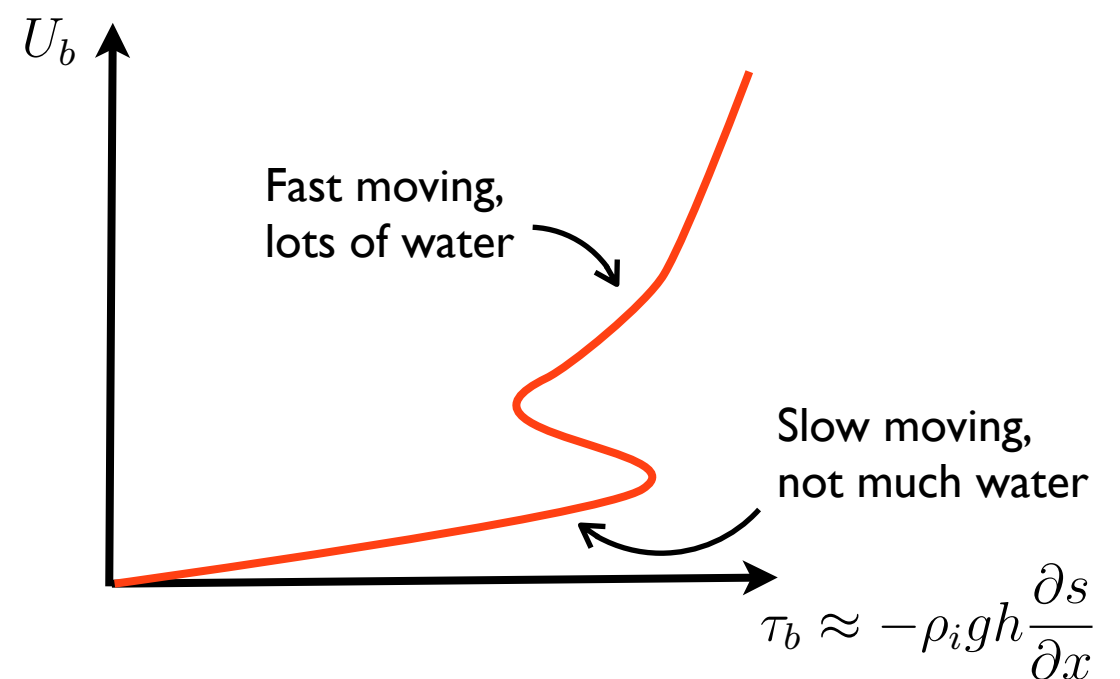


Model

$$Q = \frac{G + \tau_b U_b - k U_b^{1/2}}{\rho_w L} A$$

$$N = c / Q^{1/3}$$

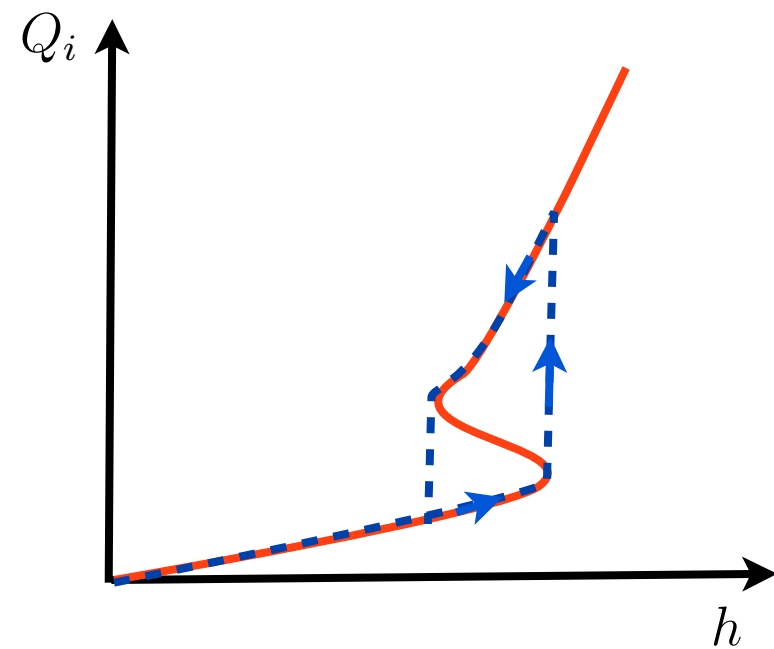
$$\tau_b = C U_b^p N^q$$



The relationship between ice thickness and speed can become **multivalued**

Surges and ice streaming

Temporal variability



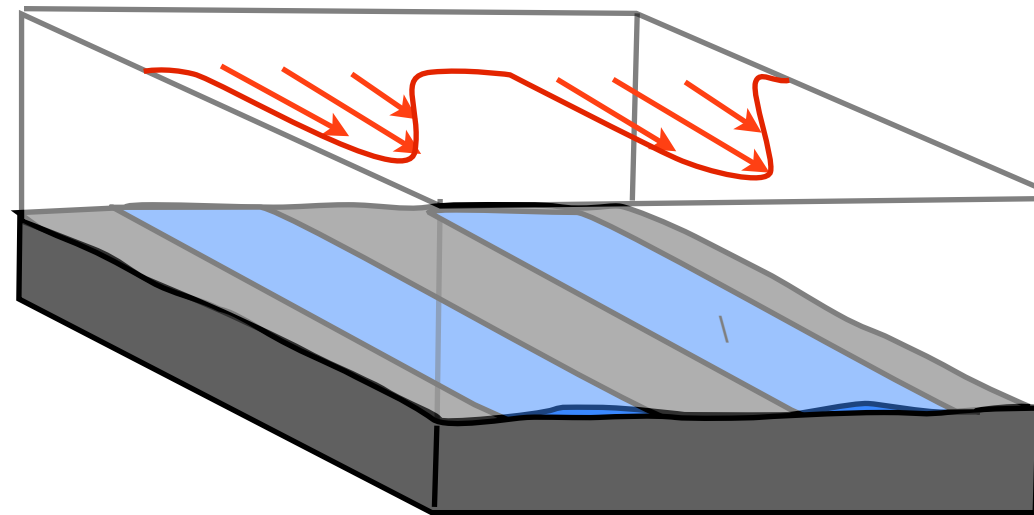
$$\frac{dh}{dt} \propto A - Q_i$$

accumulation

ice flux

⇒ Surges

Spatial variability



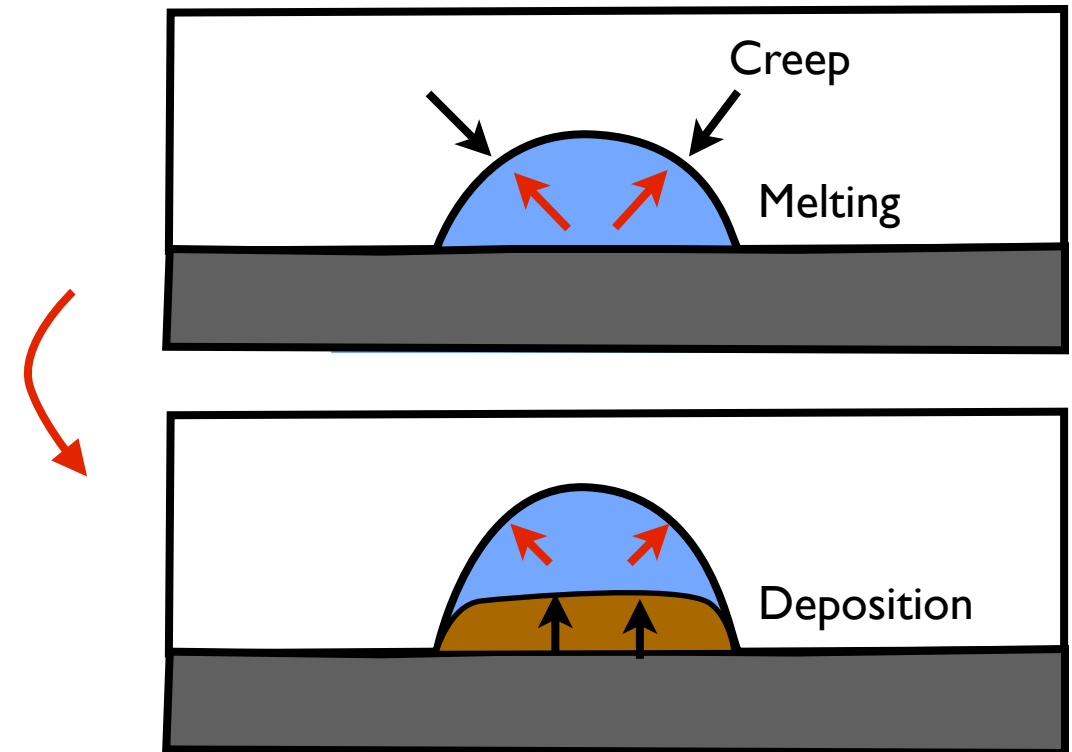
⇒ Ice streams

Meltwater deposits

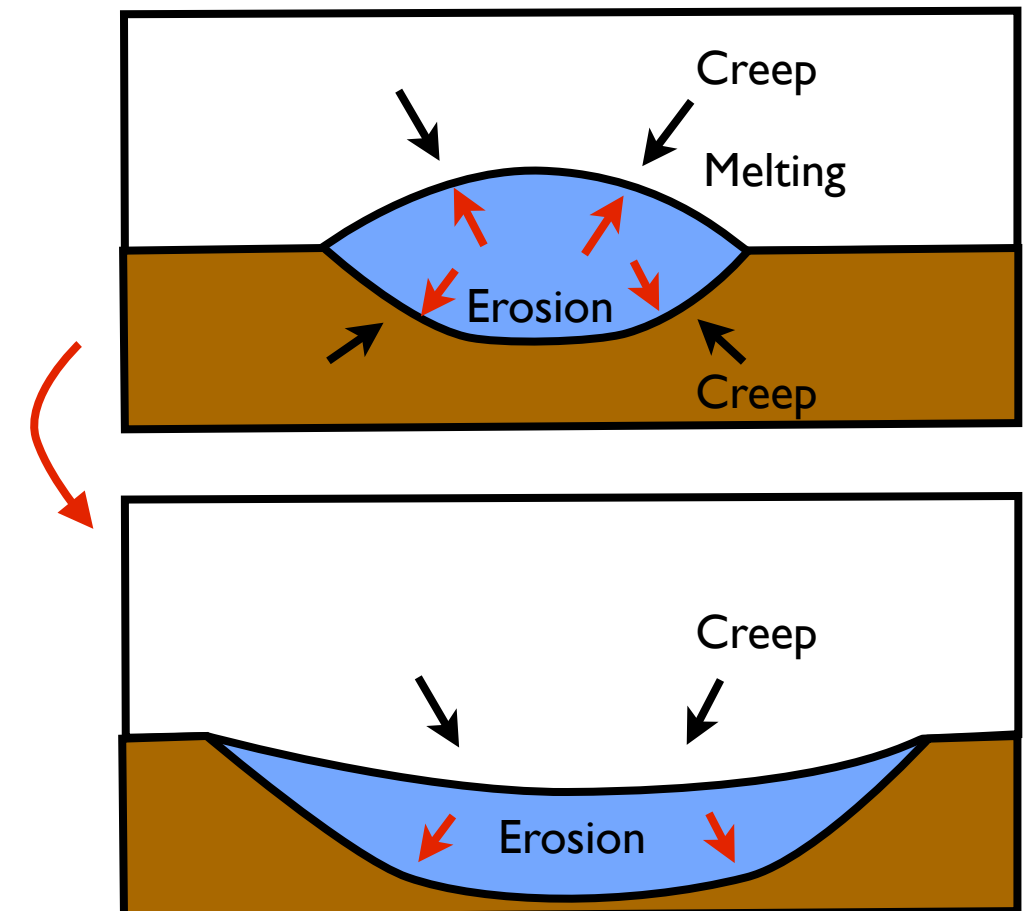
Meltwater deposits

Deposition of sediments in Röthlisberger channels can build **eskers**

- Most likely under falling water speed, near margin
- Sediment is flushed from the surrounding bed

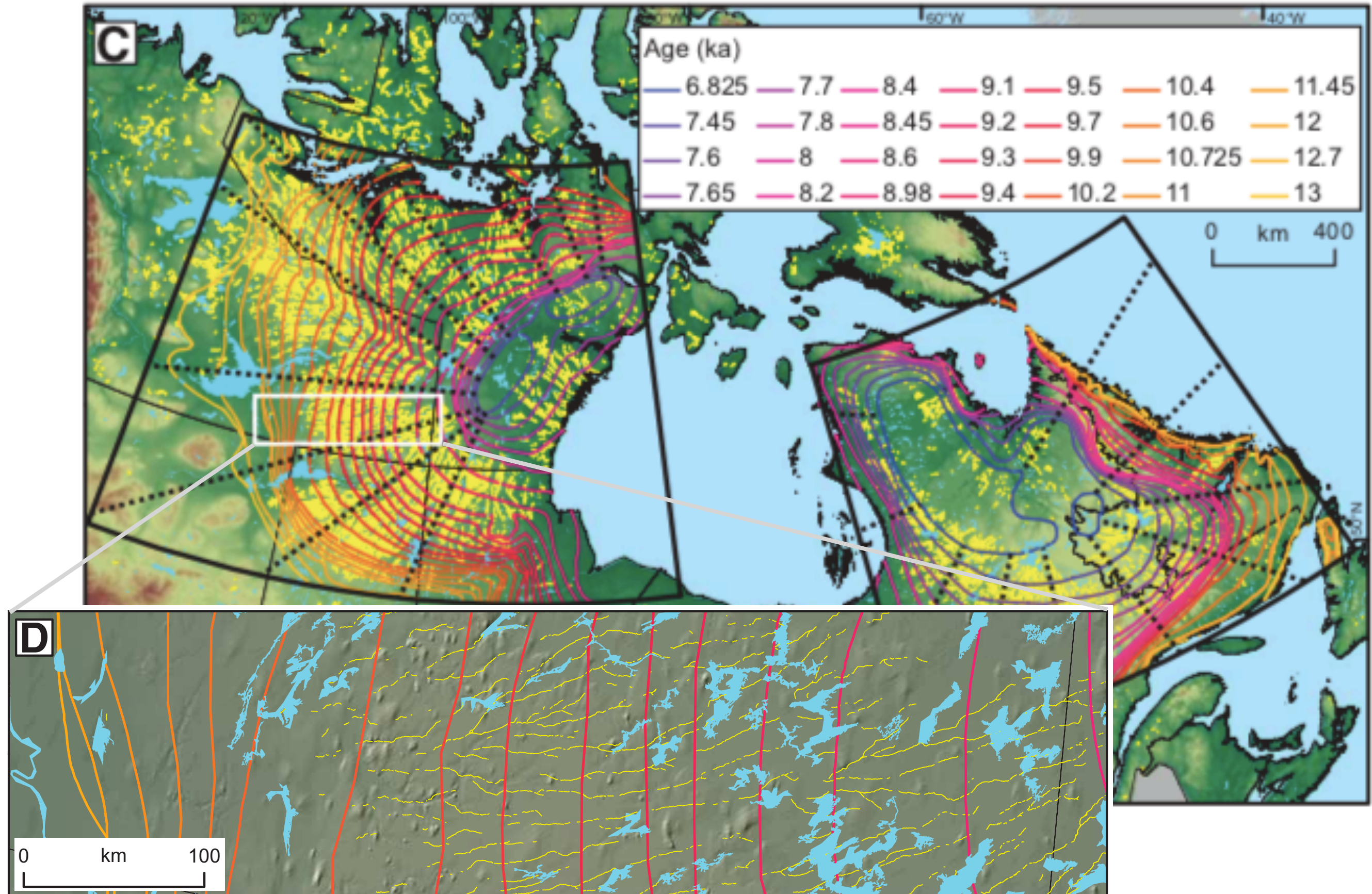


Erosion of sediments from canals can create **tunnel valleys**

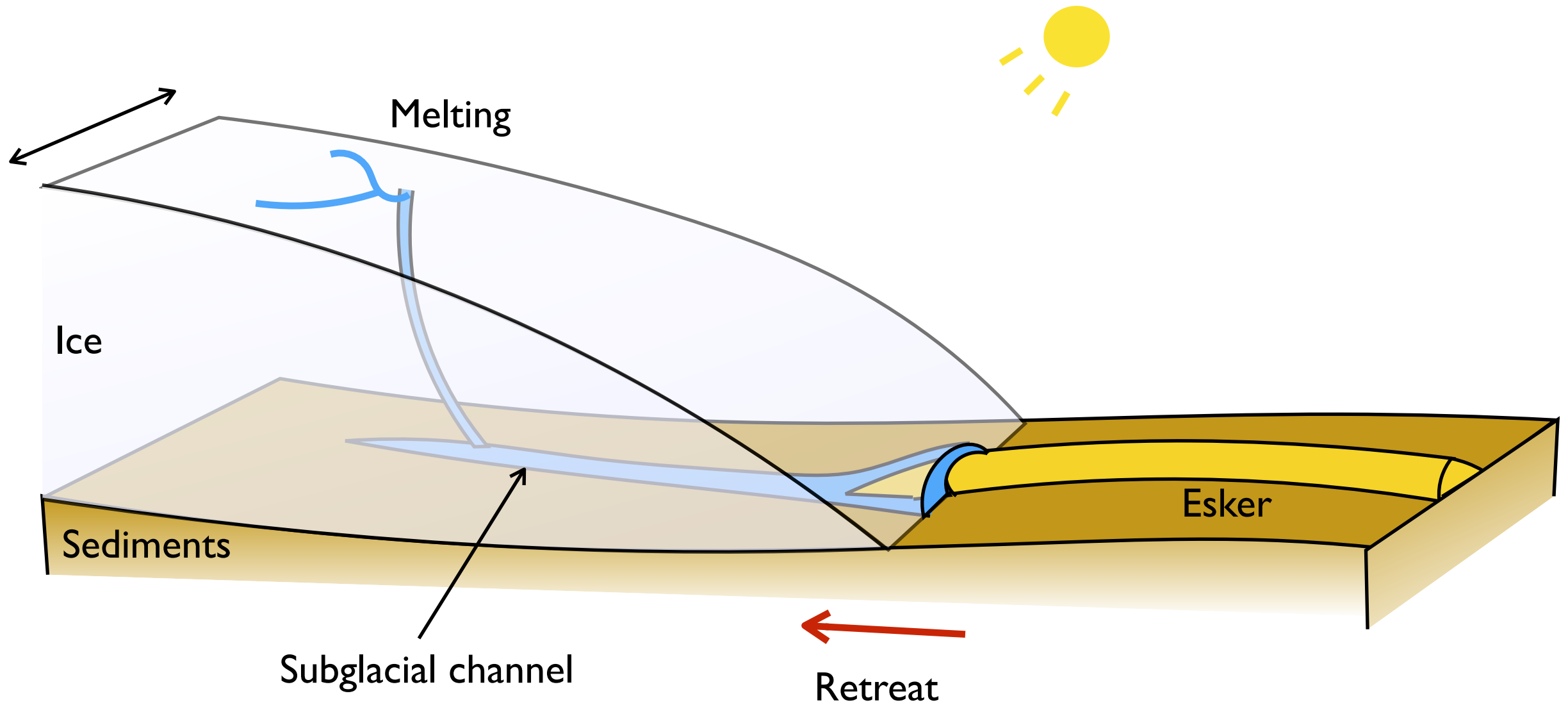




Eskers



Eskers

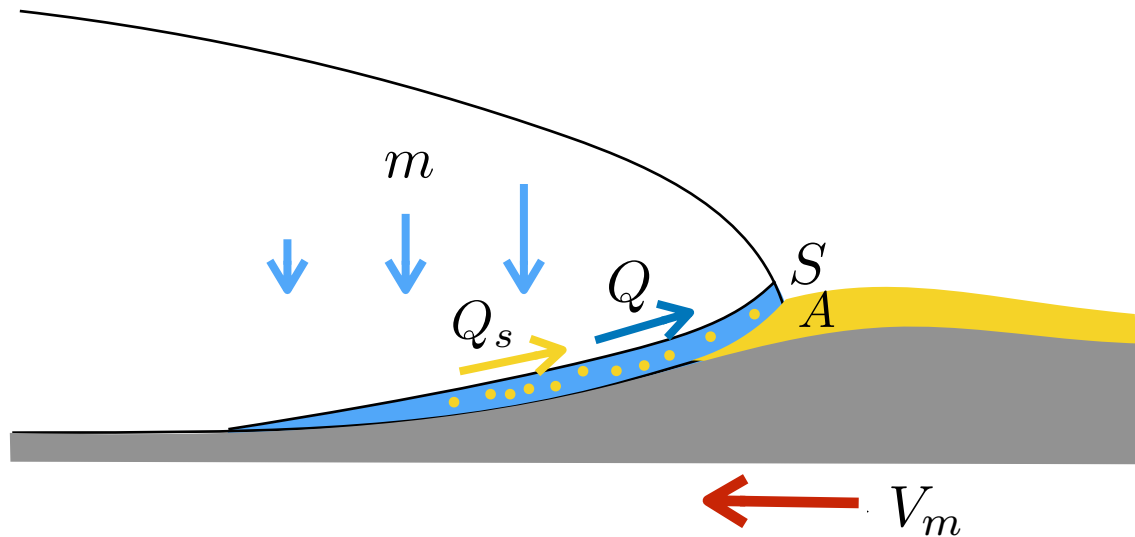


Model

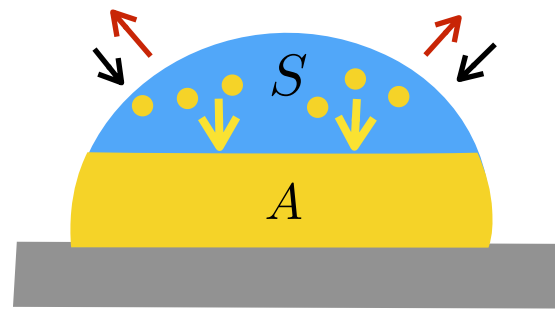
An extended version of the Röthlisberger channel model that incorporates sediment transport.

Eskers

Sediment **deposition** acts to clog the channel:



Hewitt & Creyts 2019



$$\frac{\partial S}{\partial t} = \frac{Q\Psi}{\rho_i L} - \tilde{A}SN^n - D$$

$$\frac{\partial A}{\partial t} = D$$

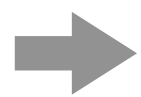
Boundary-layer analysis leads to an approximation of the cross-sectional area:

$$A = C Q_m^{-4/5} Q_{sm}^{29/15} V_m^{-1}$$

Meltwater flux

Sediment flux

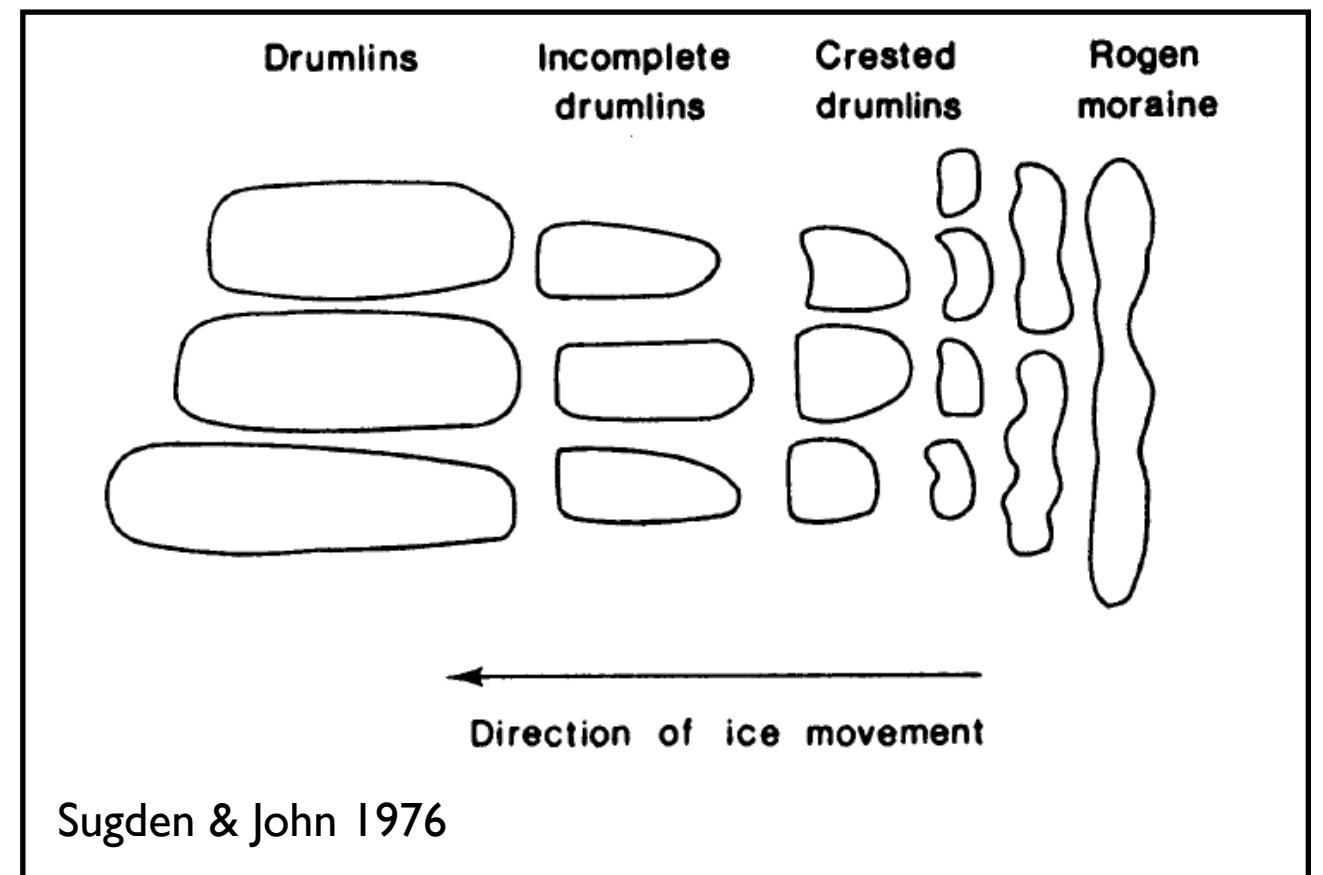
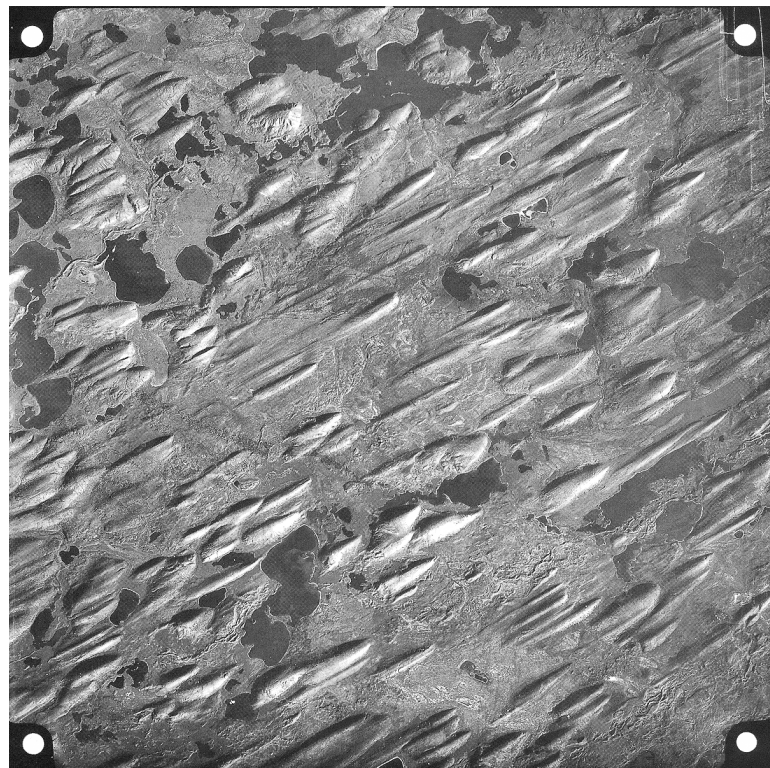
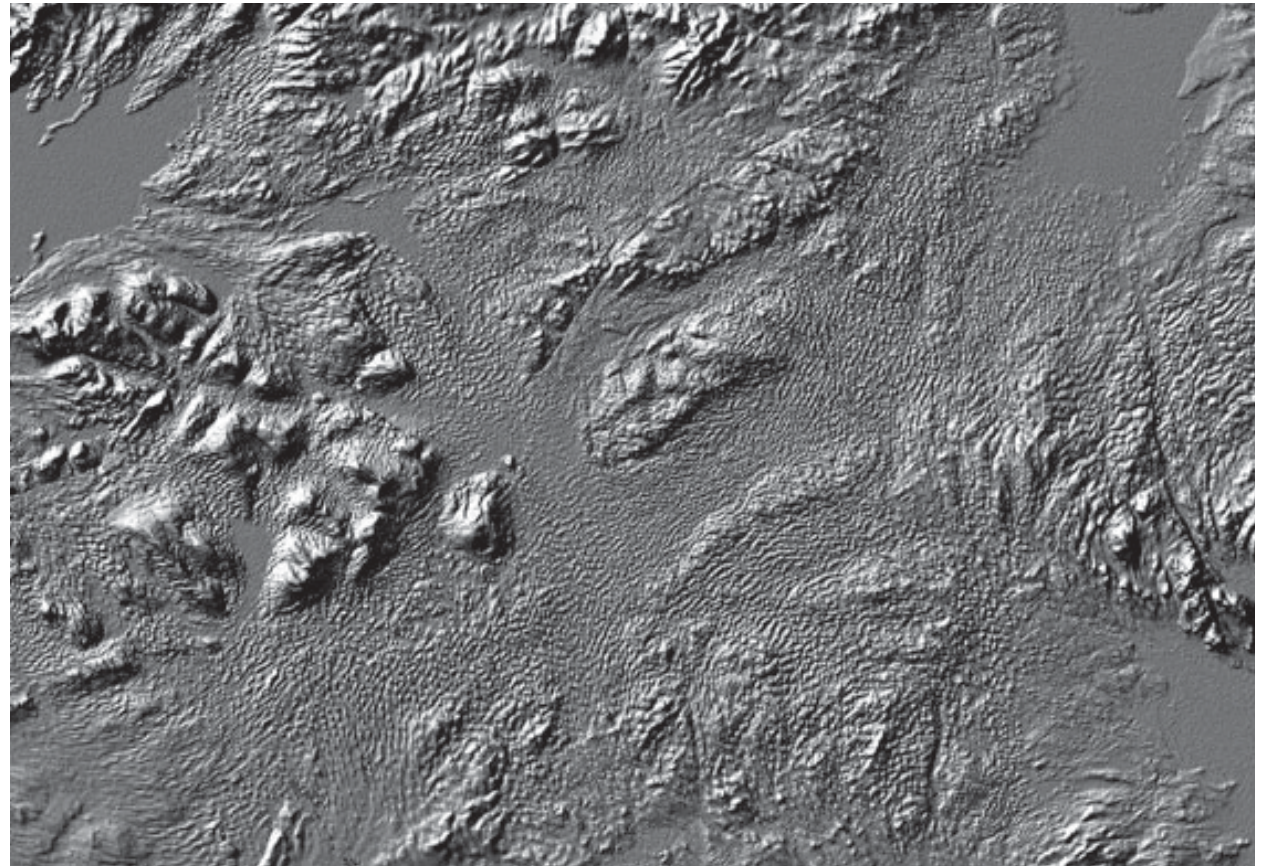
Retreat rate



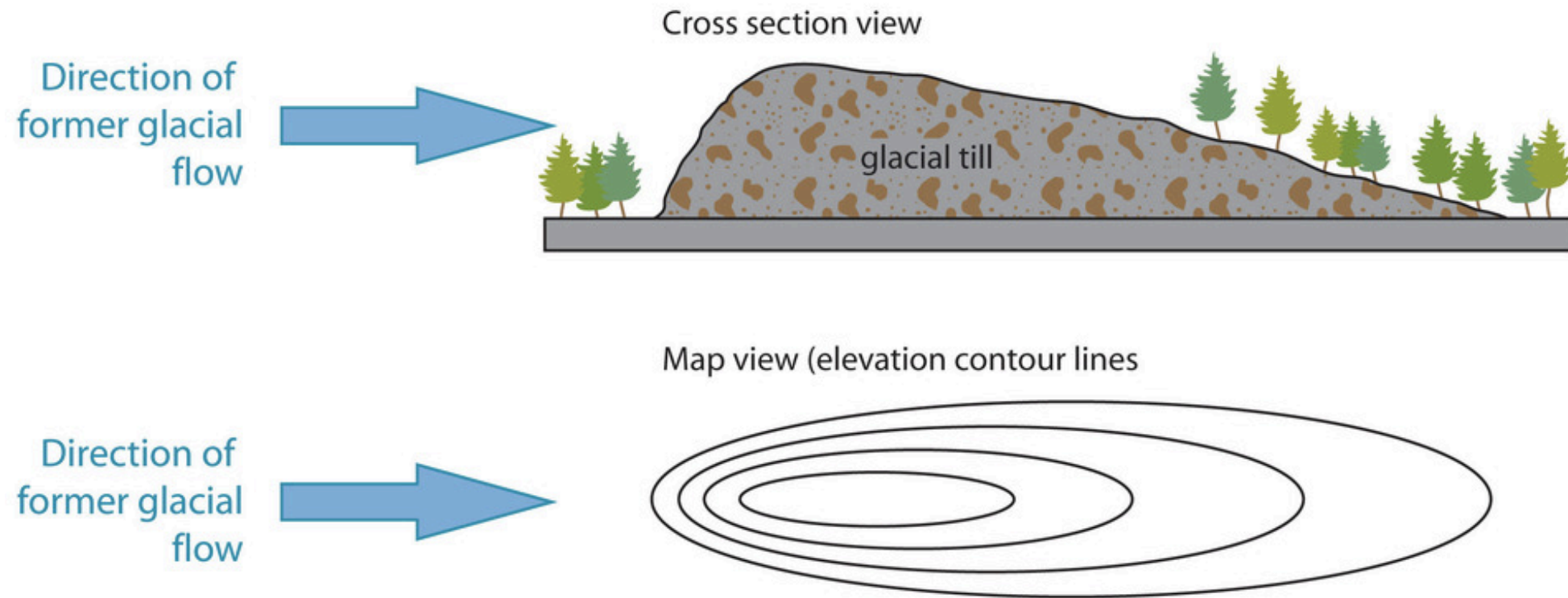
Esker size **increases** roughly quadratically with **sediment flux**, **decreases** roughly linearly with **meltwater flux**, and **decreases** linearly with **ice-sheet retreat rate**.

Deformational deposits

Subglacial bedforms



Drumlins

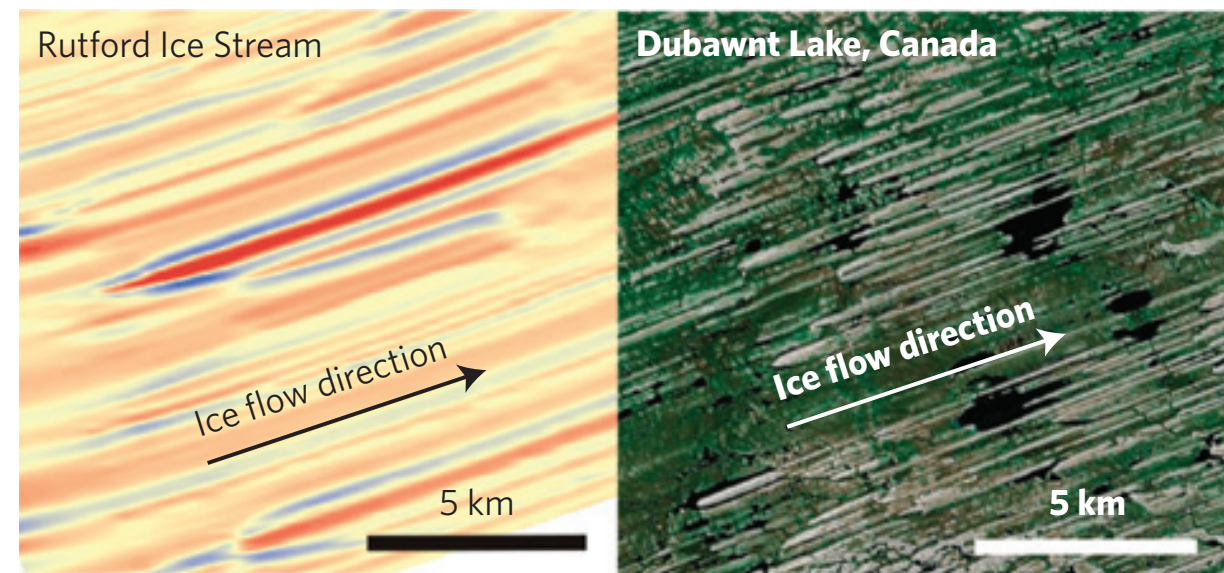
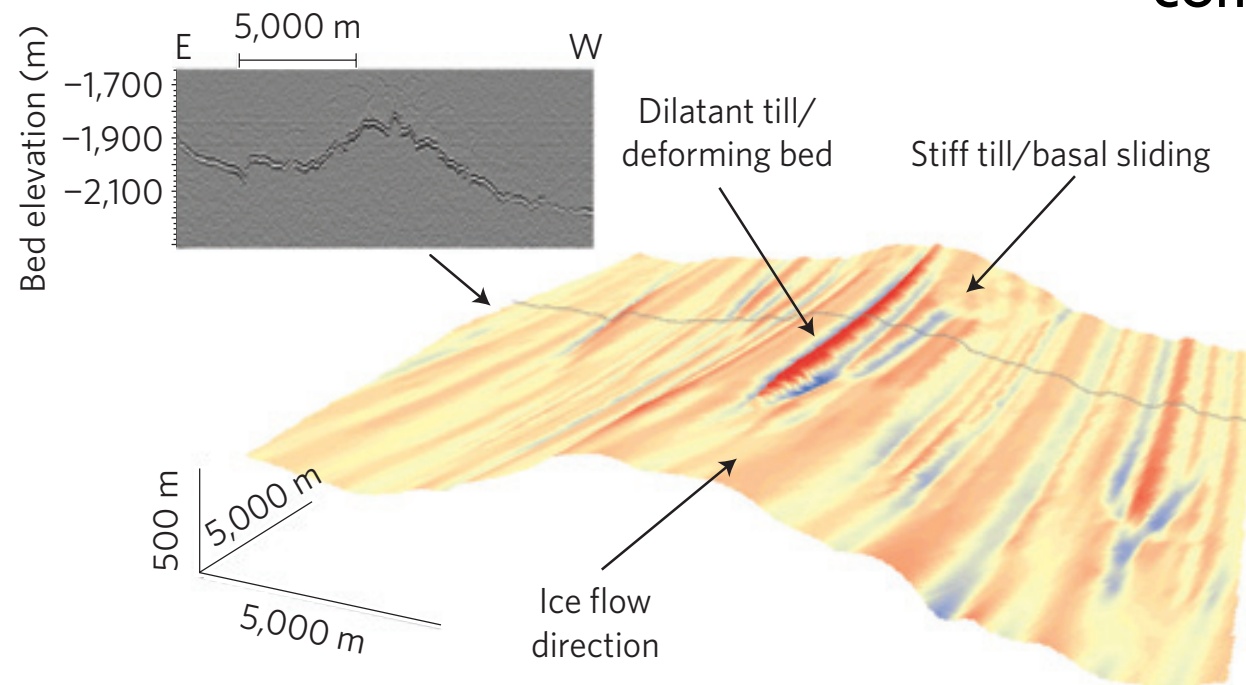


Subglacial waveforms formed by deformation/erosion by ice flow

Mega-scale glacial lineations

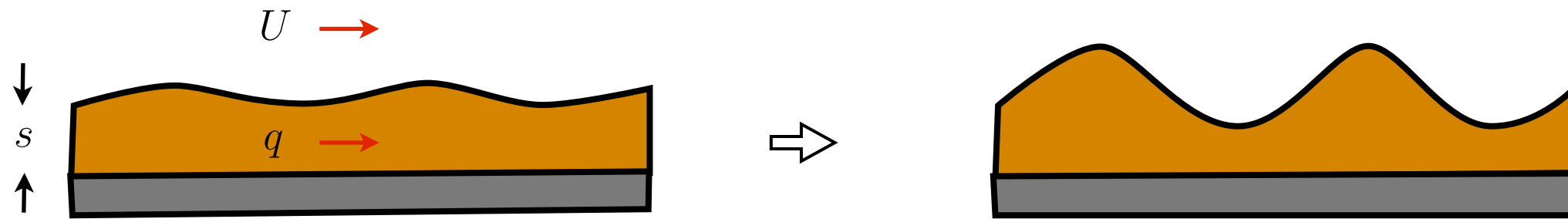
Radar profile of bed beneath Rutford ice stream (West Antarctica)

compared with de-glaciated bed of former ice stream.



King et al 2009

Instability theory Hindmarsh 1998, Schoof 2007, Fowler 2000, 2001, 2002, 2009, 2010, 2010

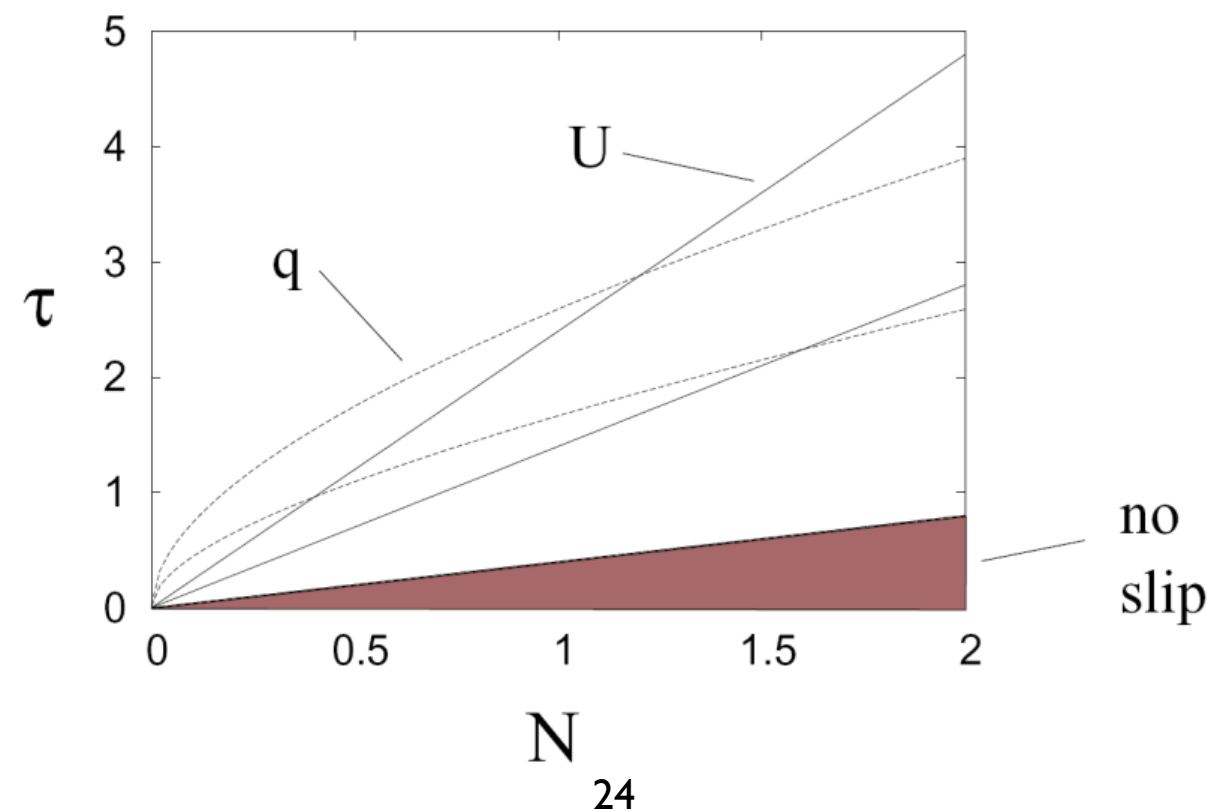


Conservation of sediment $\frac{\partial s}{\partial t} + \frac{\partial q}{\partial x} = 0$

Sediment flux $q = q(\tau, N)$

Sliding law $\tau = \tau(U, N)$

Instability if $\left. \frac{\partial \tau}{\partial N} \right|_U > \left. \frac{\partial \tau}{\partial N} \right|_q$



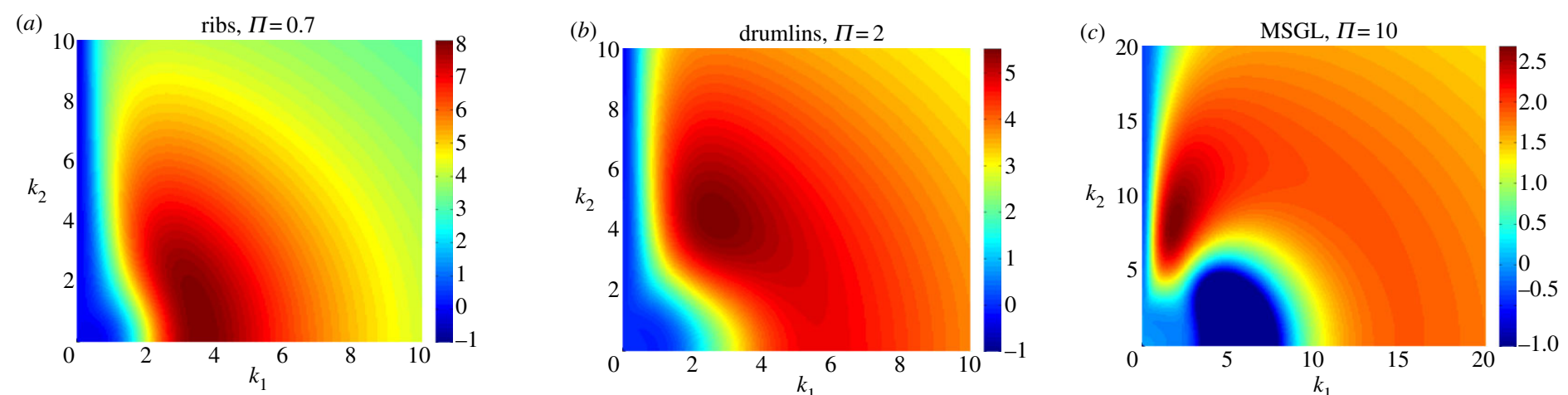
Instability theory Fowler & Chapwanya 2014

Modified theory may explain evolution of ribbed moraine, drumlins, and mega-scale lineations

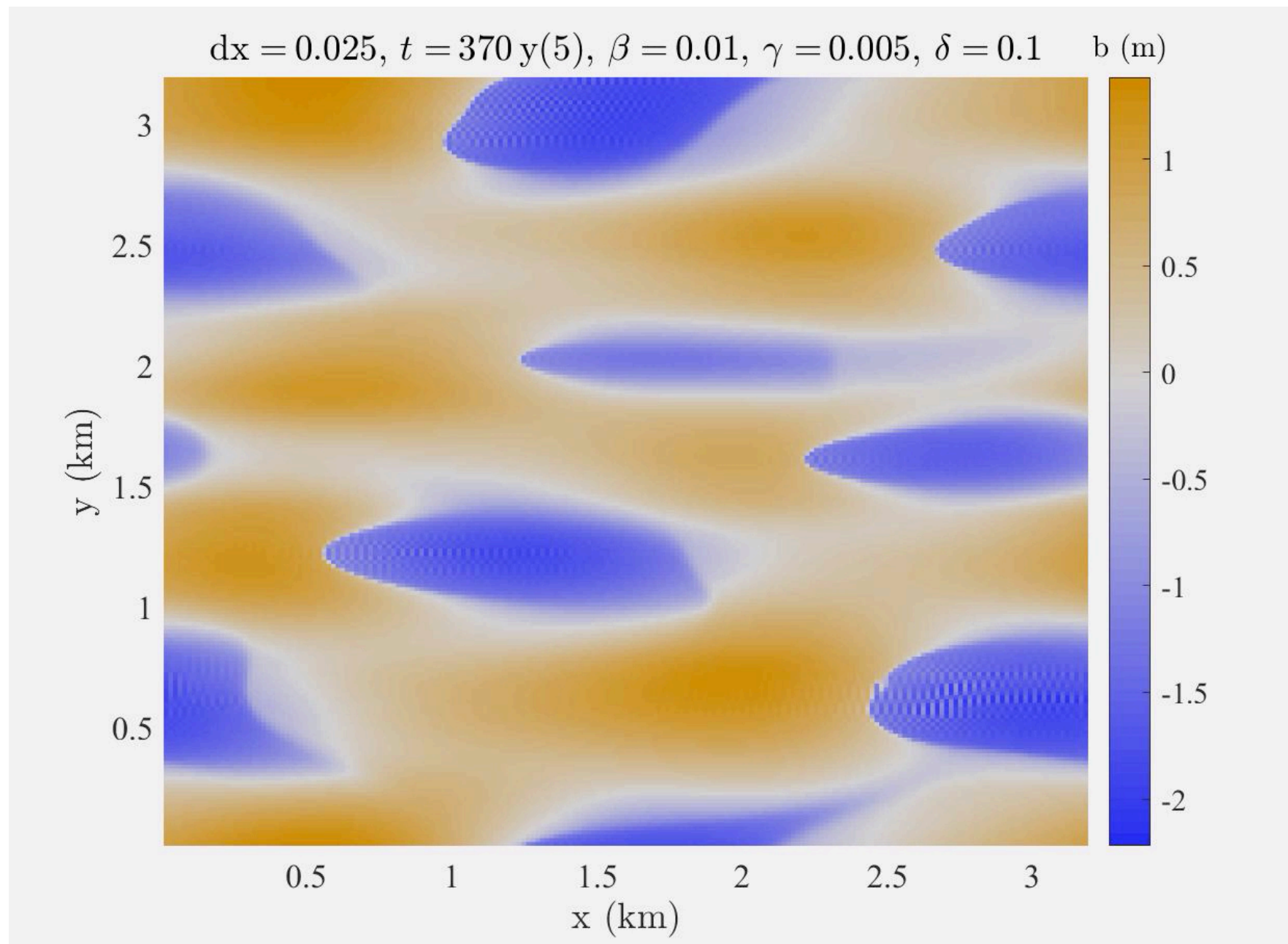
$$\left. \begin{aligned}
 \sigma(h^3)_x &= \nabla \cdot [h^3 \nabla \psi], \\
 \Psi &= s - N + \Phi, \\
 A &= \frac{1}{2} \left[\frac{f(\bar{u}, N)}{\mu} - N \right]_+, \\
 \varepsilon r h_t &= 1 - \Pi h N, \\
 b &= s - \delta h, \\
 b_t + \nabla \cdot [A \bar{u} \mathbf{i}] + \sigma \gamma [B(\tau_e) h]_x &= \beta \nabla \cdot [A^3 \nabla N] + \gamma \nabla \cdot [B(\tau_e) \{h \nabla \Psi + \theta \nabla b\}], \\
 \tau_e &= \sigma h \mathbf{i} - \{h \nabla \Psi + \theta \nabla b\} \\
 \alpha s_t + \bar{u} s_x &= w(\Phi, N).
 \end{aligned} \right\} \quad (2.59)$$

- Ice flow over sediments causes **transverse dune-like** instability.
- Water flow carrying sediments causes **longitudinal rill-like** instability.

Growth rates

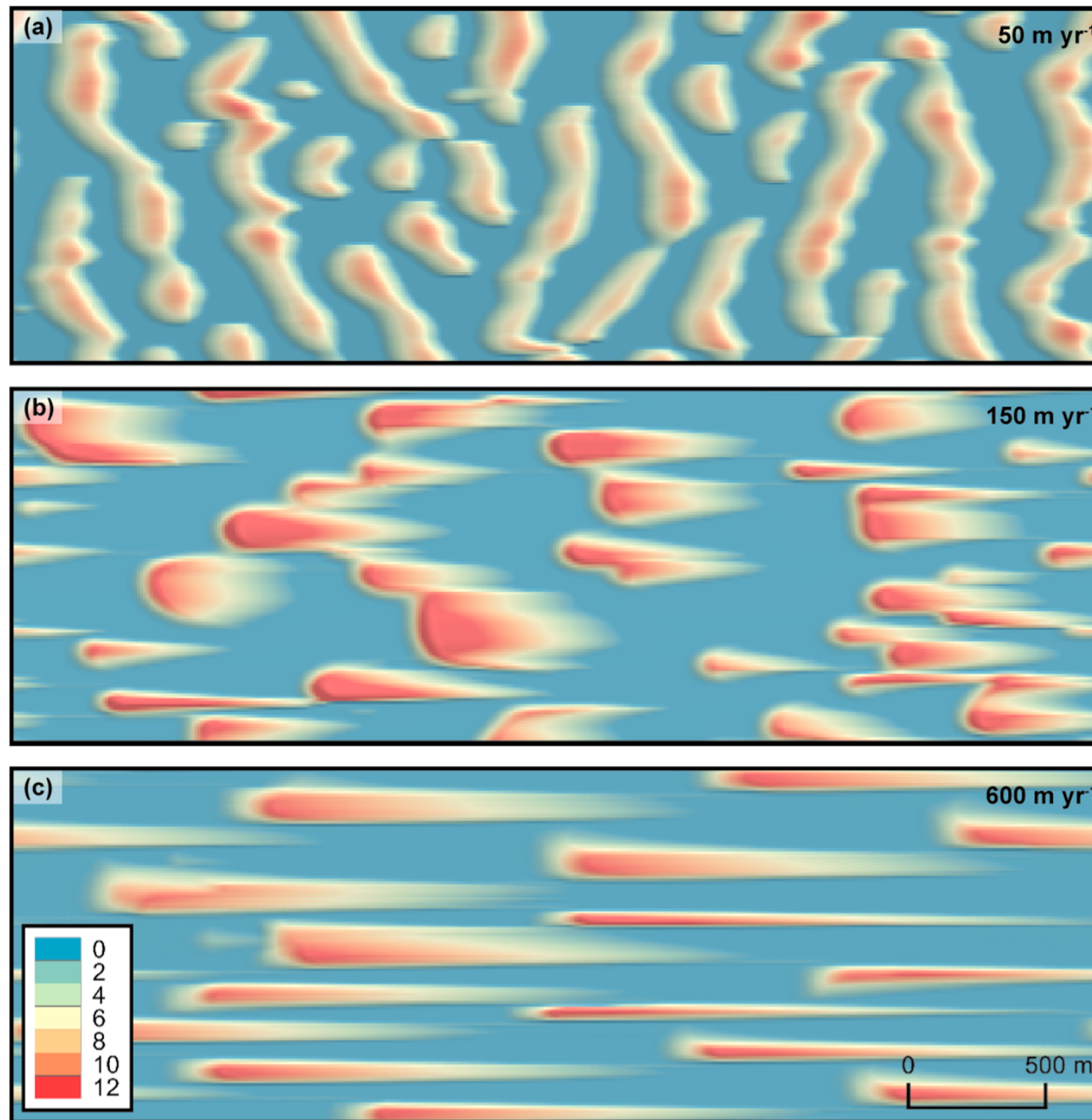


Instability theory



Fannon et al 2017

Instability theory



Barchyn et al 2016

Summary

Friction laws for soft beds are **similar to hard-bed sliding laws** (even though the local slip / deformation mechanism may be different).

Drainage over till may occur through **films, cavities** and **canals**.

Eskers form through deposition in Röthlisberger channels.

Ice flow over **deforming till** can be unstable and produce **ribbed moraine, drumlins**, and **mega-scale glacial lineations**.

An important open question is how the development of these bed-forms controls/affects ice dynamics.