

Week 7. Chris Beem

VOA(M4) and Vafa-Witten.

State of the art work of work by
Feigin - Gukov; Gukov + collaborators (Physics)

Dylan Butson | more mathematical.

— connection to lots of other things in String QFTs.
 $q = e^{2\pi i \tau}$, $\tau = \frac{\theta}{2\pi} + \frac{4\pi i}{g_{\text{YM}}^2}$
Recl: $Z_{\text{vw}}^G(M_4) = \sum_n VW(G, M_4, n) q^n$
4-manifold $C_2(G\text{-bundle} \rightarrow M_4)$

i) expected to satisfy some modularity properties
predicted from Physics.

— General state of modularity prediction, e.g. including
4-manifolds with boundary, v pretty vague.

Claim: these q -series should be characters of
(module, for) some vertex algebras,
associated to (G, M_4) , that are nice
enough to get modularity.
e.g. C_2 cofinite, or quasi-lisse, vertex algebras.

Why should we expect this?

Claim: $\exists$ a family (labelled by a simply-laced Lie algebra an $\in ADE$) of Gd QFT (2,0-theories)

Large spacetime symmetry broken: $OSp(8|4)$ superconformal algebra.

— Mysterious theories! Argued to exist w/o any other mysterious things (like string / m theory).

Don't have a Lagrangian / path-integral description.

— Using string theory (m-theory tricks, various "facts" known:

after compactifying to $d < 6$, can often identify resulting lower-dimensional theory as something more conventional. Gives geometric organization of some collection of lower-dim theories.

Early 2000's: industry of guessing lower-dim theories.

(2,0)_{any} $\rightarrow$ $X_5 \supset \int_R$ $\xrightarrow[\text{choice, here}]{R \ll \text{Vol}(X_5)}$ "5d Sym theory with gauge group G , Lie $\mathfrak{g} = \mathfrak{g}$ choice, here. $\uparrow$ mysterious

$(g_{\text{YM}}^2 \sim R)$ $\xrightarrow{\text{choice, here}}$ mysterious coupling constant.

elliptic curve param τ

4d $N=4$ Sym
with gauge group
 $G \leftarrow$ choice here.

$$(2,0)_g = X_4 \times \overline{E}_\tau \xrightarrow{\text{choice here}} \text{vol}(\overline{E}_\tau)^{1/2} \ll \text{vol}(X_4)^{1/4}$$

$$\frac{\theta}{2\pi} + \frac{4\pi i}{g_{\text{YM}}^2} = \tau$$

$\uparrow$
complexified
(coupling) constant.

"Geometric S-duality": $\overline{E}_\tau$ depends on τ

also $\uparrow$ to $PSL(2, \mathbb{Z})$ action on $\mathbb{H}$;

so expect r.h.s. to depend on τ also $\uparrow$ to $PSL(2, \mathbb{Z})$ too.

$$(2,0)_g = X_4 \times C_{g,n} \xrightarrow{\text{vol}(C_{g,n})^{1/2} \ll \text{vol}(X_4)^{1/4}} \text{Theory of de R S}.$$

$\uparrow$
Mysterious

$\uparrow$
Riemann
surface genus
 g, n moduli

$\uparrow$
Very
mysterious.

Gukov, Gaiotto / Dvali / Gaiotto, ... on gauges
for many other complicated lower dimensional limits.

Compatibility = a 2-d thing is easy to understand
2-manifolds very well. Compatibility = 3d things is
not complicated.

A simplification that is supposed to exist for these theories:

Twisting. (2,0) theory, has $\text{Spin}(5)_R$
 R -symmetry: preferred global symmetry of theory
 $\text{Spin}(5) = \text{Sp}(4)$.

— There is no topological twist of the 6-d theory.
 — But can restrict by hand to background of form $X_6 = X_4 \times X_2$

$$\text{Spin}(6) \supset \text{Spin}(4) \times \text{Spin}(2)$$

$$\text{Spin}(4) \xrightarrow{12} \text{Spin}(3) \times \text{Spin}(1)$$

$$\rho: \text{Spin}(6) \times \text{Spin}(1) \longrightarrow \text{Spin}(5)$$

R -symmetry group

is this twist: which properties of theory are trivialized in a cohomological sense?

— In this twist, X_4 becomes topological, holomorphic

Only feature of X_2 that you care about is holomorphic

"Holomorphic topological twist"

$(2,2)_g$ SCFT $\xrightarrow{\text{twist (hol/top)}}$ Holomorphic Topological
 field theory,
 labelled by g .

Explanation to $VOA(M_4)$ story:

$(2,2)_g$ HT $\sim$ Kähler $\mathbb{P}^1$
 $X_4 \times X_2$

$\text{vol}(X_4)^{1/4} \ll \text{vol}(X_2)^{1/2}$
 $\equiv$

use hol
surfs +
twist

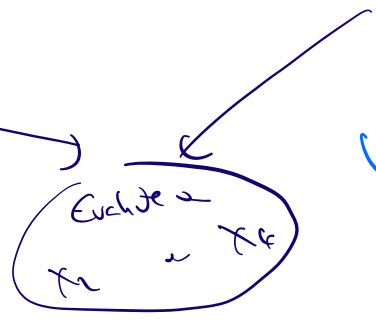
hol twist of

$(0,2)$ SCFT if X_4 gen

$(0,4)$ SCFT if X_4 Kähler.

$T(X_4, g)$
 evaluated $\sim X_2 / \mathbb{P}^1$

2-d theory



VW
twist

$\text{vol}(X_2)^{1/2} \ll \text{vol}(X_4)^{1/4}$
 $\approx$ there
 surfs +
 no twist.

4d $N=2$ SCFT if X_2 gen.

4d $N=4$ SCFT if $X_2 = \mathbb{P}^1$.

top twist of

$T(X_4, g)$
 evaluated $\sim X_4$

However, as theory is holomorphic / topological, answer
 should be independent of volumes j so
 all 3 things should be equivalent.

is Homomorphic twist of

$$(0, 2) \approx (0, 4) \text{ SFT}$$

Algebra of local operators,
becomes a vertex algebra.

Evaluating $\approx \langle E \rangle$ gives
the vacuum character of
vertex algebra.

- roughly

$$\text{VOA}(X_g)$$

Vafa-Witten
partition function

determinant / trace
 $q = e^{2\pi i \tau}$

$$\text{Compatibility} \approx X$$

= "taking mod of X
very small"

Key thing: after
taking twist, answer independent
of the size, so need not
take the limit.

Caveats / nuances in this picture.

- if 2-d theory $T(X_4, g)$ is nice enough (compact SCFT at low energies) then $\mathcal{L}$ holomorphic
- with of $T(X_4, g)$ has a natural, simple sub-VOA "left-moving chiral algebra":
— and this is what Feigin-Gukov call $VOA(X)_{FG}$
- And the full holomorphic algebra is a module for this.
- The theory they say is a module, is $clp = VOA$ in it, over right
- The full VOA may not be simple (can have paper ideal) Nice enough case.
- (Chi doesn't know: which X give nice enough theory; e.g. which compact X .)
— Feigin-Gukov don't restrict to "nice enough" X .

Two special classes of interesting cases:

Suppose X_4 non-compact (not equivariantly compact)



compact fixed point set, for some four action

extend to $M^3 \times \mathbb{R}^+$

$dM_4 = M_3$

(Question: can you do V-variant of $\mathbb{C}$ moduli with boundary?)

(2.0) HT theory
reduced = this

$T(M_3, \mathfrak{g})$ extended in
solid tori



$\Rightarrow$ solid tori partition further
Hilbert VA lives in d (solid tori)
inject the operator in center of tori.

— So: get lots of some moduli

expectation: solid tori that have
moduli symmetry.

$\mathbb{R}^1 \times \mathbb{R}^7 \leftarrow$ holomorphic
 $\leftarrow$ topological.
(HT)



$T(M_3, \mathfrak{g})$

$T(X_4, \mathfrak{g}) \leftarrow$ boundary condition
for $T(M_3, \mathfrak{g})$

Hence: VW invariant for G mod with boundary
 $\Rightarrow$ create moduli expectation.

fergus - Anton: $V-W$ invariant on
 G mod with boundary; need to show a
flat connection on the boundary. ($\Rightarrow$ fiber
neighborhood of mod?)

Case when things are most concrete, with
actual coordinates:

Suppose X is finite, $\mathbb{C}^n \times \mathbb{C}^n \rightarrow \mathbb{C}^n \times \mathbb{C}^n$

(fixed set is point)

expect to have a reasonable equivariant version
of the theory. $\Rightarrow$ Expect vertex algebra
with parameters (equivariant parameters) that you
can set to values.

equivariant localization?

$\downarrow$

Physics is the R -background

Expectation, because $i = 6-d$.

$$(Z_1^0)^{HT} \quad g = se_N^2$$

$$\mathbb{C} \times \mathbb{R}^4$$

$\underbrace{\varepsilon_1 \varepsilon_2}_{\text{equivariant under } (\mathbb{C}^*)^2}$

AGT correspondence

$\leadsto$ Beem - Rastelli - Van
Ree



Get

Vertex
algebra

at $h=1$

$\mathbb{C}$

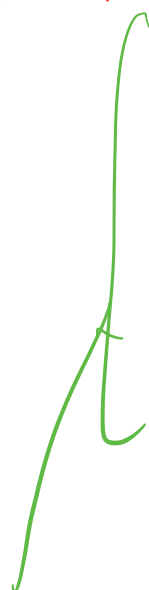
$\mathbb{C}$

vertex
algebra $\rightarrow$

W_N algebra with central
charge

$$C = (N-1) + (N^3 - N)(b + \frac{1}{b})$$

$$b^2 = \varepsilon_1 / \varepsilon_2$$



This is equivariant $VOA(m\epsilon)$

$$\text{of } (\mathbb{R}^4 \supset \mathbb{C}^{*2})$$

— For more details, see X_4 ,

let $\{p_1, \dots, p_n\}$ be the $U(1) \times U(1)$ fixed points

expect a) a Abelian algebra of Equivalent VOA $(X_E, \tau_1, \tau_2, S^1_N)$

(Feynman Graphs)

extension



periodic central charge.

$$\bigotimes_{i=1}^N W_N(C(\tau_{i,1}, \tau_{i,2}))$$

$i=1$

T

for: each W_N induces

$W_n(\dots)$

$\uparrow$
Heisenberg

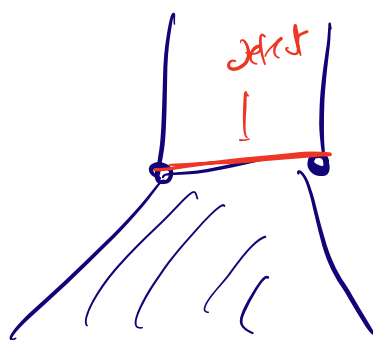
Bottom:

— Expect $\bigotimes_N VOA(X_E, \tau_{i,j})$ is a

extension of $\bigotimes_{i=1}^N W_N(\dots)$

(Bottoms think of X_E as divisor is (43)).

$$\mathcal{O}(-1) \otimes \mathcal{O}(0) \rightarrow \mathcal{P}$$



Bottom.

6-d theory which
is a 5-d theory on $\mathcal{O}^1$.

divisor: $\mathcal{O}(-2) \rightarrow \mathcal{P}$.

Physics question: which model of

$$\bigoplus_{n \in \mathbb{N}} W_n(c(\dots))$$

appears in $\text{VOA}(\mathcal{H}, \mathcal{E}, \mathcal{E}_n)$.

$$\bigoplus^T (\text{possibly many } \mathcal{H}_i \text{ of } \bigoplus W_n)$$

— Get prediction for the vector space

of $\text{VOA}(\mathcal{H}_{\infty})$ as a model for

$$\bigoplus W_n(\dots)$$

— The way to get VOA inside $\text{VOA}(\mathcal{H}_{\infty})$

$$\text{VOA}(\mathcal{H}_n)$$

$\supset$

$$\bigoplus W_n$$

$\uparrow$

characterize

$\uparrow$

characterize by symmetry

$$\text{by } (\text{ultra VOA}) = \bigcup (\mathcal{H}_n)^{\otimes n} \text{ VOA(mix } \mathcal{H}_n \text{ of } \mathcal{H}_n) ?$$

$ag = ag^L$: nice theory: Nakajima-Grojnowski?

One thing to advertise:

— (one which is not toxic, but should have nice structure if work equivariantly wrt. a single $V(\mathbb{A})$).

(Tanaka-Thomson: equivariant to make things work at all. BT-ask: what if the $\mathbb{C}^*$ -module X has a group action too?)

equivariantly equivariant.

Con. $X_{\mathbb{C}} = T^*G \hookrightarrow \mathbb{C}^*$ action & fibe.

Here $\text{VOA}(X_{\mathbb{C}})$ should be $\cong$ Chiral algebra of class S
 $\downarrow$
projected

T (contributed by

(2-d q -deformed
Yang-Mills partition fn.)

T . A. Nakawaj
very explicit character formula.

U(1) -equivariant theories: don't get extra
parameters.

What difference does X Kähler make?

Fejn Gulo $VOA(ME)$ $\Rightarrow$ (moduli).
 $\uparrow$
He looks in the VOA.

T

X Kähler gives

some extra information here

(defn) vector space of $VOA(ME)$? (in terms of
moduli space? But - Ryzik: say something.